

Quasi-linear wave-mean flow interactions in large-scale planetary circulations

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The quasi-linear approximation

Quasi-linear (QL) model

(for quadratic non-linearities, for example in barotropic, QG, primitive equations....)

Retained:

- Wave-mean flow interaction (mean flow = zonal average)
- Interaction of waves of opposite zonal wavenumber (Reynolds stress)

Neglected:

- All other eddy-eddy interactions

The quasi-linear approximation

Advection of a quantity $a = \bar{a} + a'$ by a meridional flow $v = \bar{v} + v'$ (zonal average/eddy decomposition):

$$\frac{\partial a}{\partial t} = -\bar{v} \frac{\partial \bar{a}}{\partial y} - \bar{v} \frac{\partial a'}{\partial y} - v' \frac{\partial \bar{a}}{\partial y} - v' \frac{\partial a'}{\partial y}$$

Equation for the mean flow unchanged

$$\frac{\partial \bar{a}}{\partial t} = -\bar{v} \frac{\partial \bar{a}}{\partial y} - v' \frac{\partial a'}{\partial y}$$

Equation for the eddies linearized around the mean flow

$$\frac{\partial a'}{\partial t} = -\bar{v} \frac{\partial a'}{\partial y} - v' \frac{\partial \bar{a}}{\partial y} - \left(v' \frac{\partial a'}{\partial y} - \bar{v}' \frac{\partial a'}{\partial y} \right)$$

Quasi-Linear
Approximation

Why a quasi-linear approximation?

- Truncation that captures wave mean flow interactions while conserving 1st and 2nd order inviscid invariants (see Wanming Qi poster)
- Linear framework useful for **analytical theories** (*SSST of Farrell and Ioannou; Zonostrophic instability of Srinivasan and Young, 2012; Kinetic theory of Bouchet and coauthors, 2013, see Cesare Nardini poster*)
- First step to develop **statistical closures** (QL equivalent to a second order closure, *CE2 of Marston and Tobias*)

Why a quasi-linear approximation?

- QL works in **barotropic flows** when there is **time scale separation** between mean flow and eddies and **no unstable modes** (e.g. Bouchet and coauthors, see Cesare Nardini poster)
- “Quasi-linear reasoning” ubiquitous in the large scale dynamics literature. Large-scale atmospheric flows look linear over a wide range of parameters (Schneider and Walker, 2008). Some theoretical explanation: Schneider and Walker, 2006.
- **How does QL work in large-scale planetary baroclinic flows with large-scale forcing and “strong” surface friction?**
(Whitaker and Sardeshmukh, 1997; Zhang and Held, 1999; Delsole, 2001)

An idealized dry GCM

- FMS GFDL pseudospectral dynamical core
 - **Radiation:** Newtonian relaxation toward temperature profile
 - **Convection:** relaxation of vertical lapse rate toward $0.7 \times (\text{dry adiabatic})$
 - Diffusion of momentum and dry static energy in the **boundary layer**
 - Uniform surface, no seasonal cycle
 - Run at T85 (256×128) with 30 vertical sigma-levels
 - 600 days average after 1400 days spin-up
- OR
- **Radiation and convection:** Newtonian relaxation to a statically stable temperature profile, (Held-Suarez)
 - Rayleigh drag in the **boundary layer**

(Held and Suarez, 1994; Schneider and Walker, 2006)

Removal of the eddy-eddy interactions in the GCMs

At each time step:

Advection of a quantity $a = \bar{a} + a'$ by a meridional flow $v = \bar{v} + v'$ (zonal average/eddy decomposition):

$$\frac{\partial a}{\partial t} = -v \frac{\partial a}{\partial y} = -\bar{v} \frac{\partial \bar{a}}{\partial y} - \bar{v} \frac{\partial a'}{\partial y} - v' \frac{\partial \bar{a}}{\partial y} - \boxed{v' \frac{\partial a'}{\partial y}}$$

transformed into

$$\frac{\partial a}{\partial t} = -v \frac{\partial a}{\partial y} = -\bar{v} \frac{\partial \bar{a}}{\partial y} - \bar{v} \frac{\partial a'}{\partial y} - v' \frac{\partial \bar{a}}{\partial y} - \boxed{\bar{v}' \frac{\partial a'}{\partial y}}$$

Mean Zonal circulation

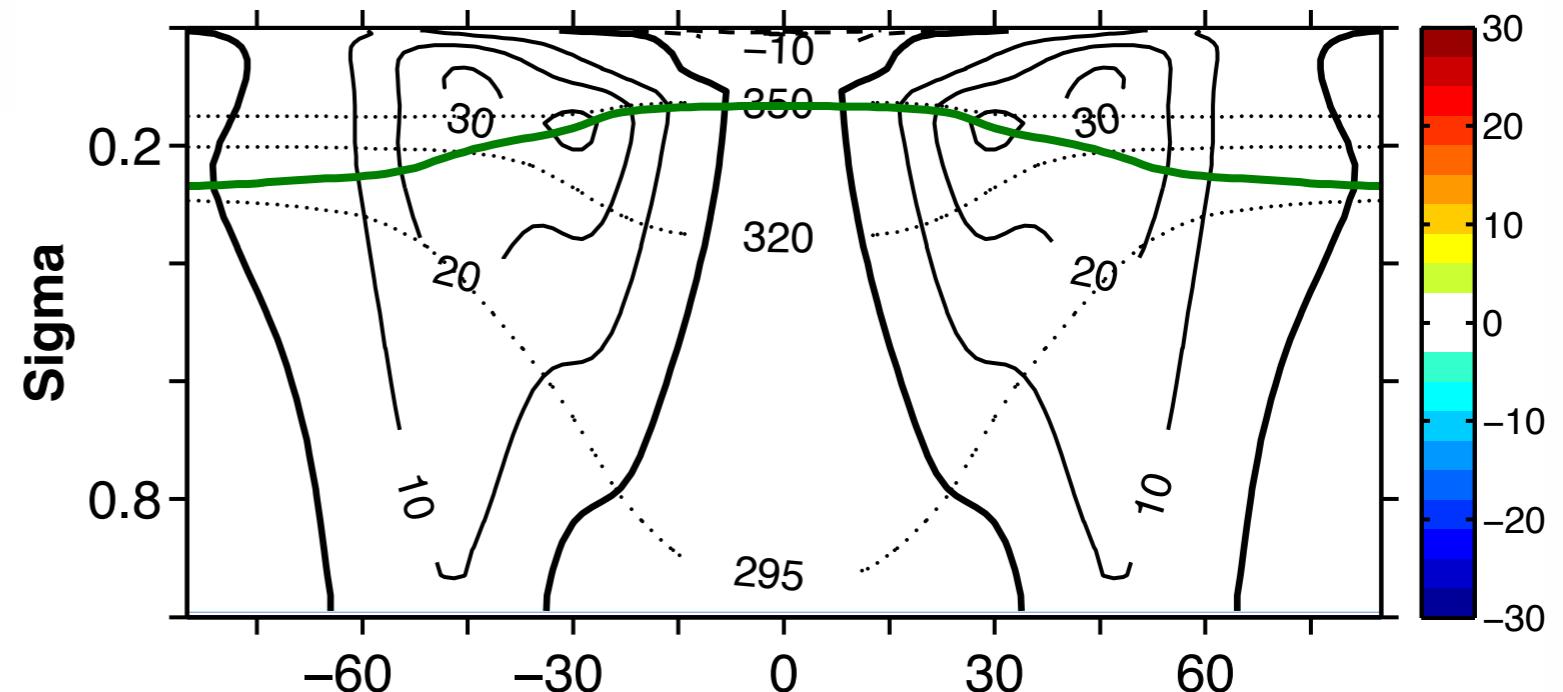
O'gorman and Schneider, 2007

Full

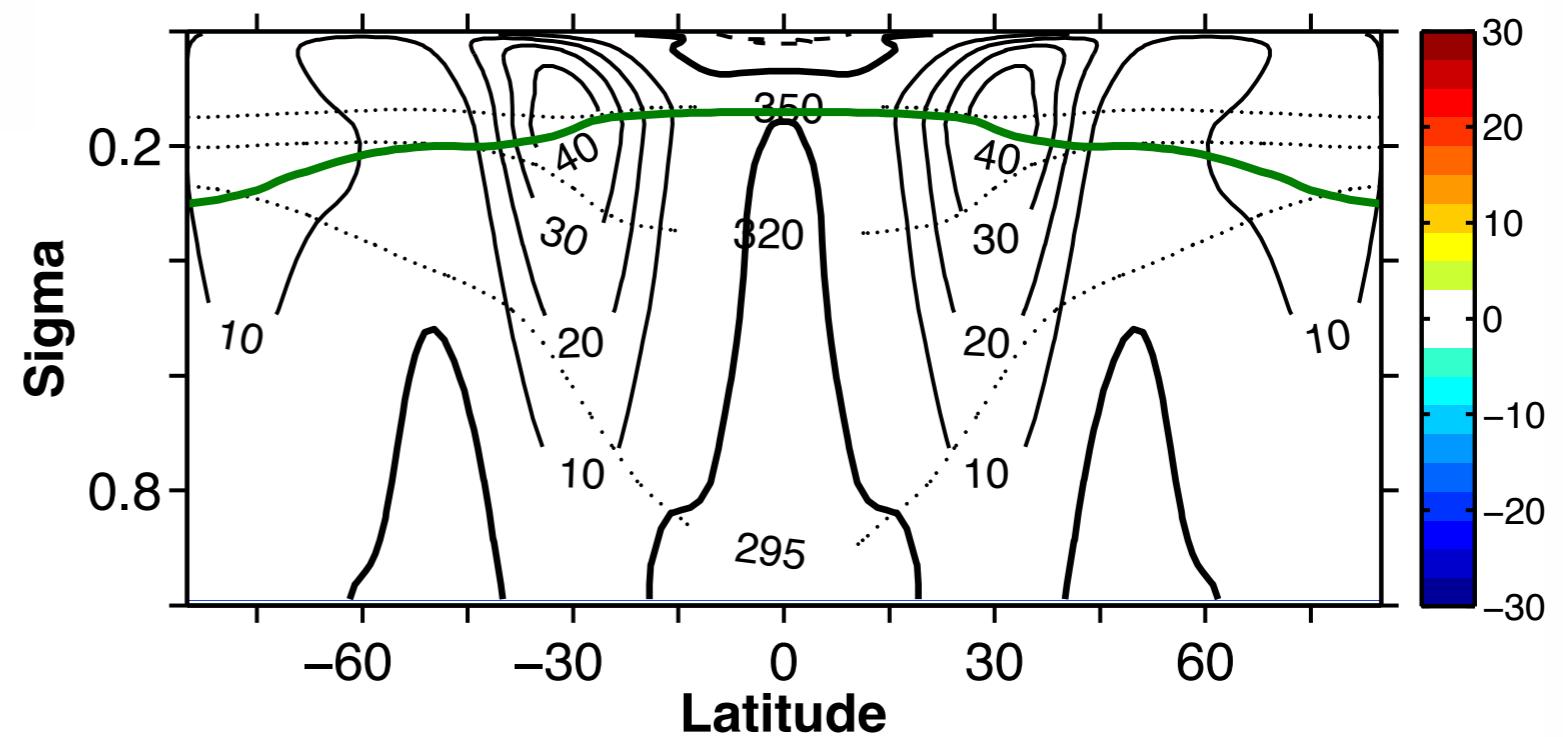
Contours:
zonal mean flow
(m/s)

Dotted lines:
potential
temperature (K)

Grey line:
tropopause



QL



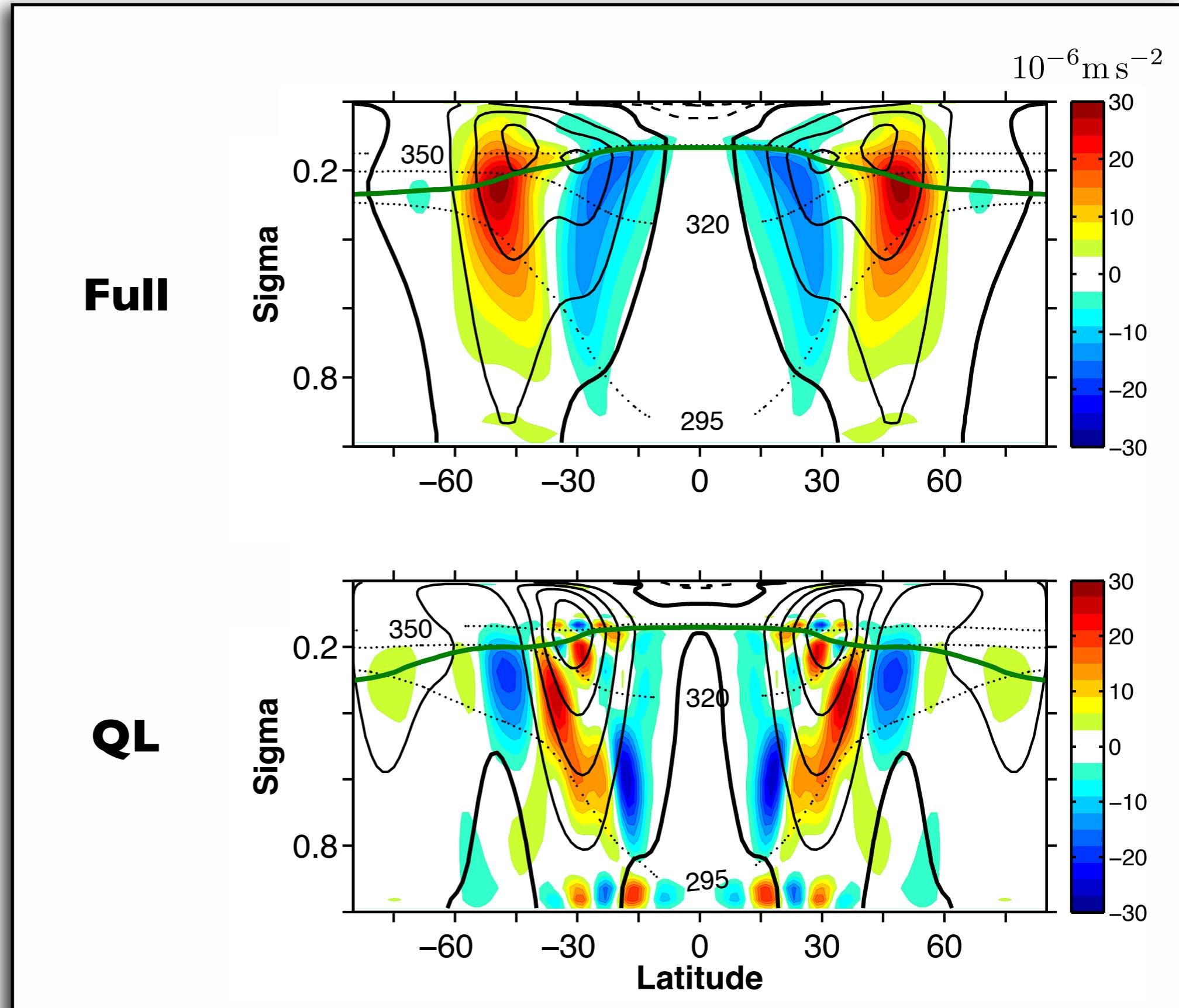
Eddy momentum flux convergence

Colors:
convergence of
eddy momentum
flux $\overline{u'v'} \cos \phi$

Contours:
zonal mean flow
(m/s)

