

# Double Ionization of Atoms in Intense Laser Fields

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## Acknowledgements:

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- J.H. Eberly and R. Panfili of the University of Rochester
- Student assistants Z. Smith, A. Karim, L. Breen

# Background--Calvin College

- Calvin College is a comprehensive college of 4,000 students in Grand Rapids, MI
- Physics and Astronomy Department has 8 tenure-track faculty (7.5 postns)
  - 5 experimental physicists
  - 2 astronomers
  - 1 theorist
- Graduate ~5 physics majors per year. About half immediately go on to grad school.

# My background

- I started at Calvin in 1983, fresh from a postdoc
- Spent at least part of each of the first 3 summers back at Colorado, where I'd postdocked.
- Received my first grant for research at Calvin in 1987 (Research Corp + NSF-RUI), and hired my first undergraduate assistant in Summer 1987
- Continuous NSF support since then

# Professional Development

- Sabbatical 1992-93 to University of Rochester
  - Funded in part through NSF ROA program
  - Collaboration (with J.H. Eberly) still ongoing!
- Sabbatical also marked a transition from work that was analytically based to work that was more computer based.
  - Brought “supercomputer techniques” back to Calvin
  - Now do my work on desktop machines

# Usual Procedures

- I like to hire 2 students each summer.
- Ideally, one of the students is experienced and can help the other
- The students work on closely related threads of the same project
- I challenge each student to take ownership of the thread, and to see it through to completion by working on a part-time basis through the academic year

# The Motivation

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PHYSICAL REVIEW LETTERS

29 AUGUST 1994

## Precision Measurement of Strong Field Double Ionization of Helium

B. Walker,<sup>1</sup> B. Sheehy,<sup>1</sup> L. F. DiMauro,<sup>1</sup> P. Agostini,<sup>2</sup> K. J. Schafer,<sup>3</sup> and K. C. Kulander<sup>3,4</sup>

Experiments in the early 1990s showed considerably more double ionization occurring than would be expected for independently behaving electrons.

- Somehow the electrons share energy and escape together

first ionization 16 photons;  
second 34 more

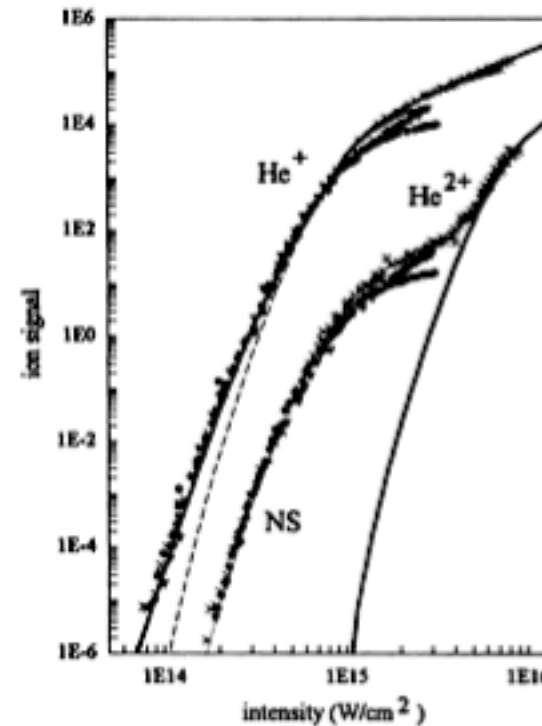


FIG. 1. Measured He ion yields for linear polarized, 100 fsec, 780 nm light. Calculations are shown as solid (SAE) and dashed (ac-tunneling) lines. The measured intensities are multiplied by 1.15. The solid curve on right is the calculated sequential He<sup>2+</sup> yield.

# The Challenge

We can't just solve the 3-d Time-Dependent Schrödinger Equation for two electrons in an oscillating electric field, even on the best computers.

Instead, we (& several other groups) worked with a one-dimensional model of the helium atom

- In the model, each electron could move only parallel or antiparallel to the laser polarization

Picture the electrons as being on railroad tracks

- A full quantum solution is possible.
- We learned a lot,  
including that the electron behavior seemed very classical.
- So we tried classical modeling of the one-dimensional system, and we found that the classical results reproduced all the important quantum results.