Dotted lines:
potential
temperature (K)

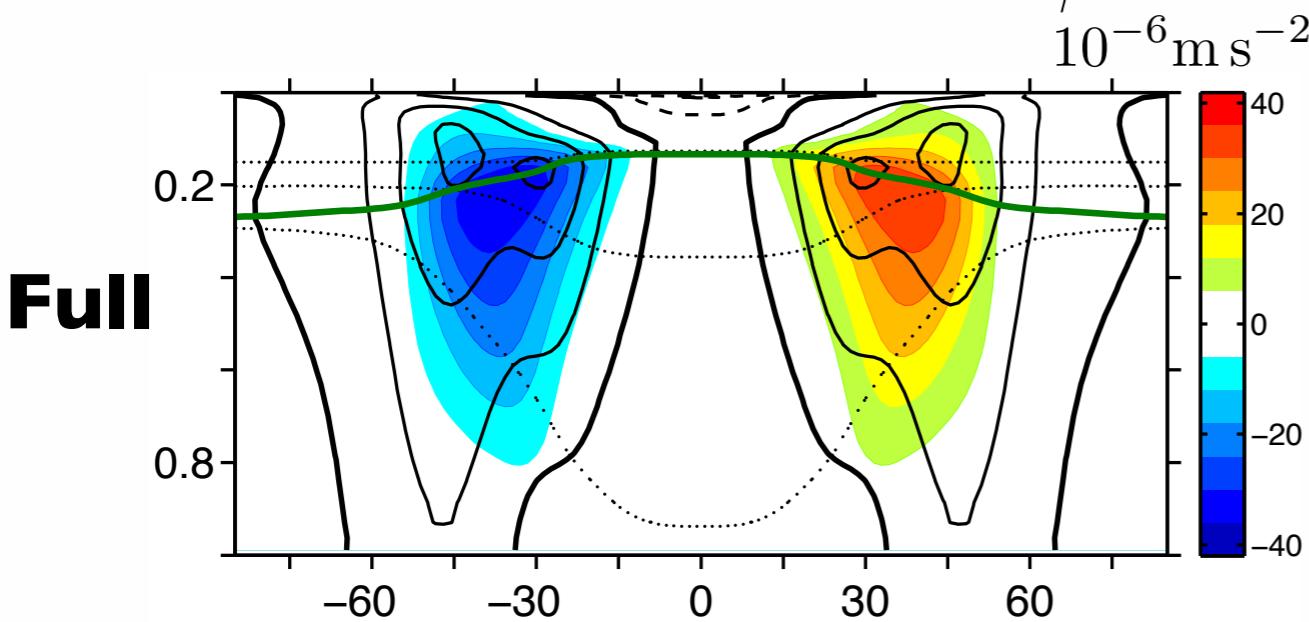
Green line:
tropopause



Momentum flux and eddy kinetic energy

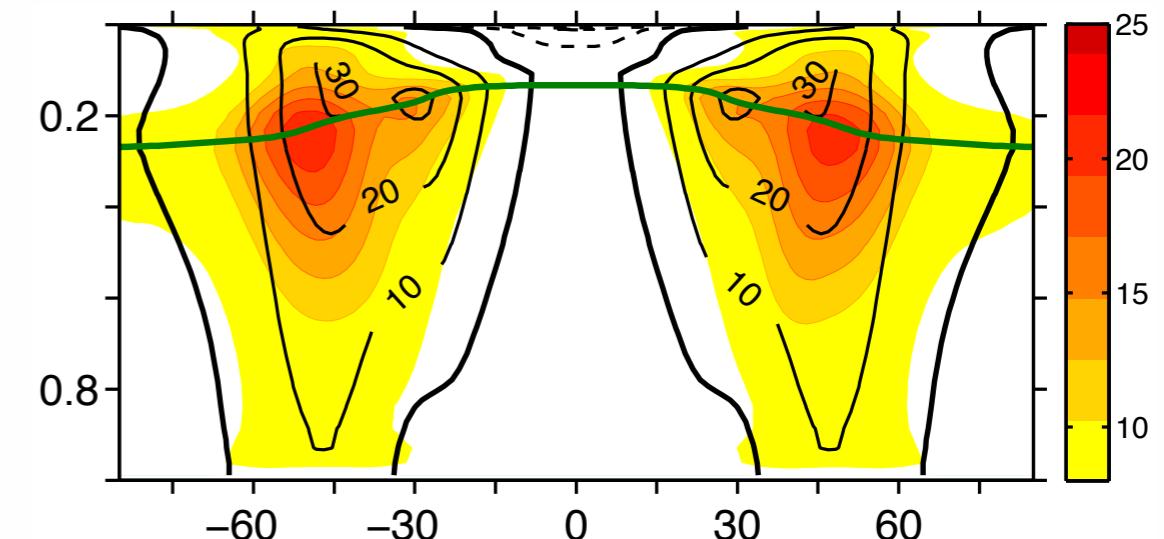
Contours: Zonal wind \bar{u} in m/s

Colors: momentum flux $\overline{u'v'} \cos \phi$

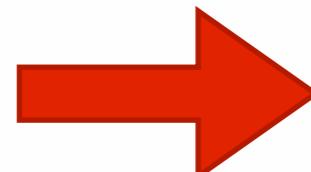
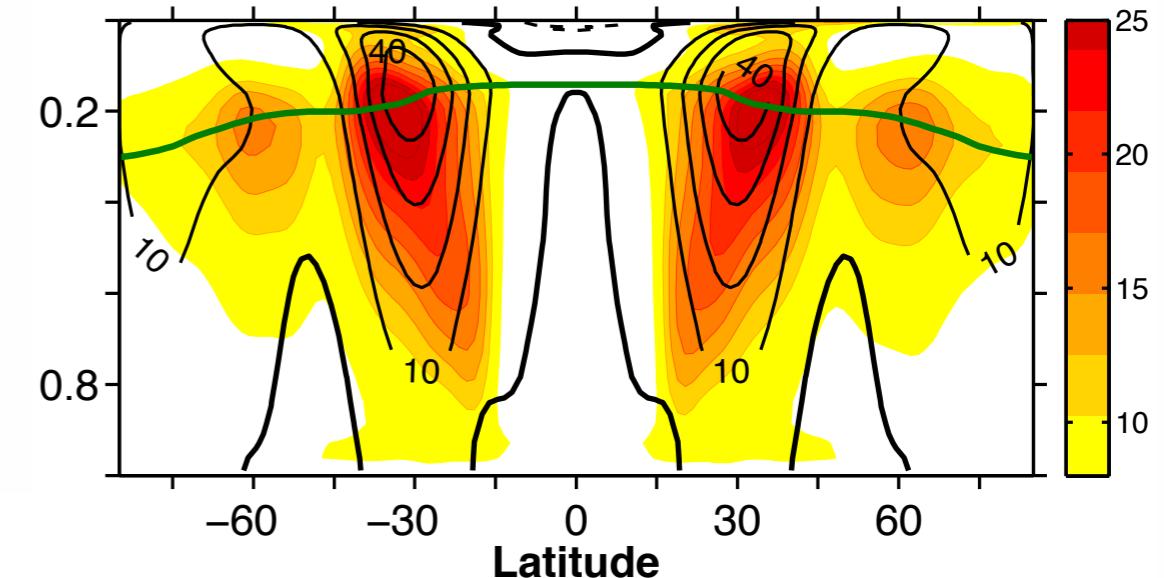
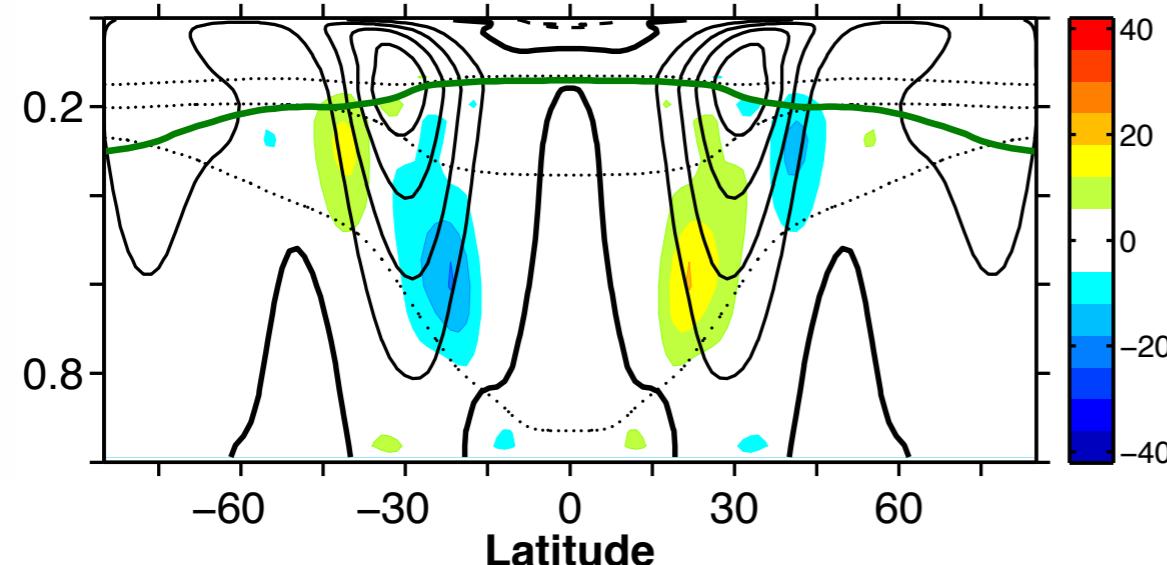


Contours: Zonal wind \bar{u} in m/s

Colors: Square root of EKE $(\overline{u'^2} + \overline{v'^2})^{1/2}$



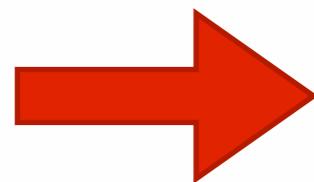
QL



- Momentum flux not enhanced in the upper troposphere
- Weak momentum flux associated with strong EKE in the upper troposphere

- Narrow, strong jets
- Momentum flux not enhanced in the upper troposphere
- Weak momentum flux associated with strong EKE in the upper troposphere

Dynamics ???



Baroclinic wave lifecycle experiments

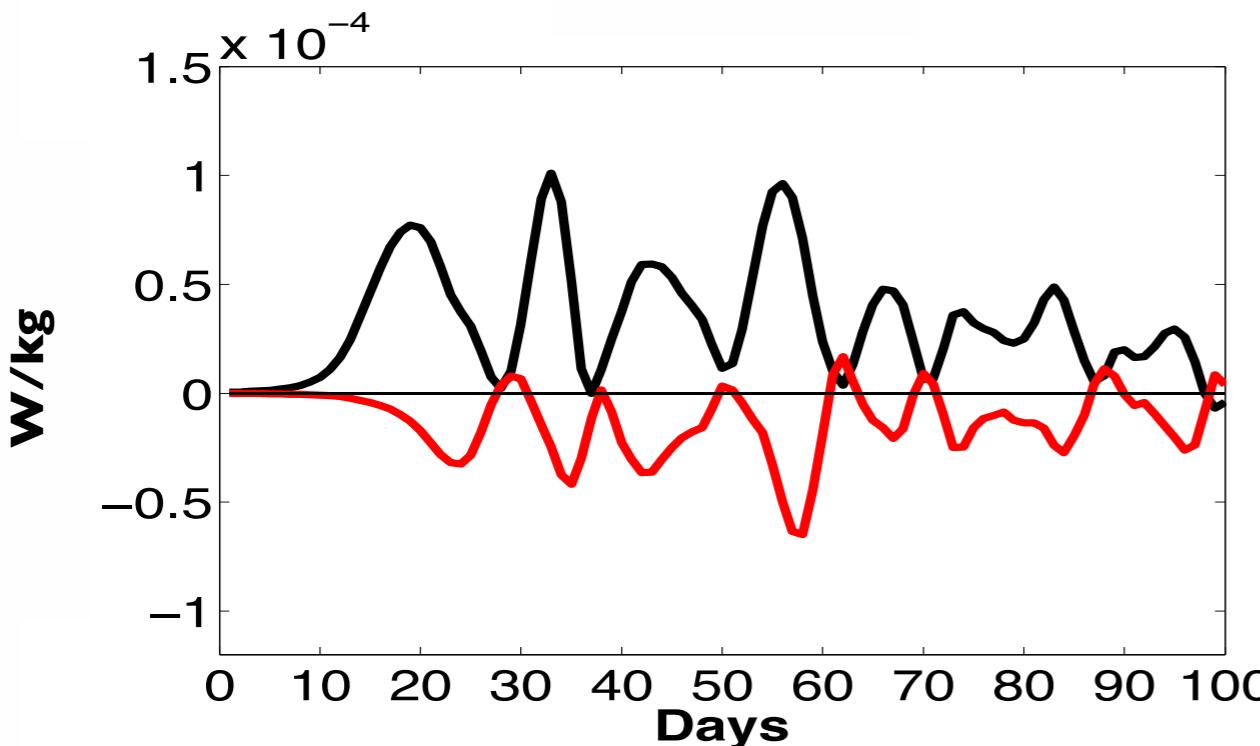
- Initialize a zonal wavenumber 6 perturbation in the zonally averaged circulation (fully non-linear model)
- Experiments run with full model and the no eddy-eddy model

Lifecycle experiments

Baroclinic conversion.

Eddy available potential energy
to eddy kinetic energy

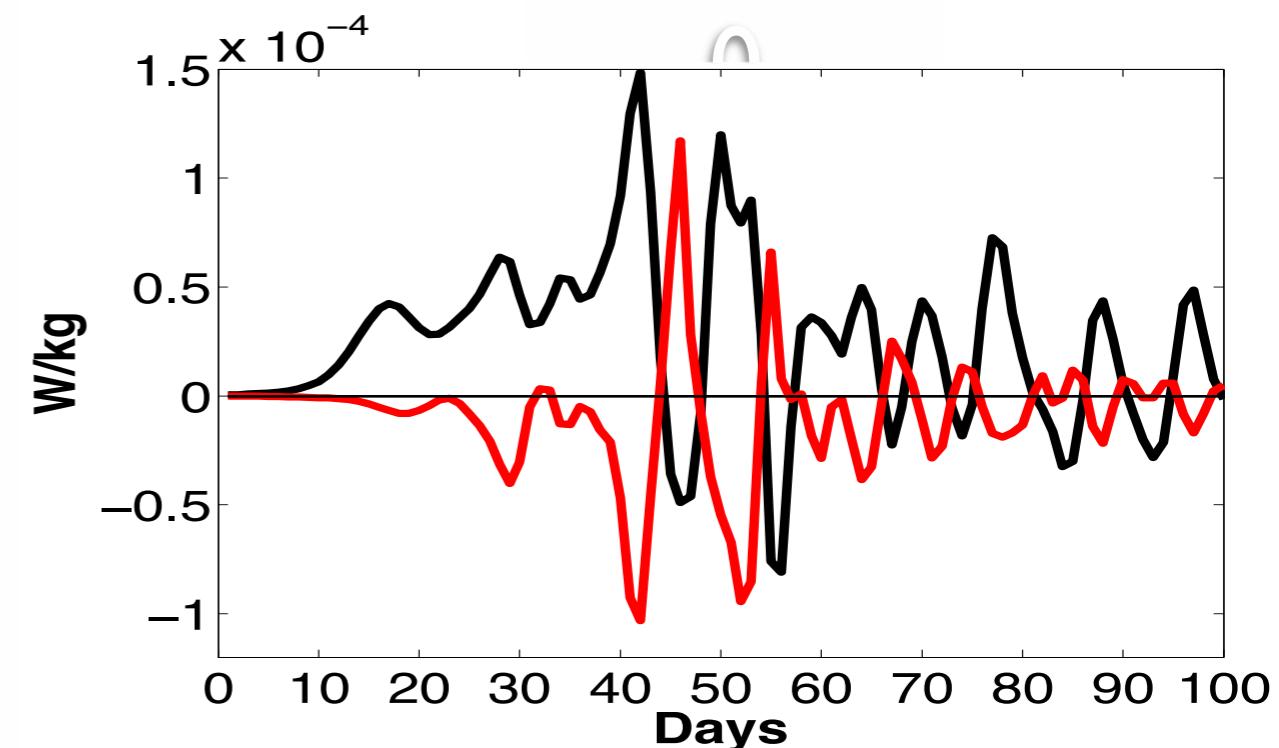
Full



Barotropic conversion.

Zonal kinetic energy to
eddy kinetic energy

QL



- Barotropic conversion large and positive at the end of some lifecycles
- Baroclinic decay significant

EP flux and QG potential vorticity flux

Arrows:
EP flux

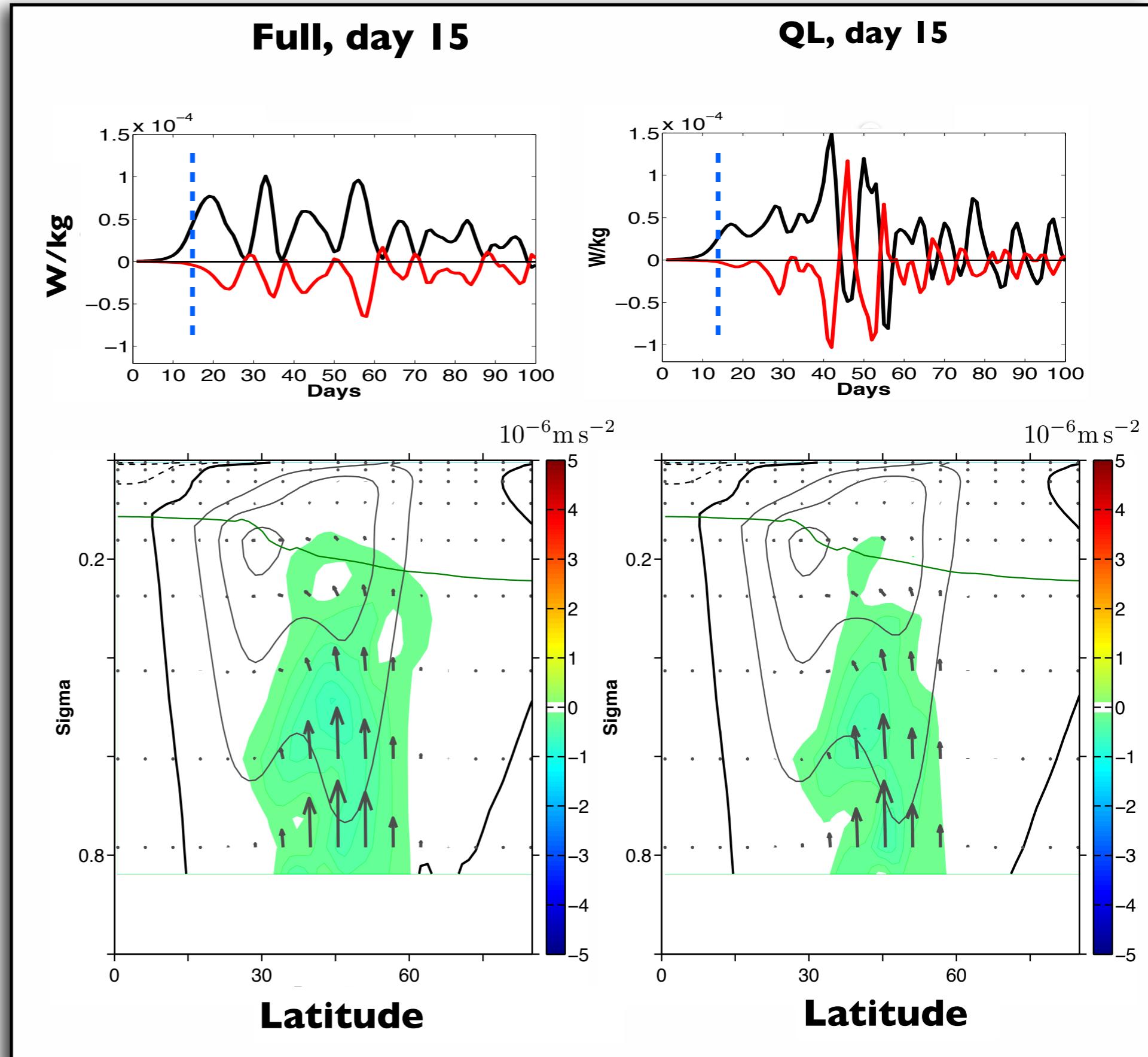
$$\mathbf{F} = R \cos \phi \begin{pmatrix} -\overline{u'v'} \\ f \overline{v'\theta'}/\frac{\partial \bar{\theta}}{\partial p} \end{pmatrix}$$

Colors:
EP flux
divergence,
or QG PV flux,
 $\nabla \cdot \mathbf{F} = R \cos \phi \overline{v'q'}$

Wave activity
equation:

$$\frac{\partial A}{\partial t} + \nabla \cdot \mathbf{F} = \mathcal{D}$$

(Nakamura and Zhu, 2010)



EP flux and QG potential vorticity flux

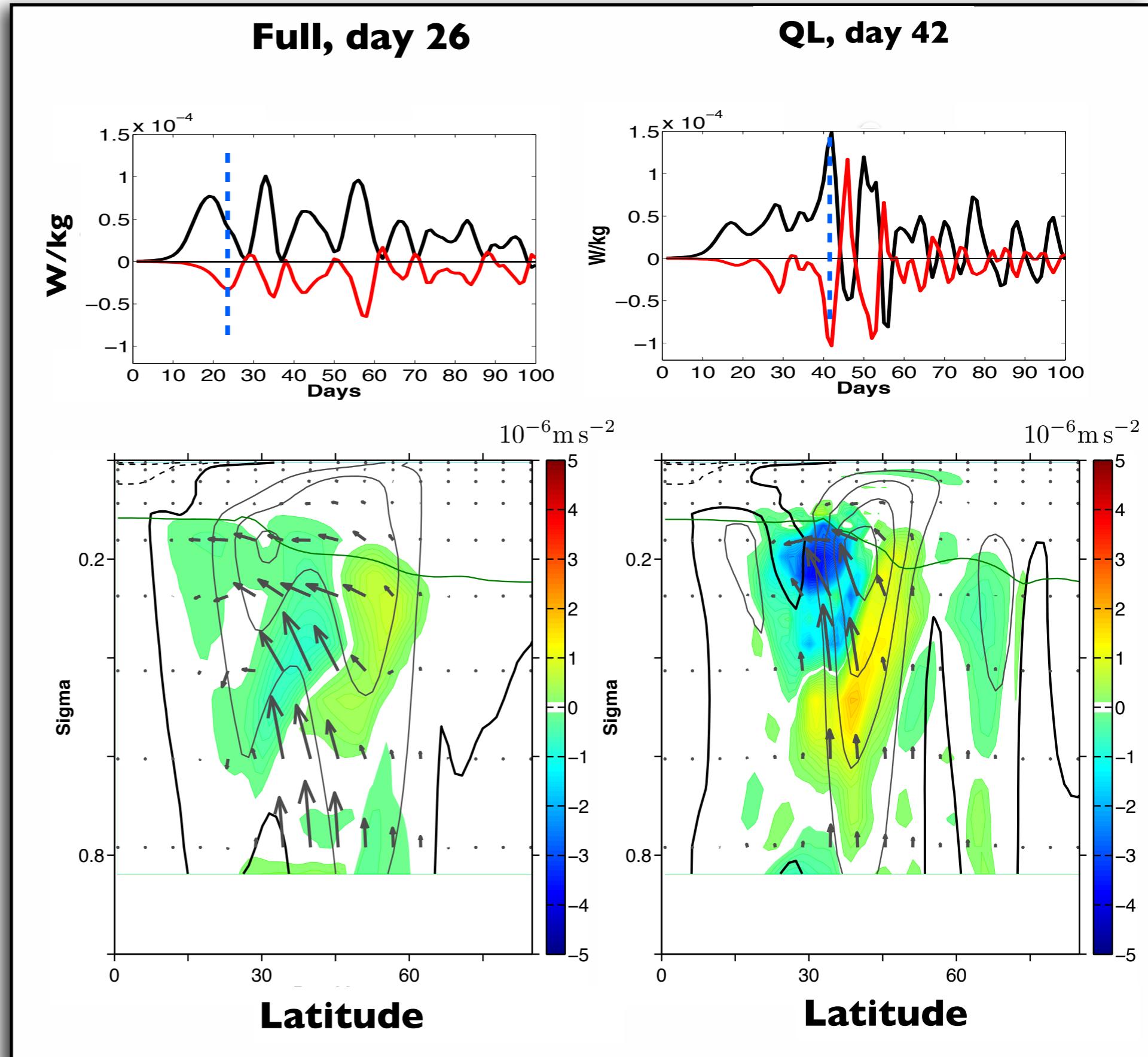
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EP flux

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EP flux and QG potential vorticity flux

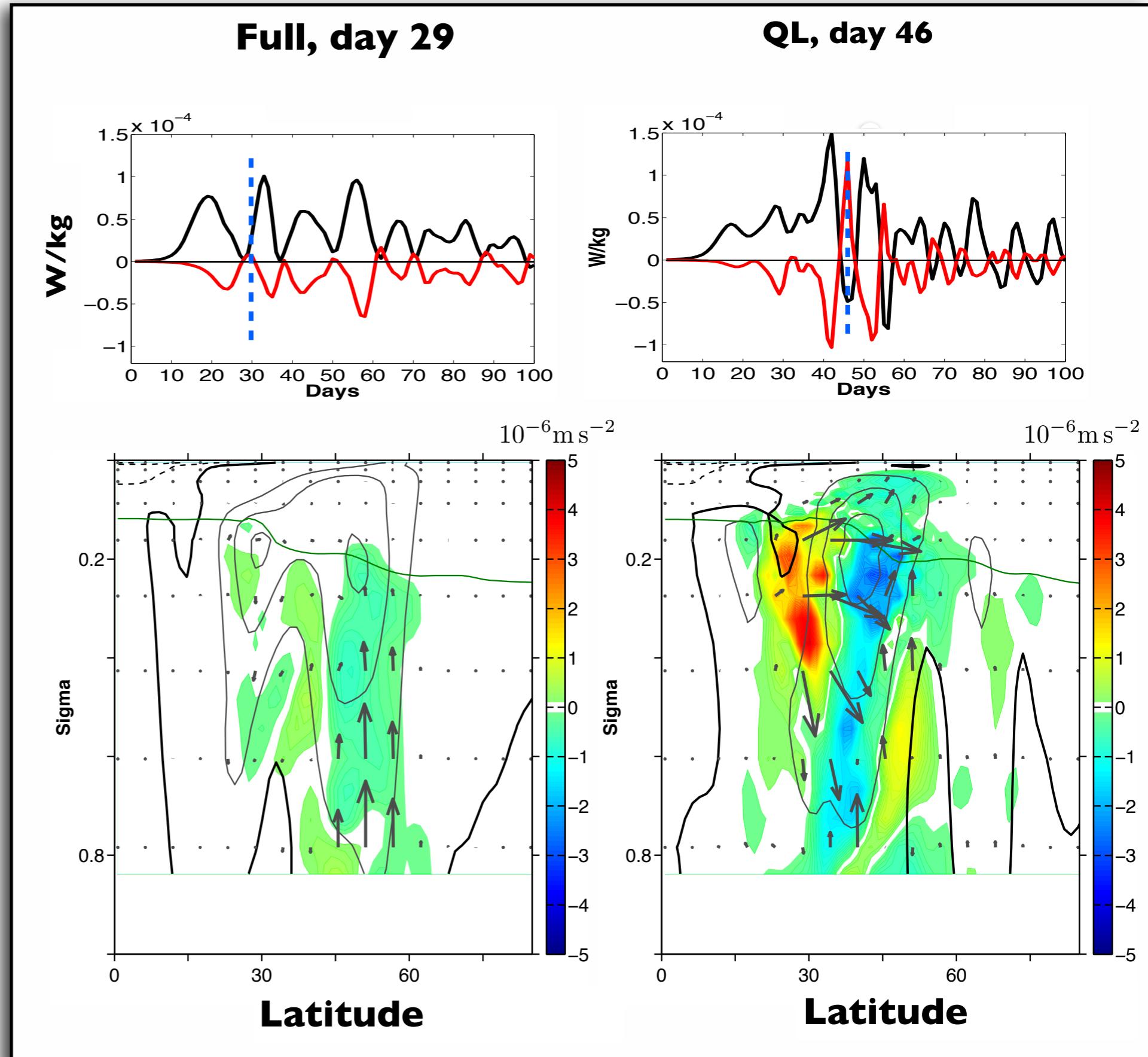
Arrows:
EP flux

$$\mathbf{F} = R \cos \phi \begin{pmatrix} -\overline{u'v'} \\ f \overline{v'\theta'}/\partial p \end{pmatrix}$$