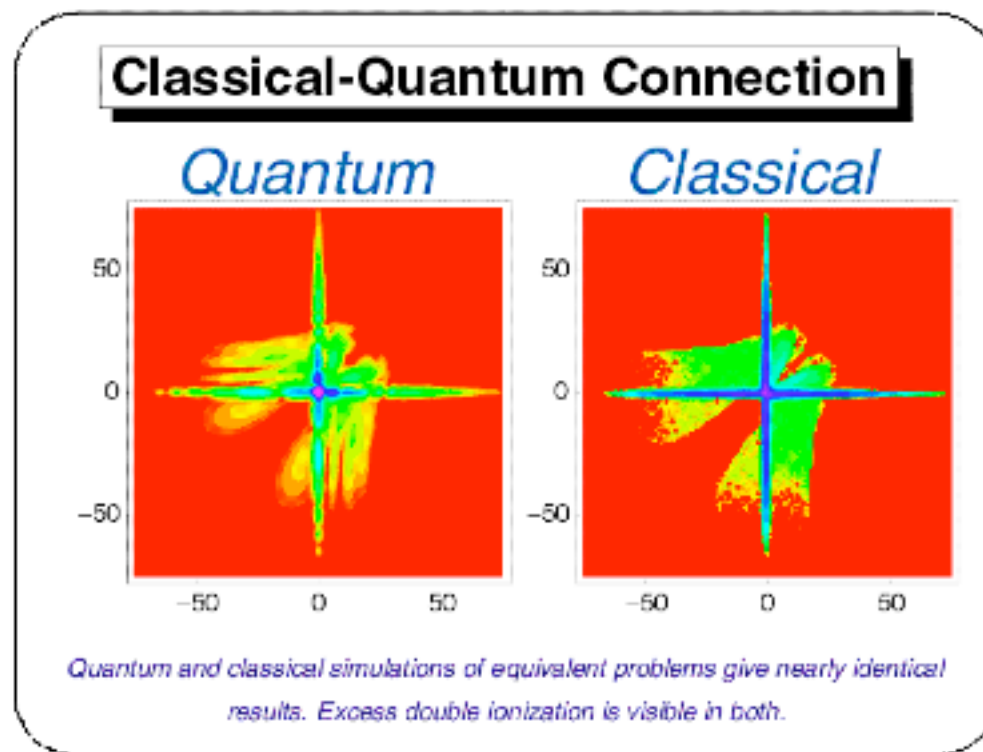
# Overview of Classical Method

- We set up an ensemble of classical two-electron atoms.
  - Each atom has slightly different initial conditions, so that collectively they imitate the quantum ground state of the atom.
  - Ensemble sizes typically 400,000.
- We evolve each two-electron atom in time through a laser pulse, using Newton's laws of motion.
  - Account for electron-electron interaction, electron-nucleus interaction, and influence of laser field (in dipole approximation).
- After each run, we can sort the trajectories and
  - Study statistical behavior;
  - Backtrack individual trajectories to learn their history.



# Classical simulations

- Classical ensembles exhibit behavior very similar to the quantum systems.
- Example below shows population distributions at a particular time.



R. Panfili, J.H. Eberly and S. Haan, Opt. Express 8, 431-435 (2001).

# Extension to 3D

The success of 1D classical simulations gives us motivation to pursue the 3D classical case.

The classical advantage--

In classical studies we find the time development for each classical atom separately.