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Wave activity
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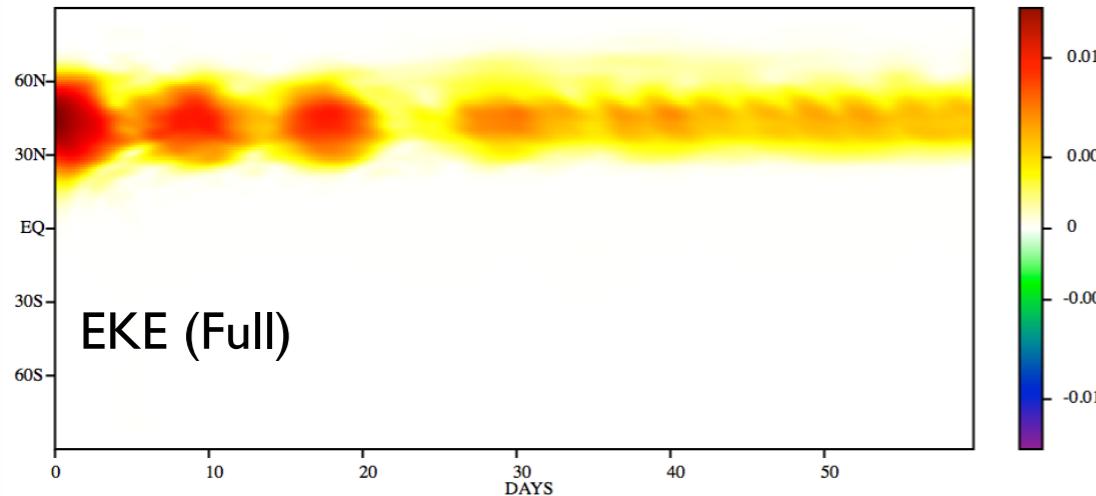


QL Lifecycle experiments

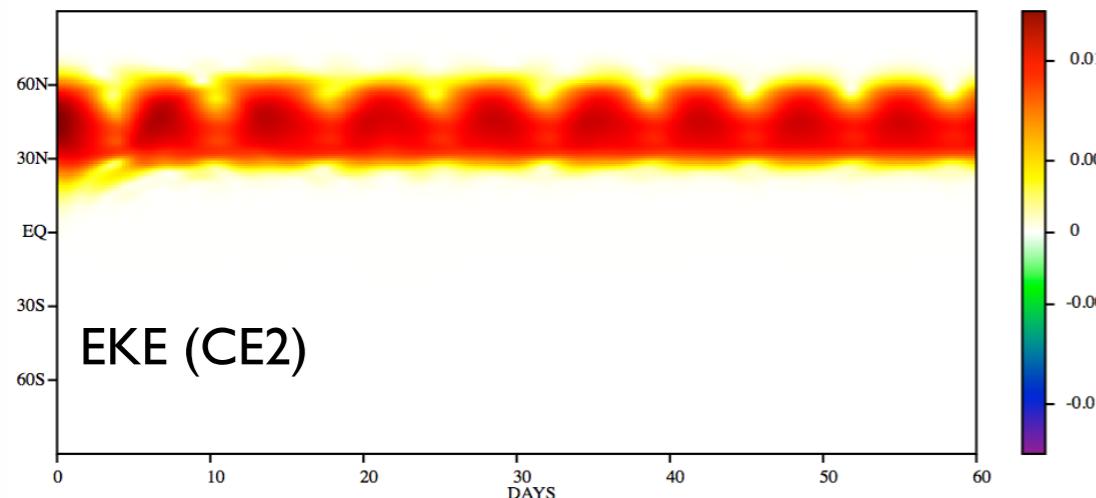
- Absence of eddy-eddy interaction prevents wave activity dissipation
- Wave activity reemitted from critical layer
- “Reflected” wave packet can decay baroclinically

Wave breaking in a barotropic model

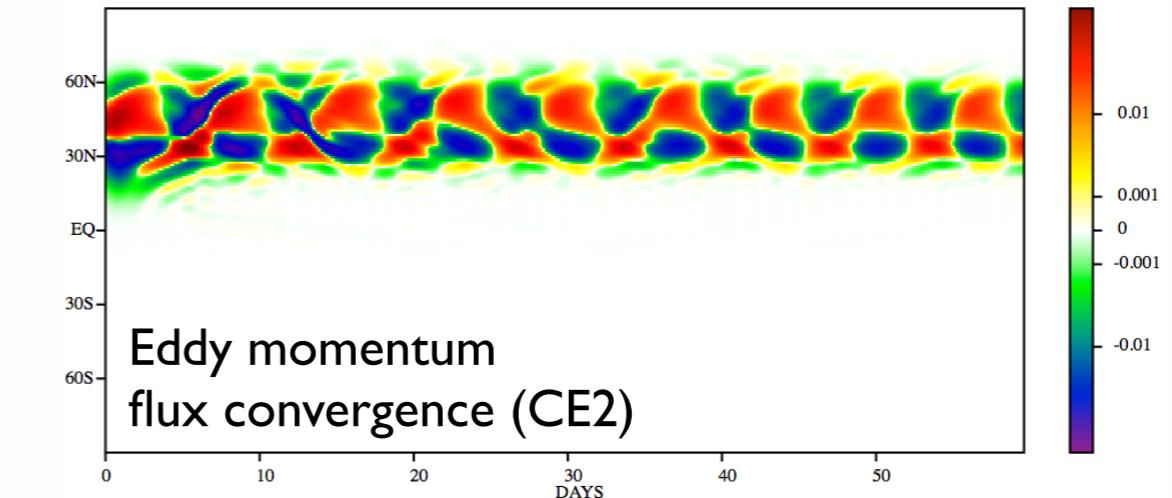
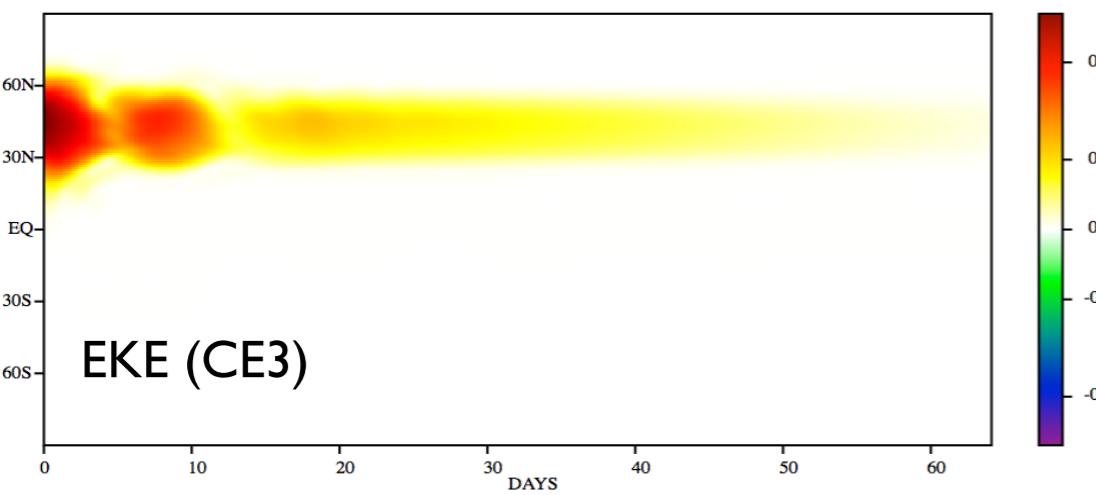
Initial wave at mid-latitudes. The only dissipation is hyperviscosity



CE2 ~ quasi-linear approx



CE3: third order closure. Negative eigenvalues of the second order cumulant are projected out for realizability.



- Rossby wave “resonance” in CE2
- More realistic Rossby wave absorption in CE3

Eddy momentum flux convergence

Colors:

Eddy momentum
flux convergence

Contours:

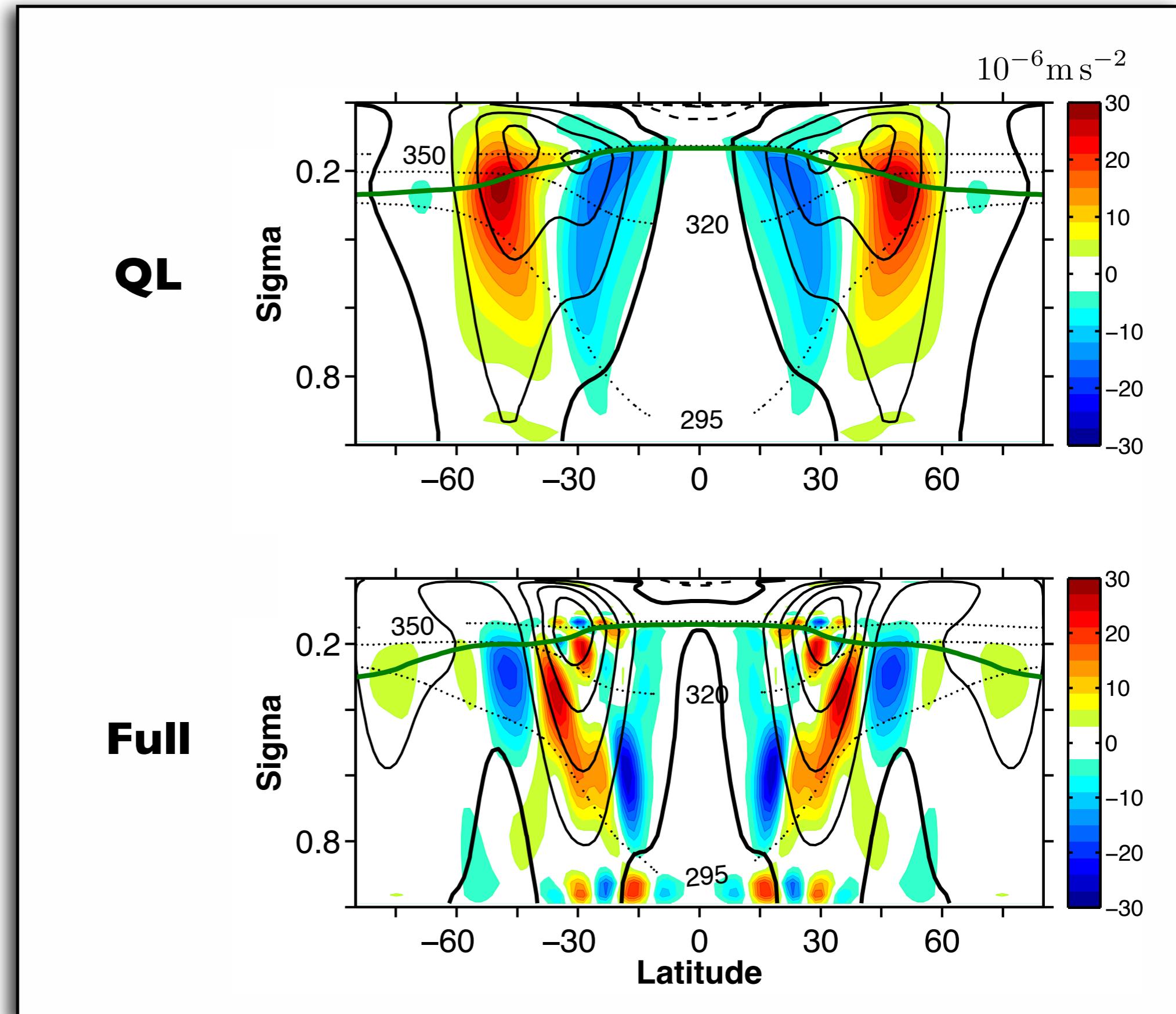
zonal mean flow
(m/s)

Dotted lines:

potential
temperature (K)

Green line:

tropopause



Restoring barotropic triads

Colors:

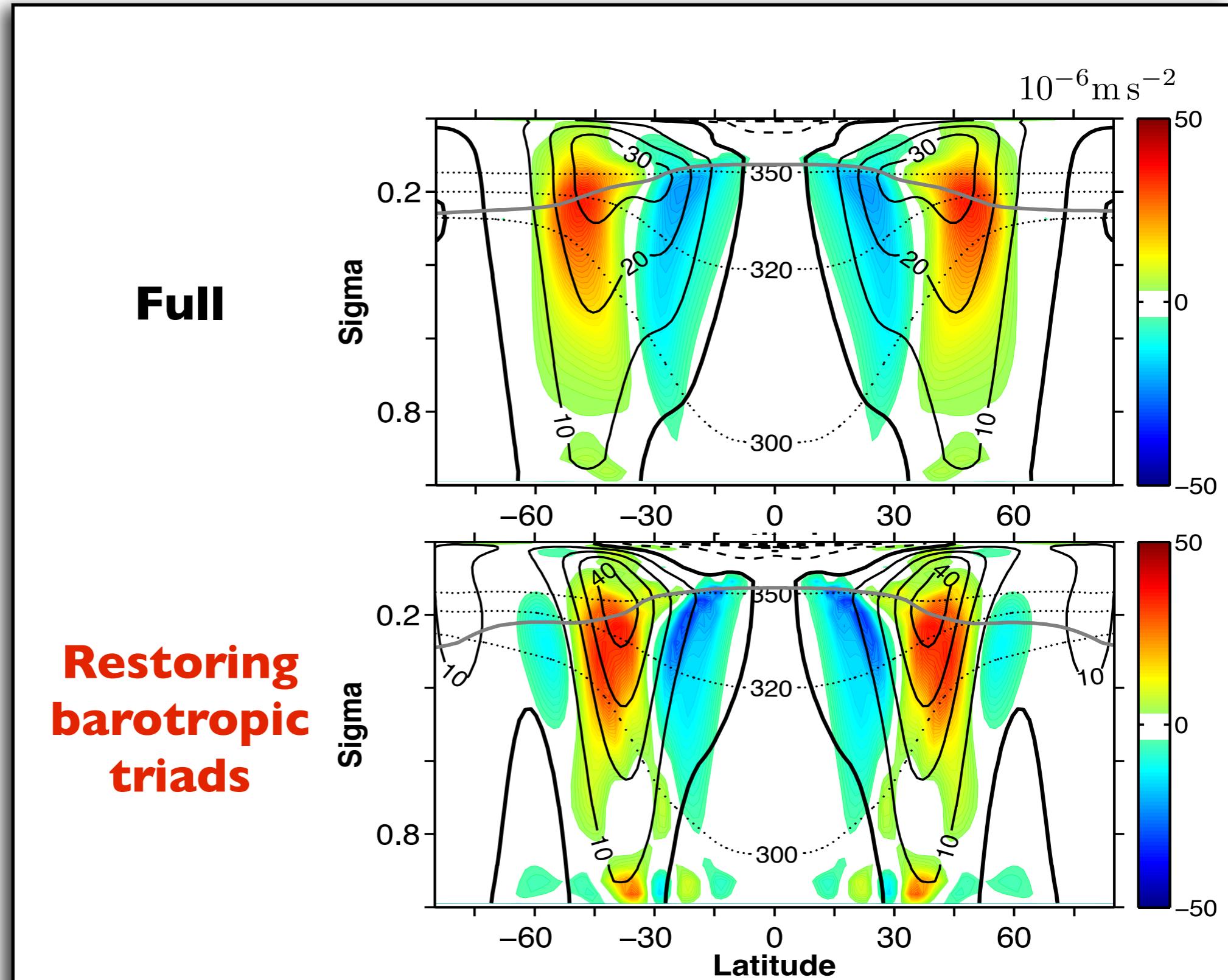
eddy momentum
flux convergence

Contours:

zonal mean flow
(m/s)

Dotted lines:

potential
temperature (K)

Grey line:
tropopause**Full****Restoring
barotropic
triads**

Viscous critical layer in the QL approximation

Held and suarez GCM, QL

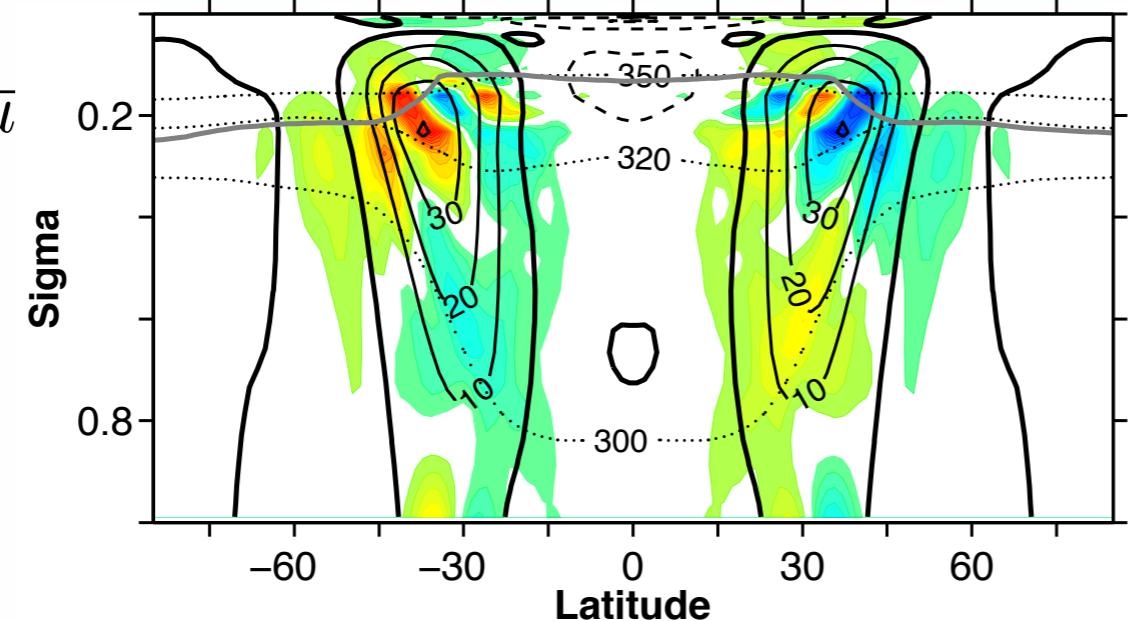
Order 4 hyperviscosity

Contours:

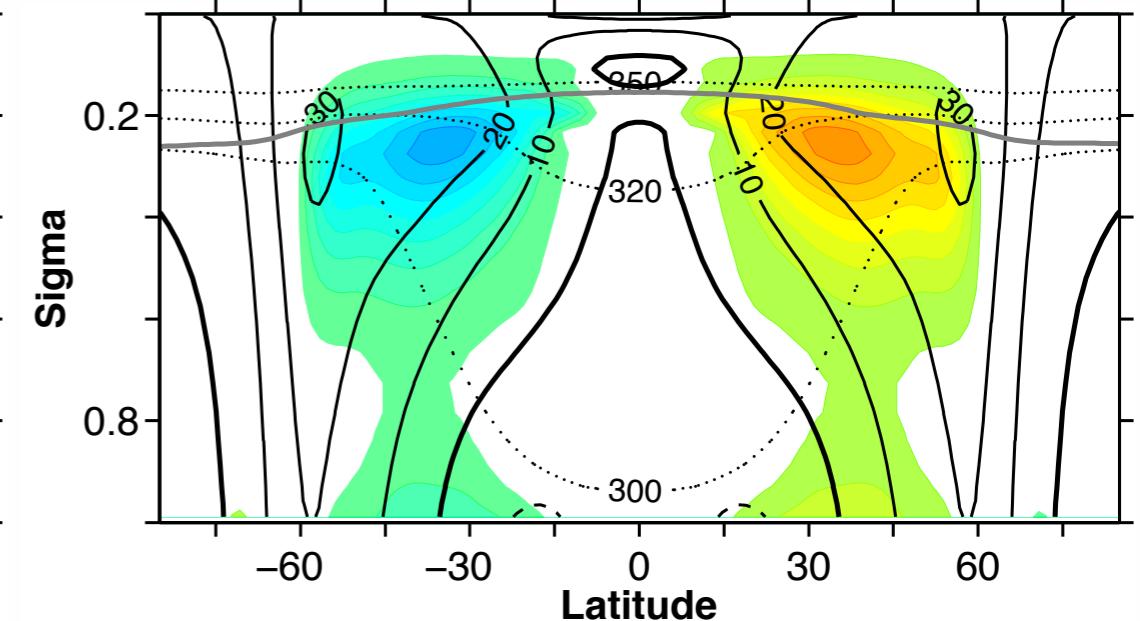
Zonal wind \bar{u}
in m/s

Colors:

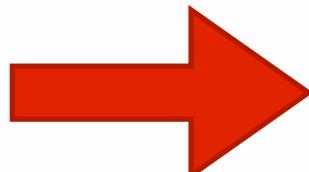
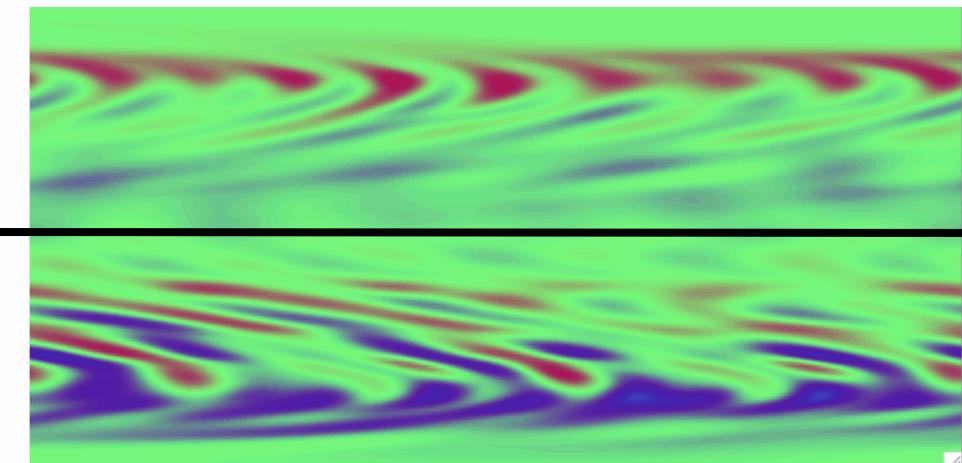
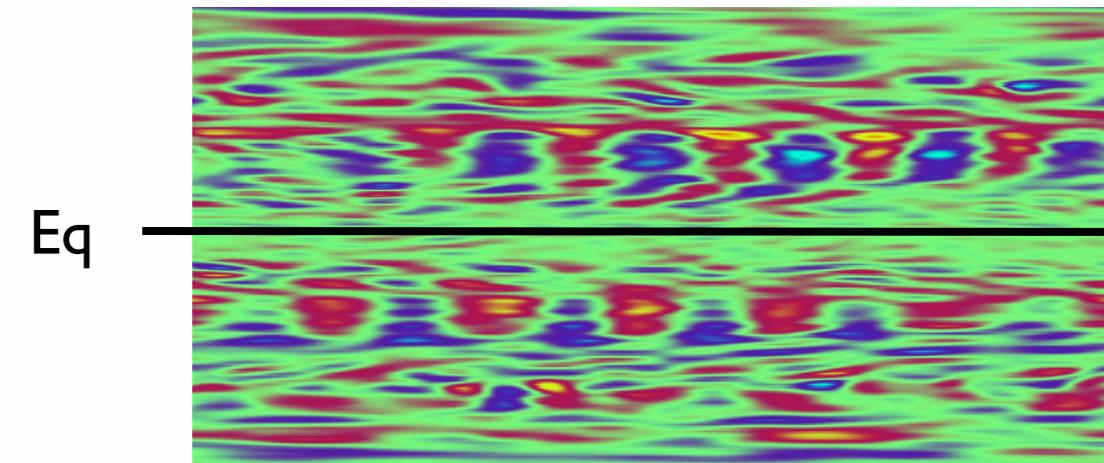
momentum
flux
 $\bar{u}'v' \cos \phi$



Diffusion



Relative
vorticity,
274 hpa



Upper-tropospheric wave absorption and EMF
enhancement captured by a viscous critical layer