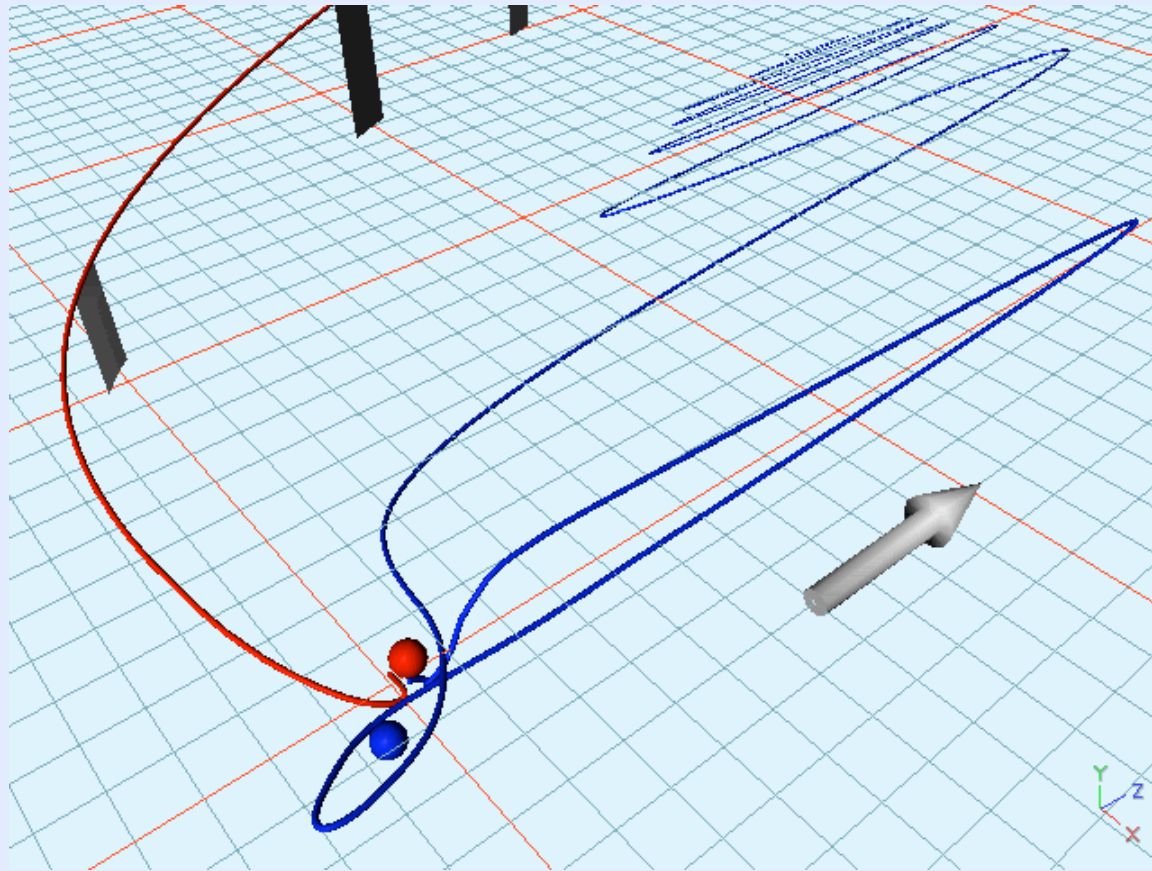
The QM wavefunction incorporates, at each time, everything that can happen

R. Panfili (Rochester grad student) wrote a first draft of a computer program for the 3d classical case.

He graduated & took a job in industry without getting the kinks out of it.