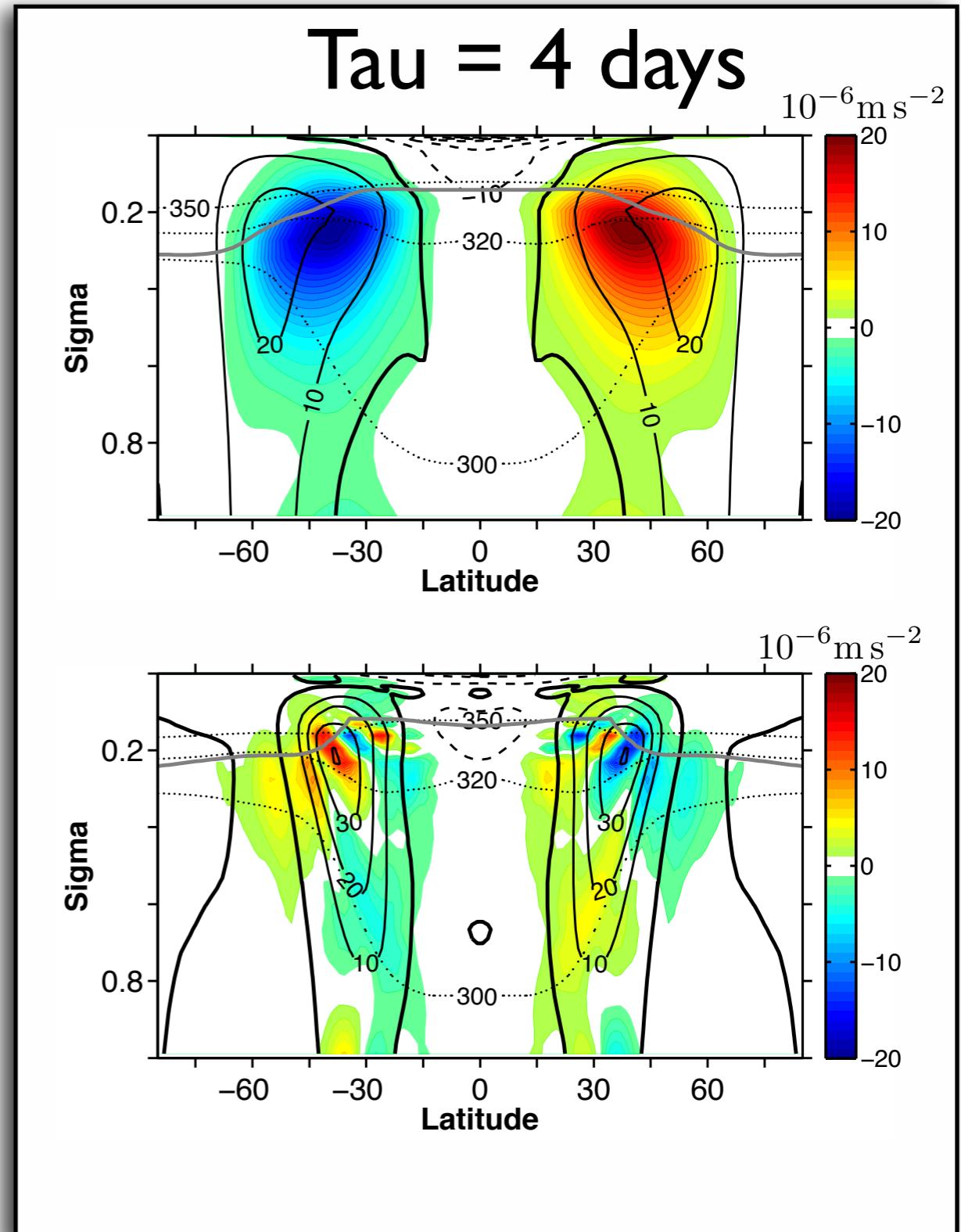
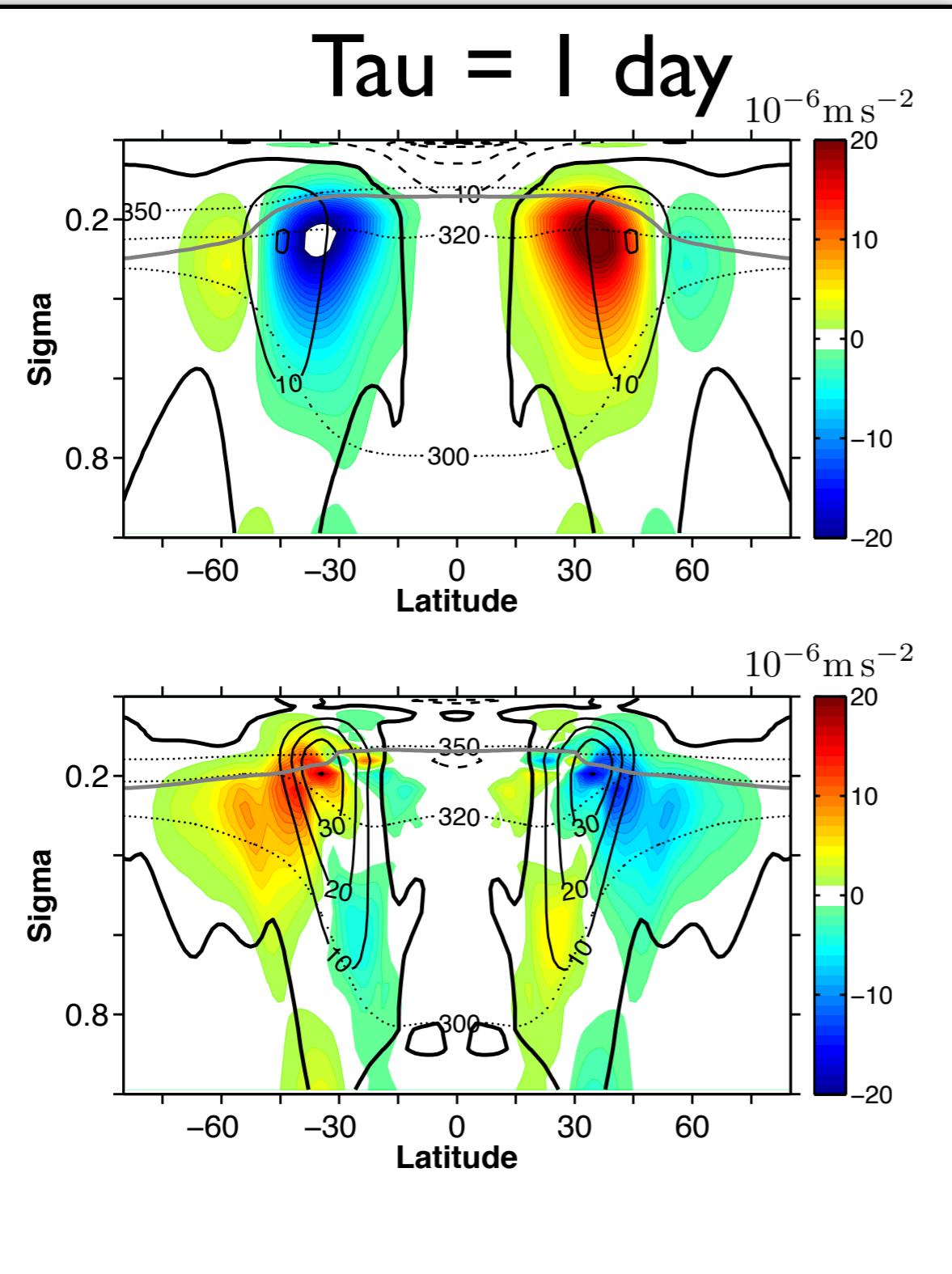
Varying surface friction

- Surface friction essential in the general circulation for momentum and energy budget.
- Complex effect on midlatitudes eddies amplitude, jet streams strength and location, energy conversions (James, 1987; Robinson 1997; Chen et al., 2007; Liu and Schneider, 2014, etc...)
- How does QL captures general circulation when varying surface friction? **We change Tau, the time scale of Rayleigh drag in the planetary boundary layer.**

Varying surface friction: momentum fluxes

Colors: momentum flux $\overline{u'v'} \cos \phi$ in m^2/s^2

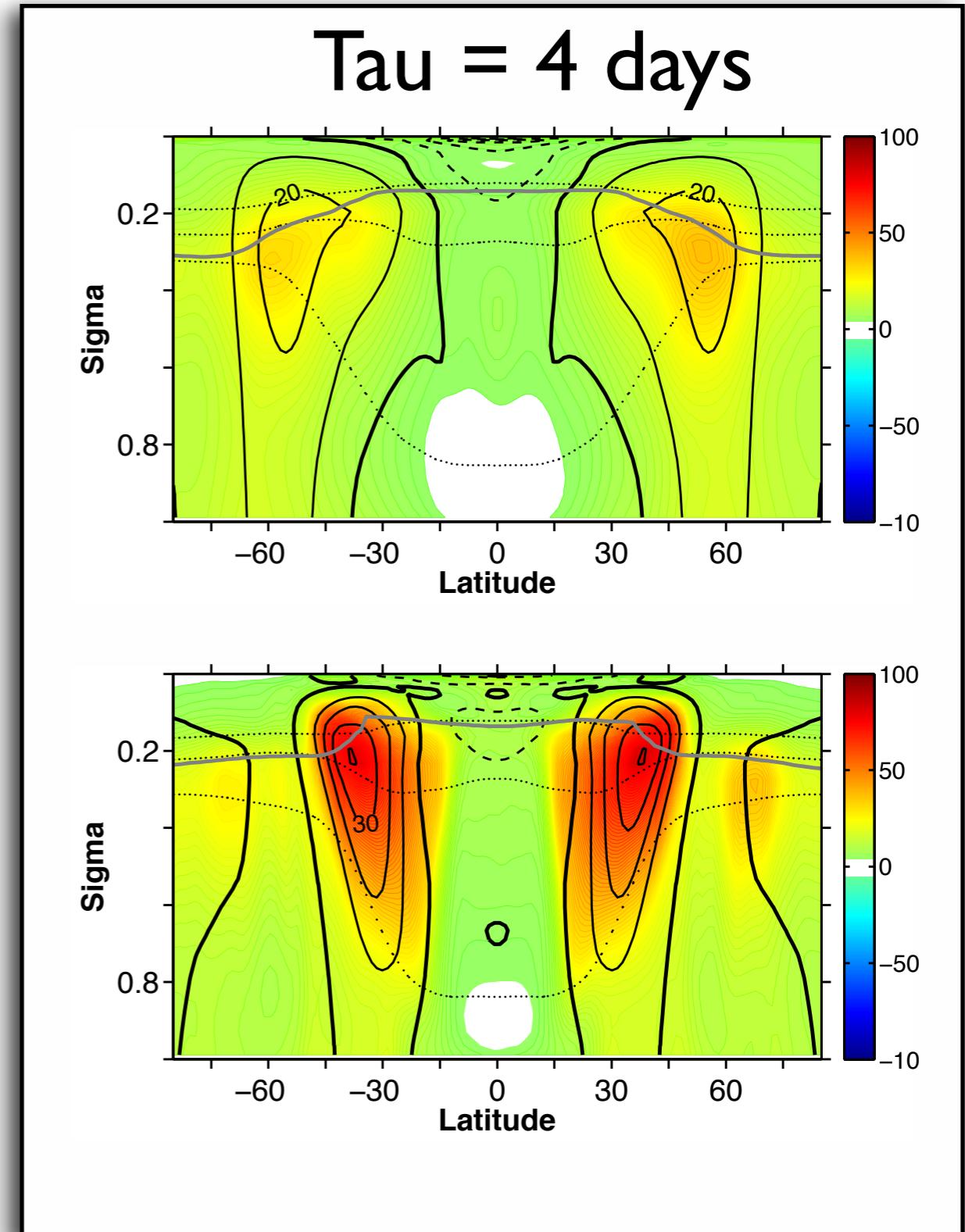
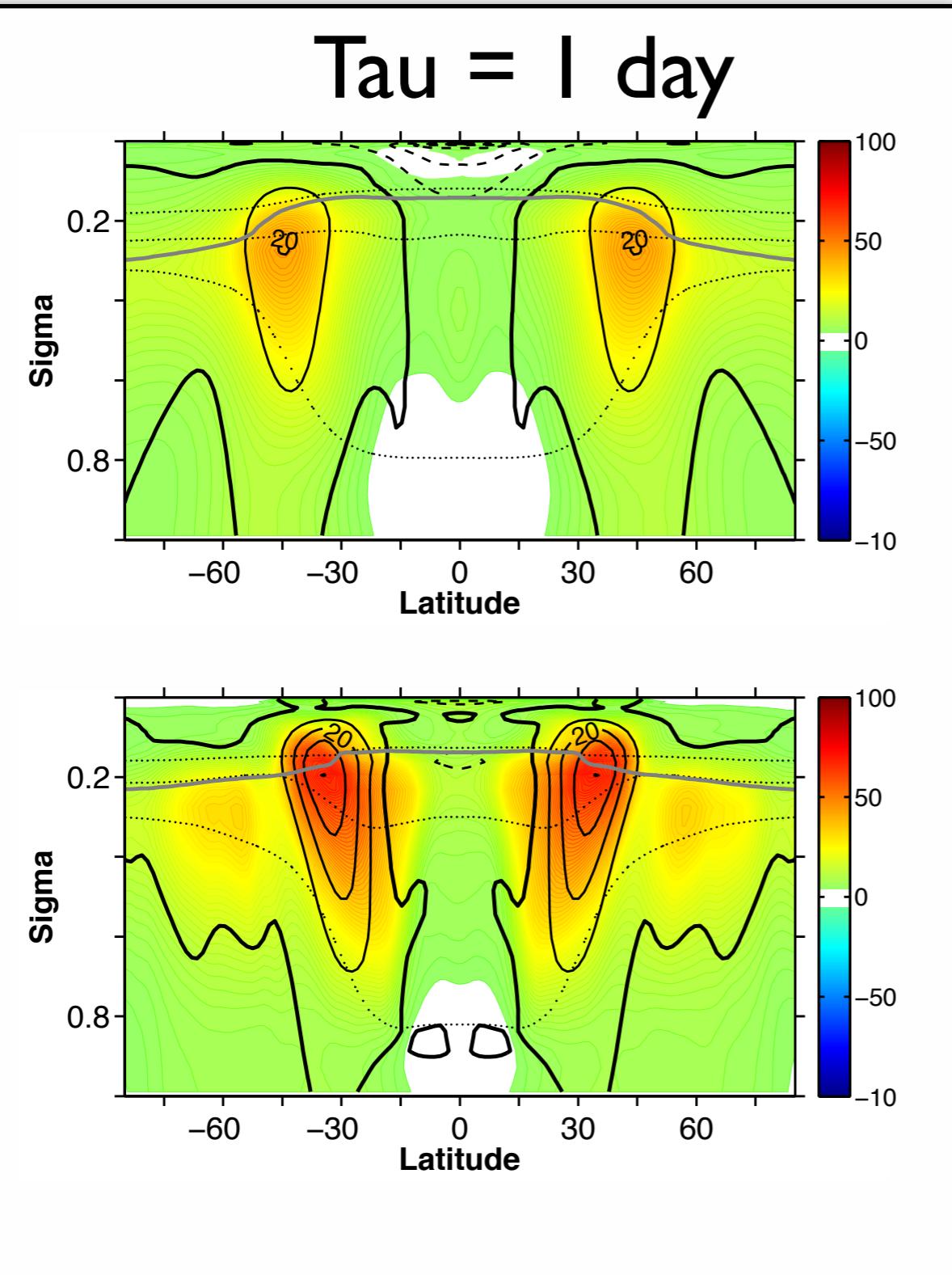
Contours: Zonal wind \bar{u} in m/s



Varying surface friction: zonal wind and EKE

Colors: Eddy kinetic energy $(\overline{u'^2} + \overline{v'^2})^{\frac{1}{2}}$ in m/s

Contours: Zonal wind \overline{u} in m/s



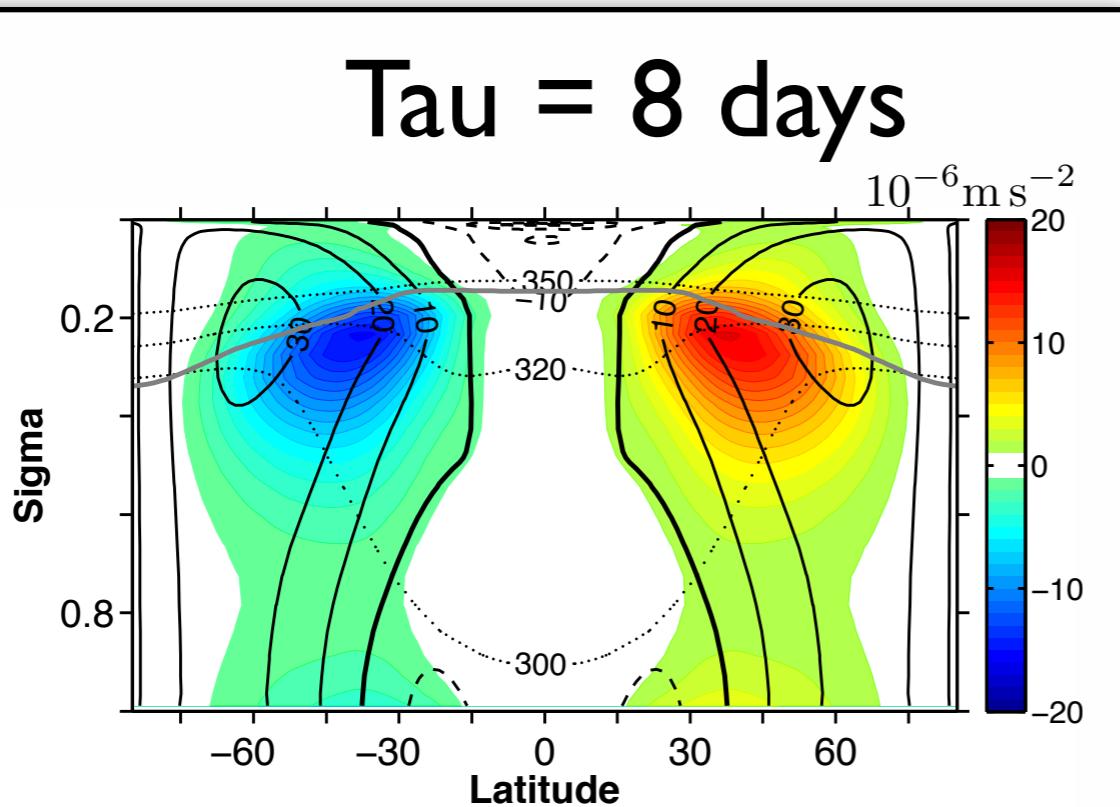
Varying surface friction: momentum fluxes

colors: momentum flux $\overline{u'v'} \cos \phi$ in m^2/s^2

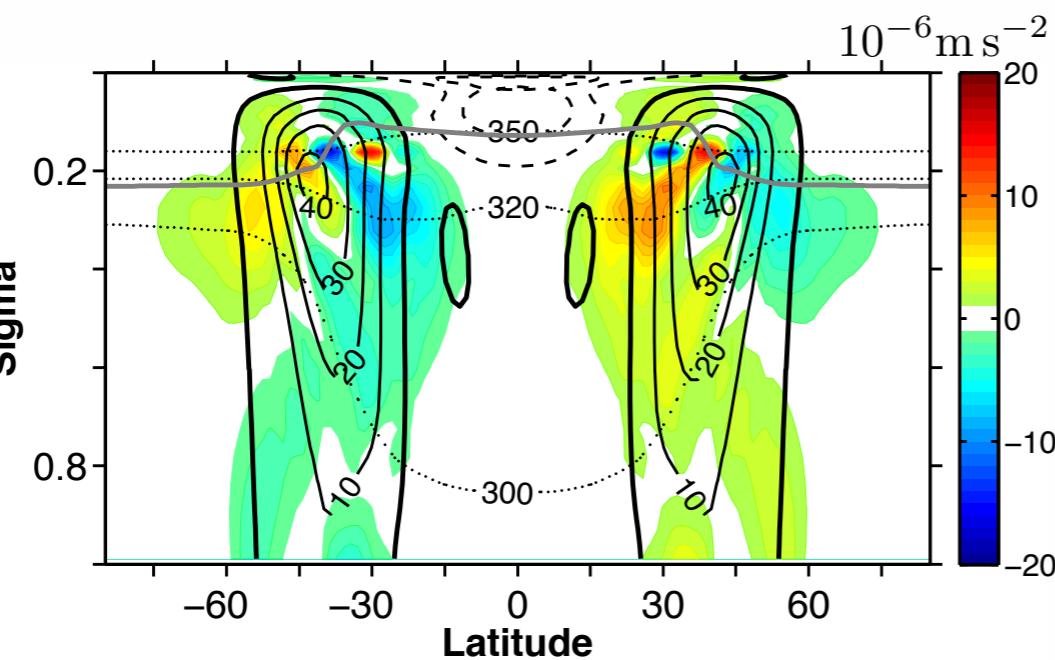
Contours: Zonal wind \overline{u} in m/s

full

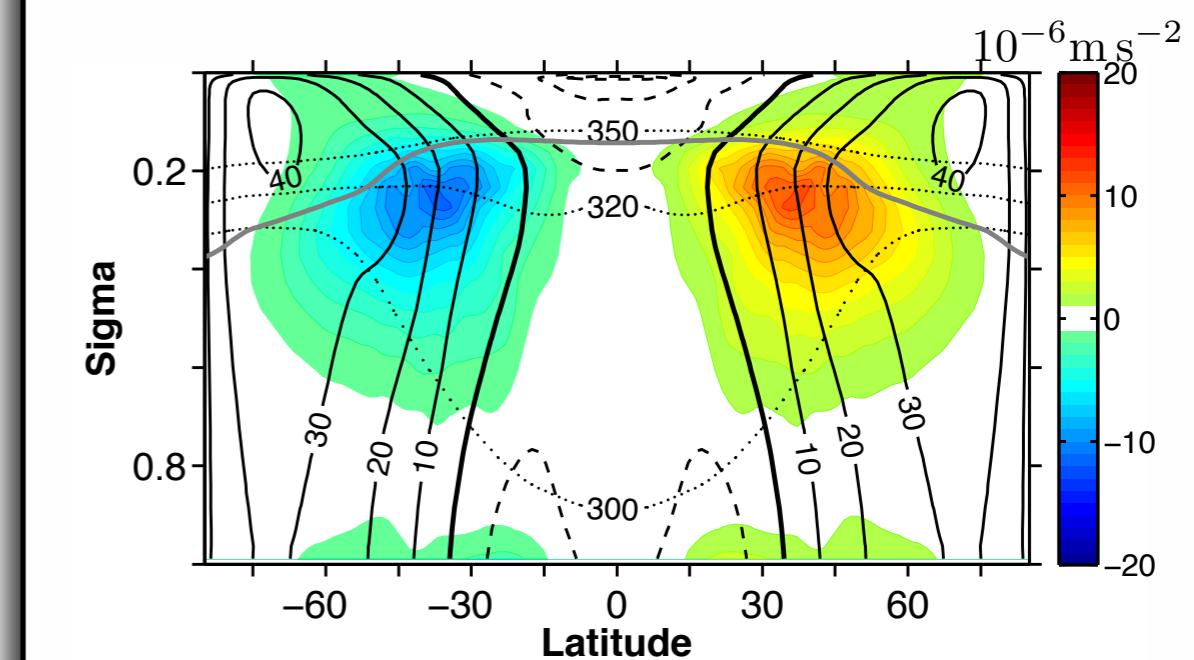
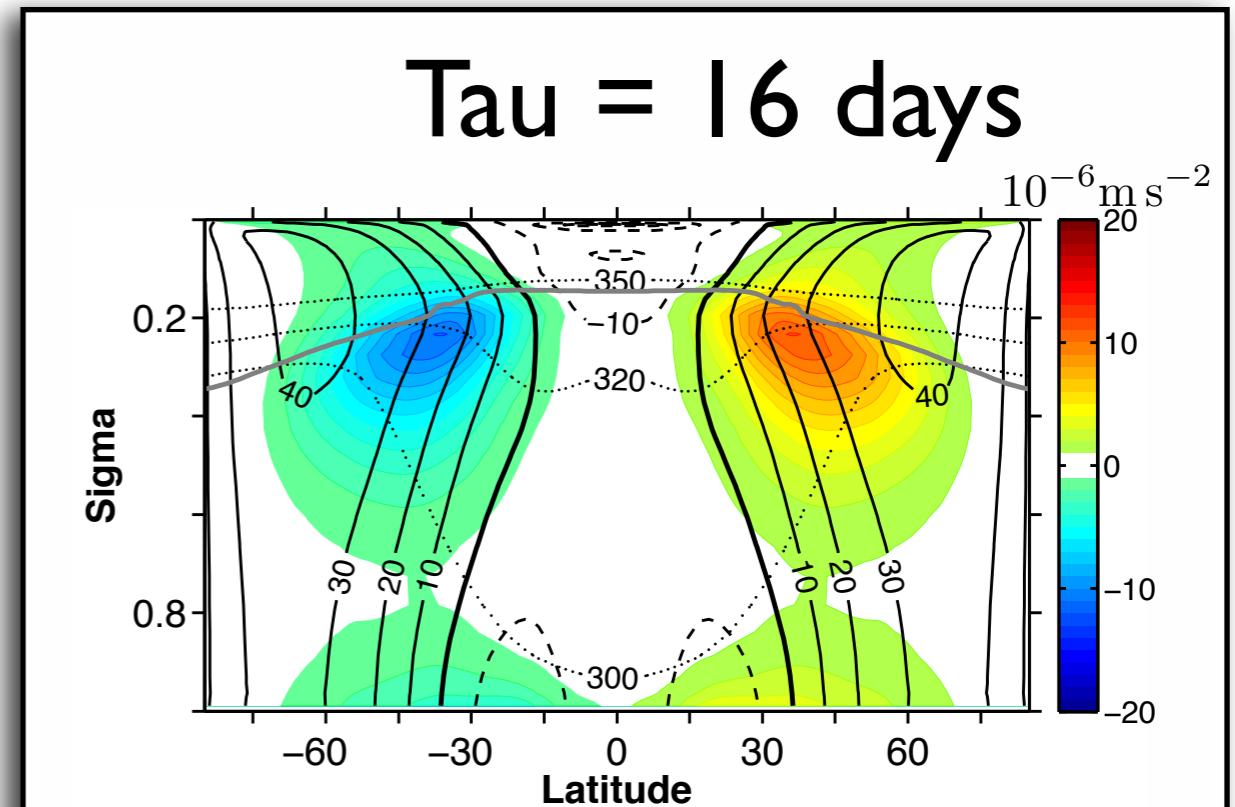
Tau = 8 days



Q_L



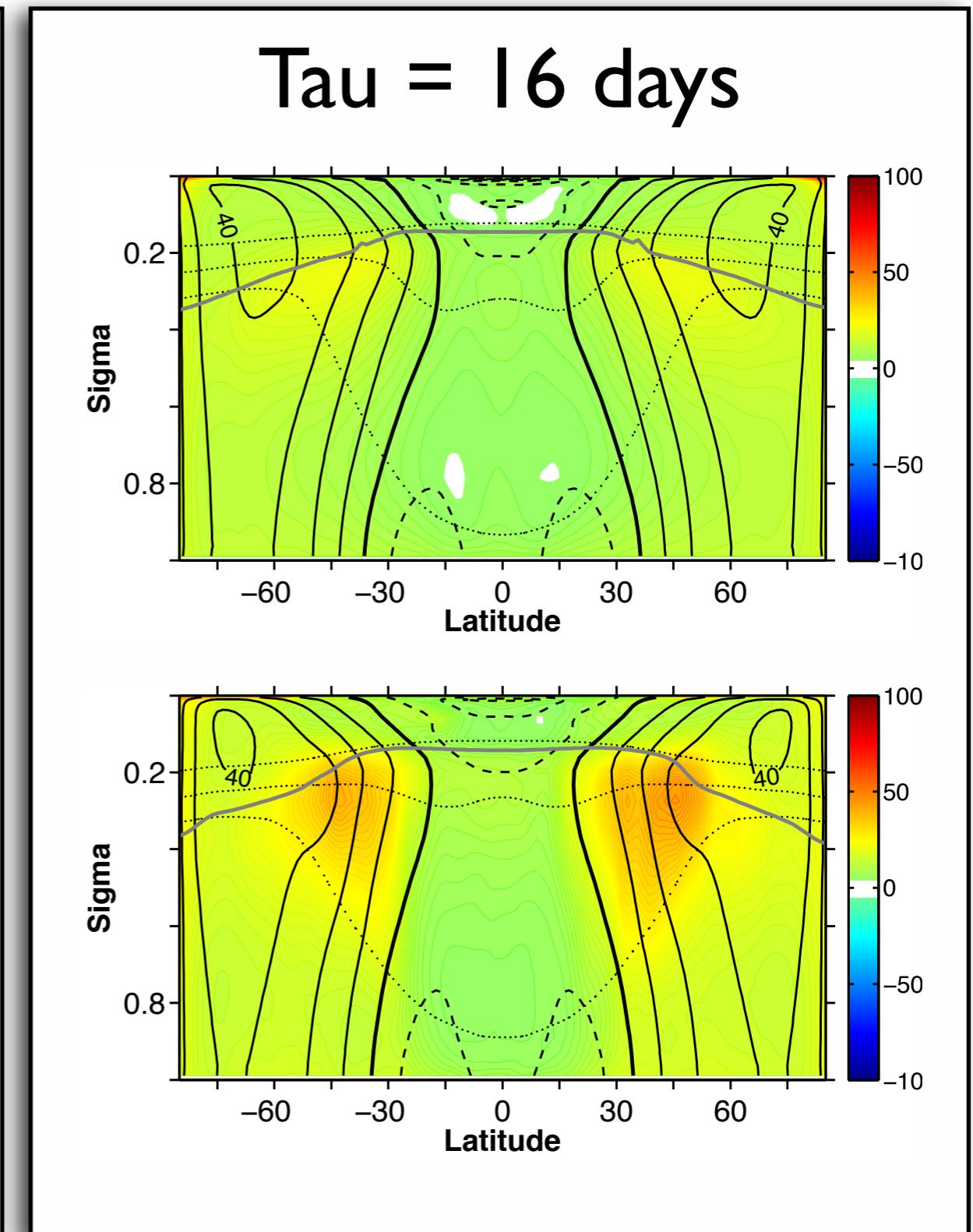
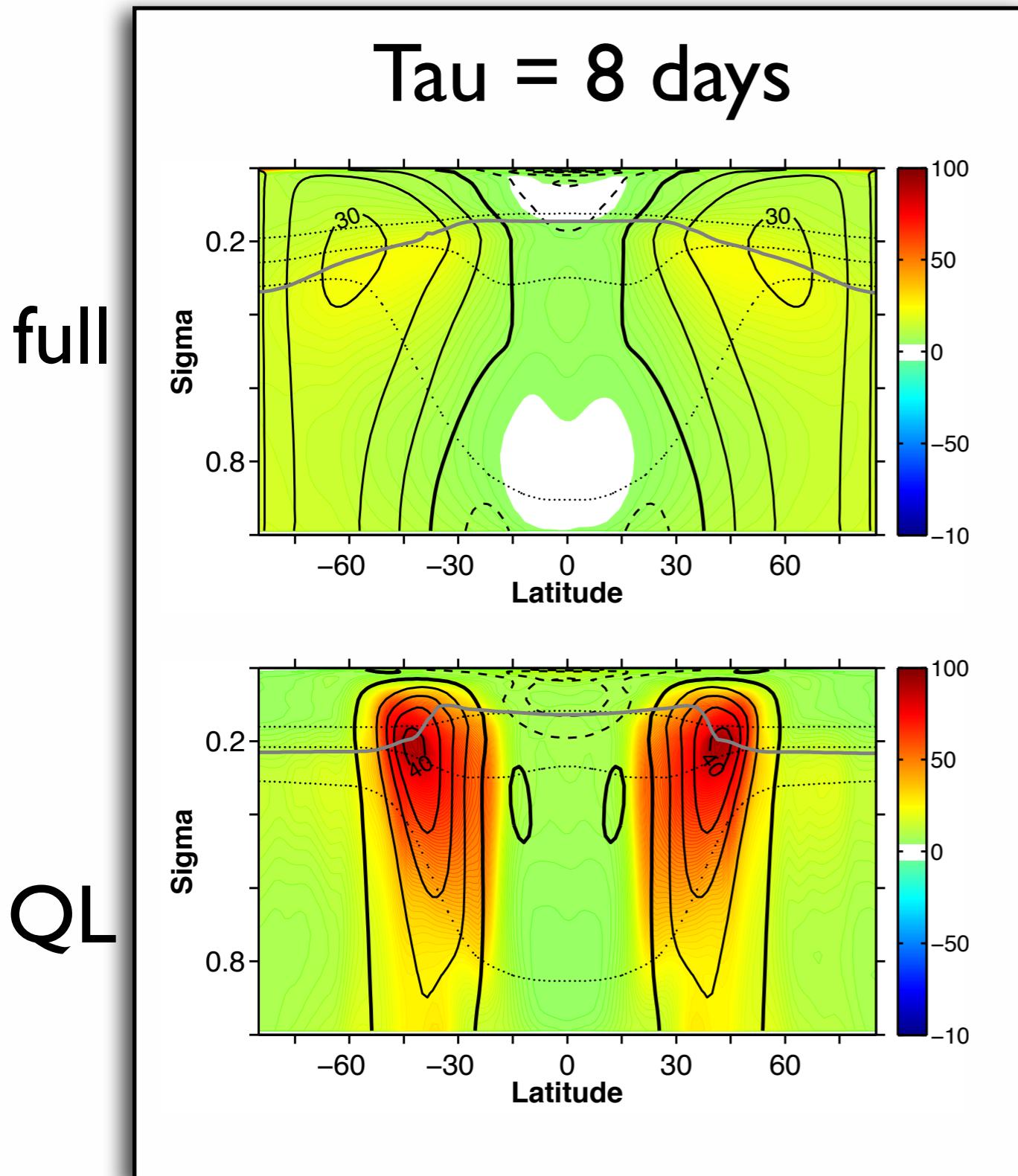
Tau = 16 days