My students and I got the kinks out

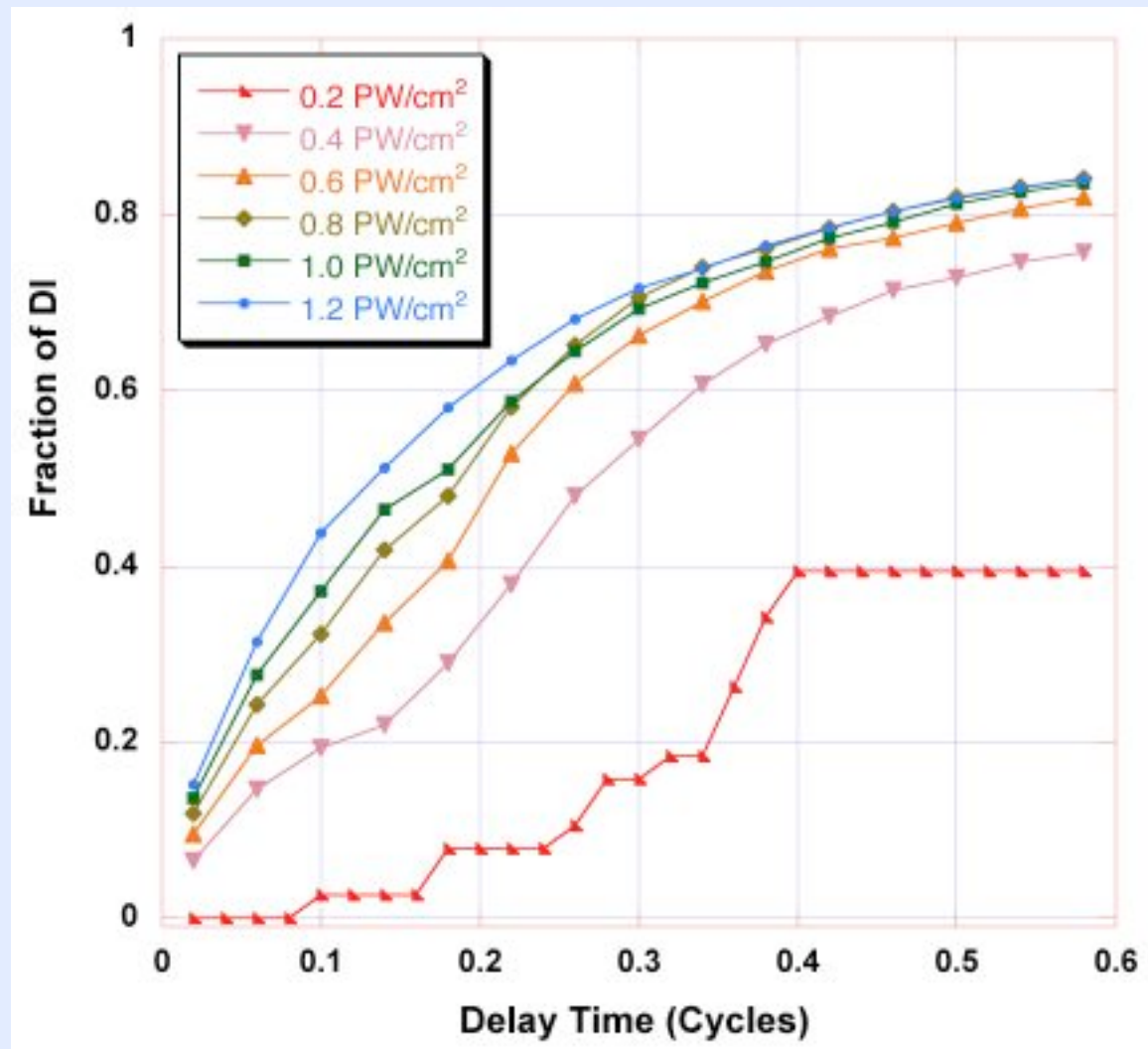
Example double ionization trajectory :  
Recollision ionization, with brief recapture of recolliding  
electron



Recollision had been proposed from the outset (by others), but most everyone overlooked this possibility of having one electron bound for a portion of a cycle after recollision.

# Delay time between recollision & double ionization

- Most DI trajectories show a part-cycle phase delay between recollision and double ionization
- Runs at right are all for laser wavelength 780 nm.



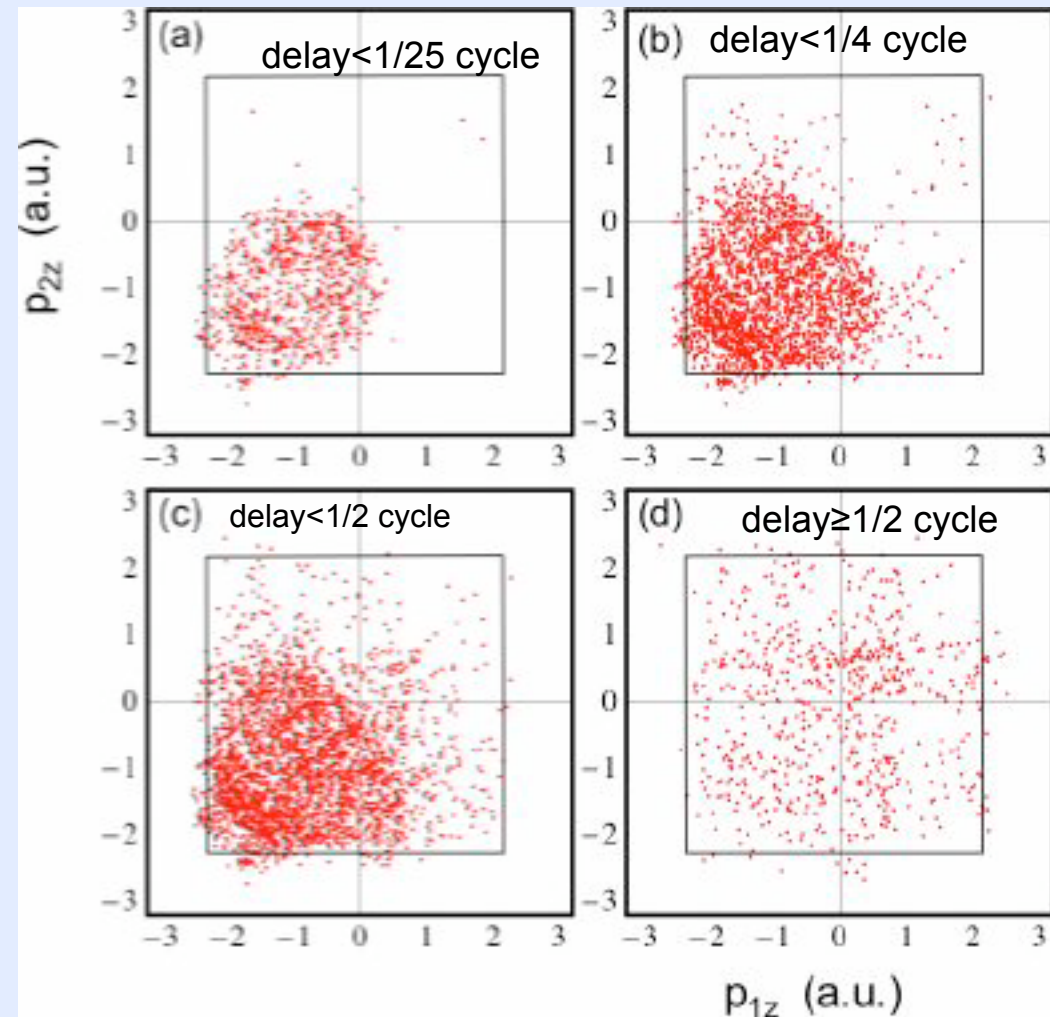
Final longitudinal momenta sorted by:  
delay times from recollision to ionization  
**and** by final direction relative to recollision direction

$I=6 \times 10^{14} \text{ W/cm}^2$

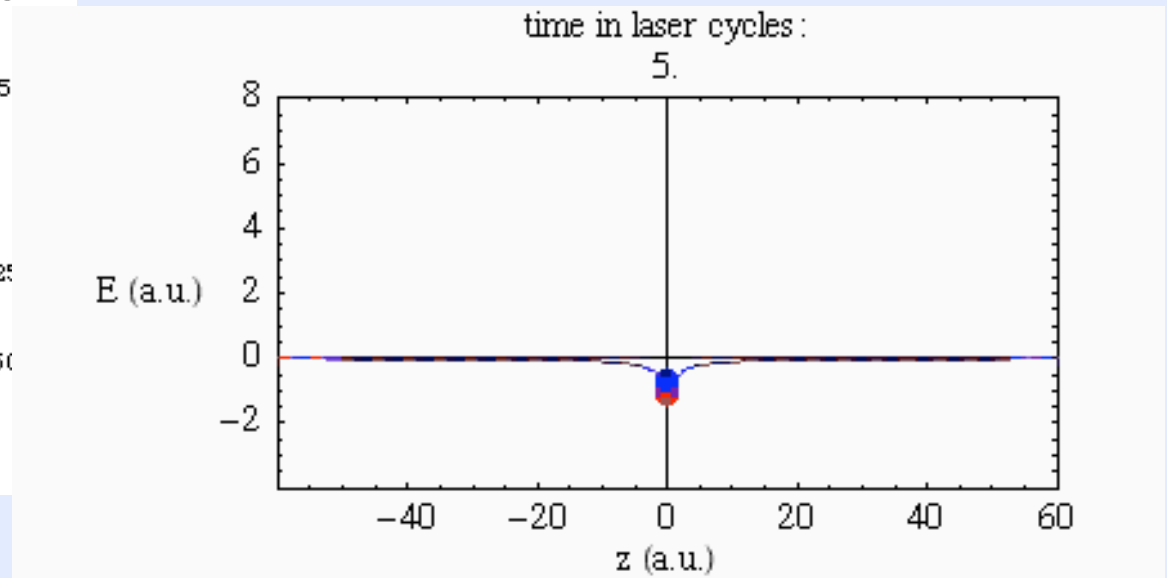
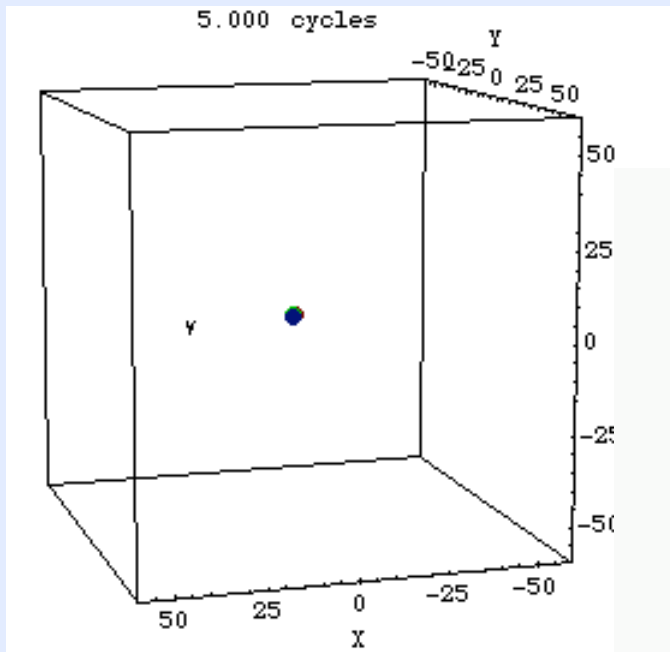
RE directions--adjust signs of momenta so all collisions occur with returning electron traveling in the +z direction.

- For small delay times, almost all final z-momenta are *opposite* from the recollision direction.

- With increased delay times, there is increased spillover into the 2nd and 4th quadrants.



- If, to first approximation, second electron ionizes after the field peaks, the electrons can have drift velocities in opposite momentum hemispheres.

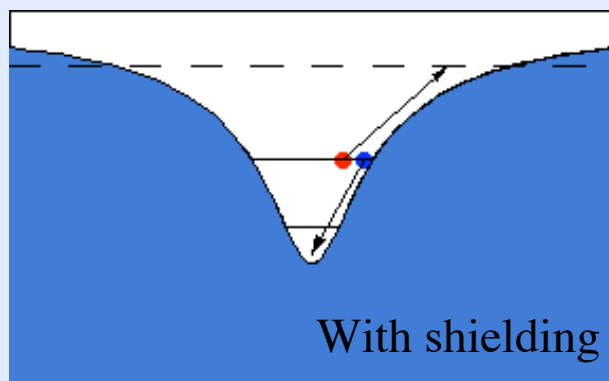
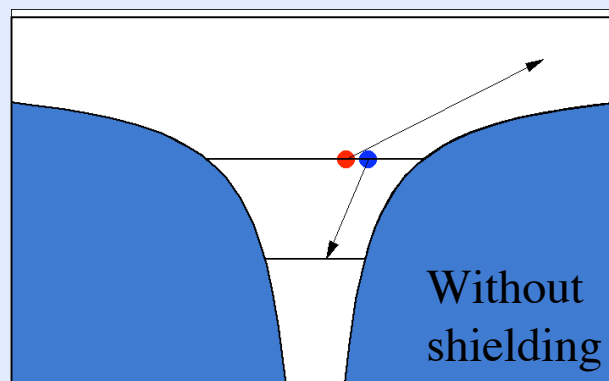


- These results all summarized in two recent papers
  - S.L. Haan, L. Breen,\* A. Karim,\* and J.H. Eberly, “Recollision Dynamics and Time Delay in Strong-Field Double Ionization,” *Optics Express* **15**, 767-778 (February 2007).
  - S.L. Haan, L. Breen,\* A. Karim,\* and J.H. Eberly, “Variable Time Lag and Backward Ejection in Full-Dimensional Analysis of Strong-Field Double Ionization,” *Phys. Rev. Lett.* **97**, 103008 (8 September 2006). Republished in *Journal of Ultrafast Science* **5**, Issue 10 (October 2006).



# A challenge for 3D: Classical 3-D Helium is unstable, even without external fields

- One electron can dive deep into potential-energy well, letting the other escape.
- The well is bottomless if use Coulomb potentials.
- Can stabilize by screening the Coulomb interaction:  
$$1/r \rightarrow 1/\sqrt{r^2 + a^2}$$
- For initial energy = -2.9035 au, setting  $a \geq 0.7$  au prevents self-ionization
  - Results I've shown  $a=0.825$  au



Plots show only the nuclear.

# Populating initial ensemble to mimic 3D ground state?

There are various possibilities related to energy and angular momentum.

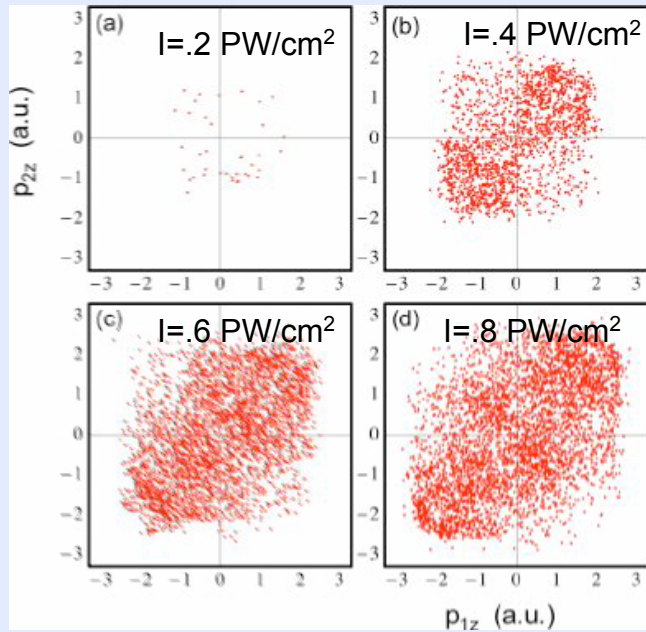
We simply:

- Set electron-nuclear shielding at 0.825 au, and electron-electron shielding at 0.05 au
- Fixed the energy at -2.9035 au (He ground state)
- Choose classically allowed positions
  - Employed several methods--He quantum ground state, Gaussian, random
  - but always spherically symmetric
- Randomly divided the available KE between the two electrons
  - (subject to constraint that each electron have total energy  $< -1$  au)
- Gave each electron zero angular momentum.
  - (If don't constrain angular momentum, we get very little double ionization.)
- Allow the system to evolve for the equivalent of one laser cycle ( $\sim 100$  au) with no laser field
  - Total energy and total angular momentum conserved
  - Position and momentum distributions stabilize

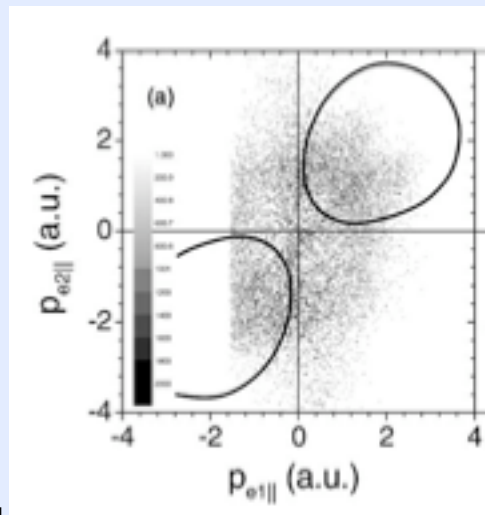
# A difficulty

- The screening of the nucleus works well for stabilizing our initial state
  - It can be justified on the basis of the uncertainty principle--the system starts in a state with low uncertainty for energy and momentum, and electron isn't allowed to know very precisely its position relative to the nucleus
- But what about at recollision?
  - We've found that *details of the final electron momentum & energy distributions depend on the form of the interaction with the nucleus*
    - One electron can scatter off nucleus at recollision
      - This electron may be either free or still bound
    - Small impact parameter  $\Rightarrow$  large deflection, if nucleus unshielded
  - In order to reproduce experimental results, we need to partly unscreen the nucleus.

# Our results:



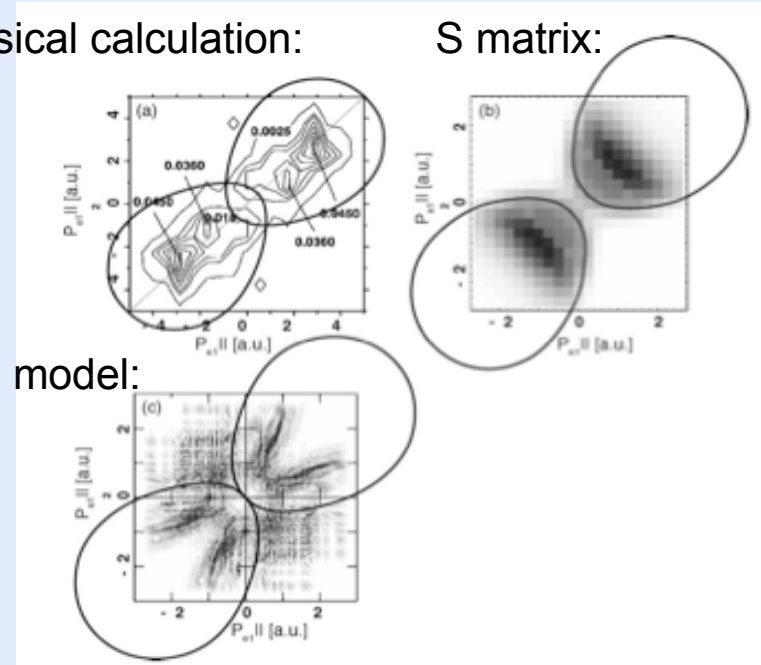
## Experiment (I=.5 PW/cm²)



S.L. Haan, Calvin College

# Other theoretical results

Classical calculation:



1-d model:

Egg-like shapes show kinematically allowed regions after recollision ionization

Experimental results--V.L.B. de Jesus, et al., Journal of Electron Spectroscopy 141, 127 (2004);  
 theoretical--J. Chen et al PRA 63,011404R (2001); A. Becker and F.H.M. Faisal, PRL89, 193003 (2002); Lein et al PRL 85, 4707 (2000)

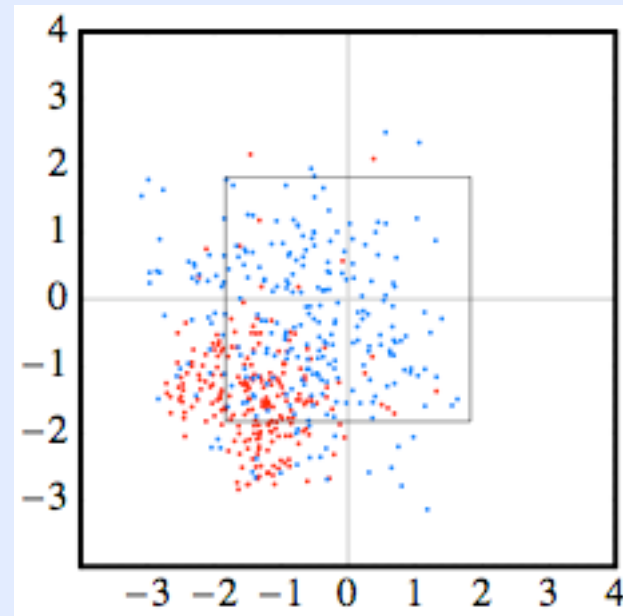
# Current work

- This summer my students and I have revised our program with a “toggle switch” on the nuclear screening

–When one electron ionizes, we reduce the nuclear shielding for both electrons

- We adjust the electrons' radial kinetic energies to preserve total energy

-Now we get trajectories with higher final energy.



# Present Status

- We're systematically adjusting the shielding at return
- I'm working on a "dynamic shielding," in which the screening of the nucleus for each electron depends on the energy of that electron
  - How, precisely, should the screening depend on the energy?
  - Should I add terms to the equation of motion to account for the "velocity dependent potential"?
- Research projects can lead us into new areas we know nothing about!