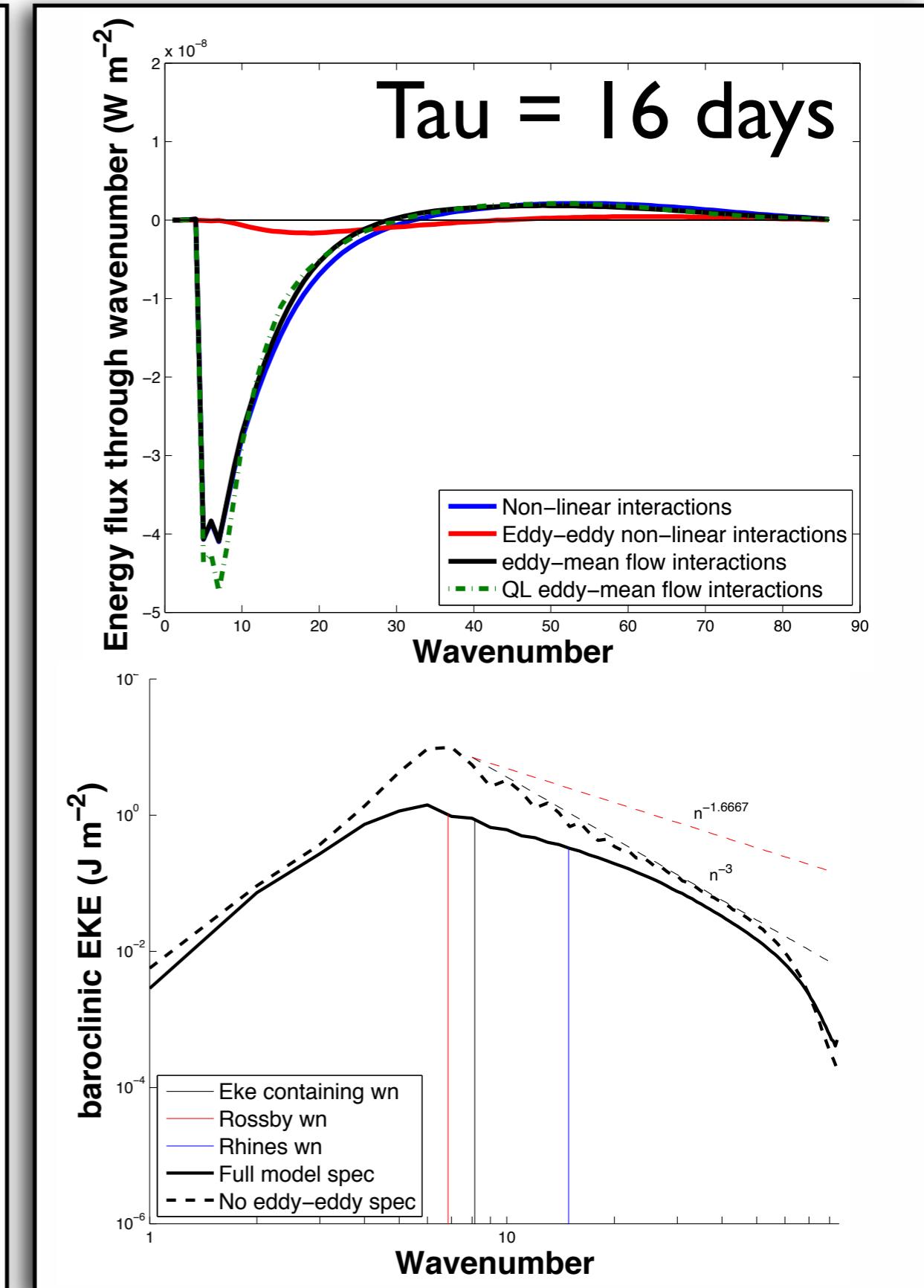
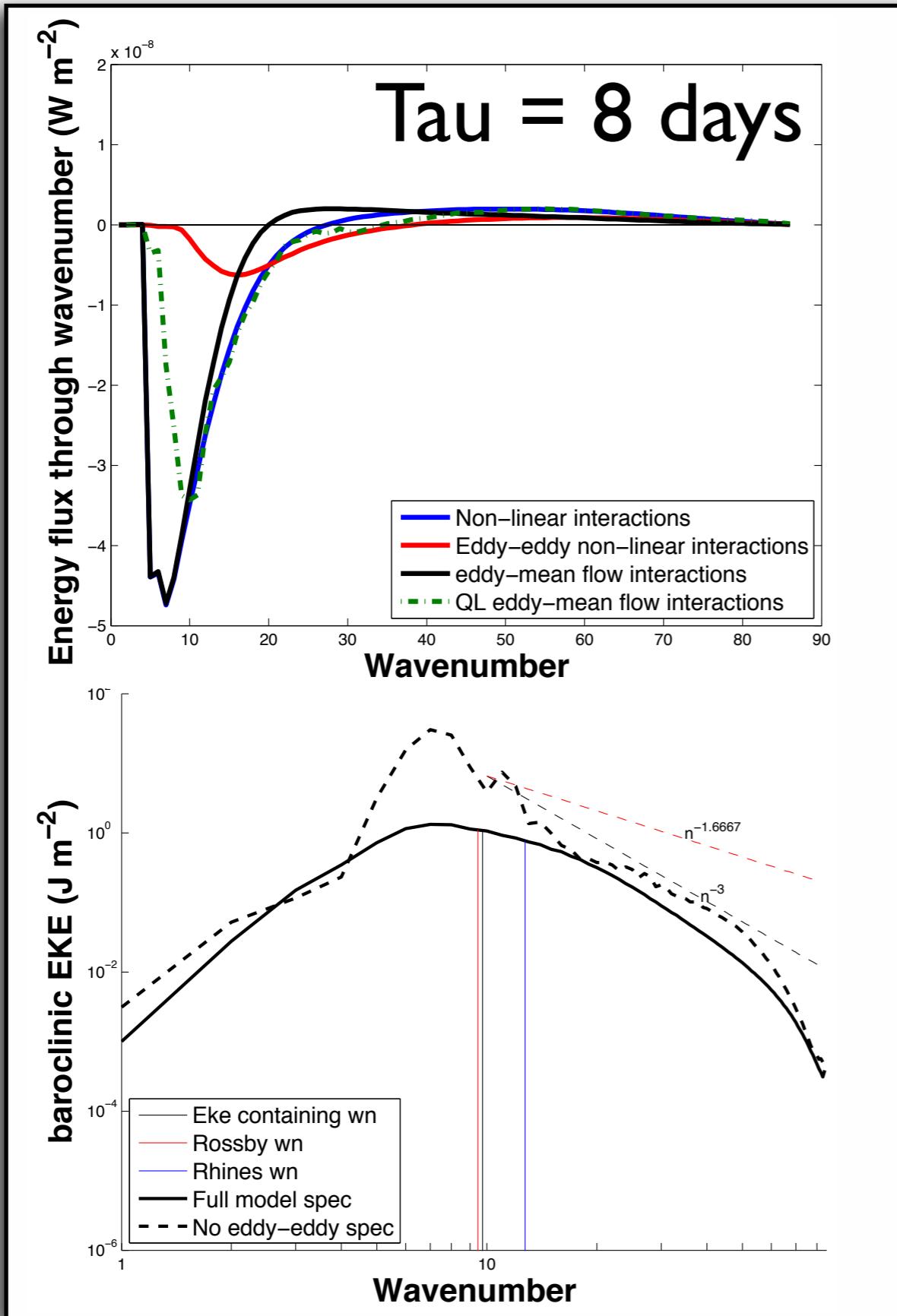
Varying surface friction: zonal wind and EKE

Colors: Eddy kinetic energy $(\overline{u'^2} + \overline{v'^2})^{\frac{1}{2}}$ in m/s

Contours: Zonal wind \overline{u} in m/s

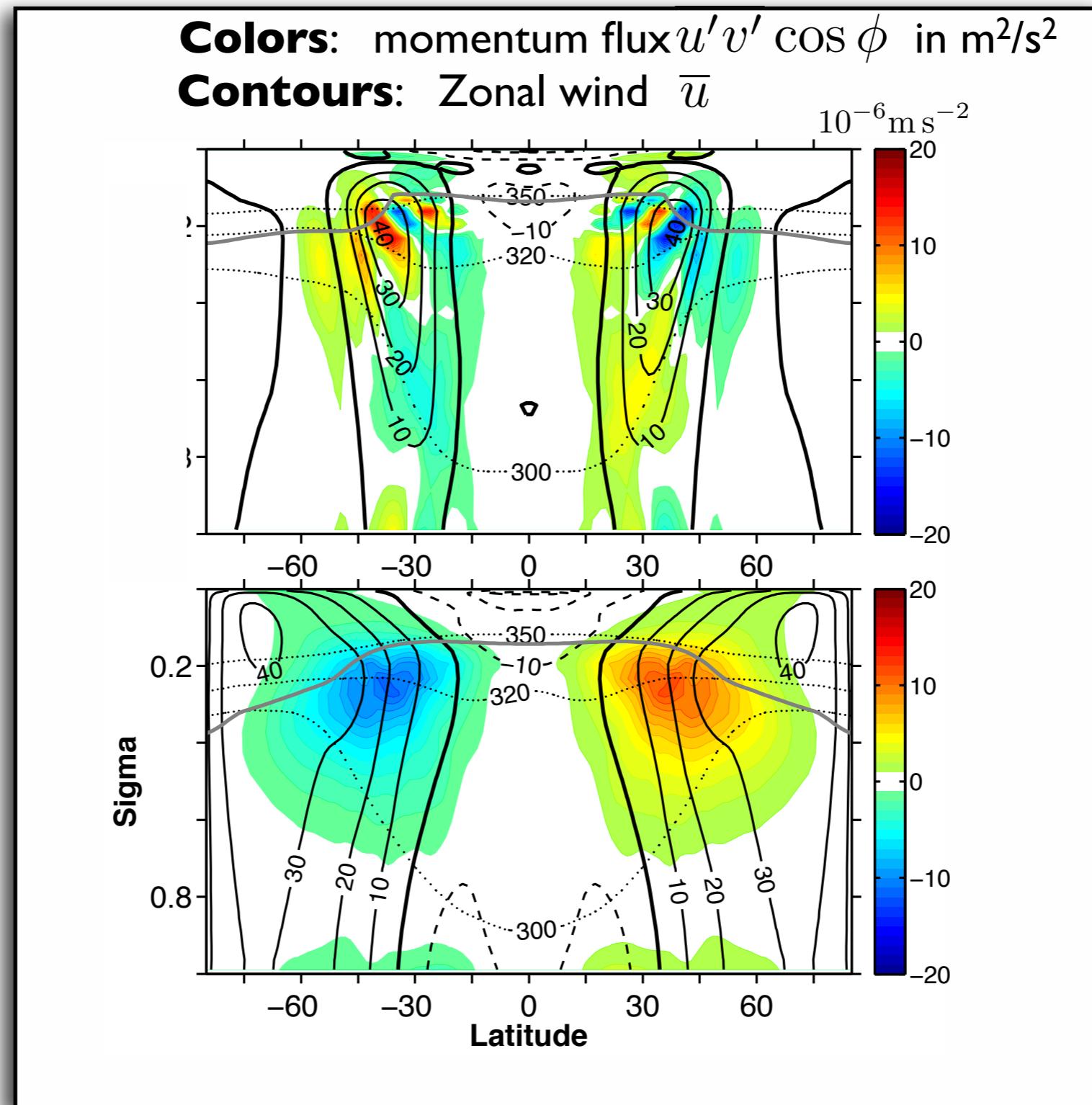


Varying surface friction: spectral budget



Multiple equilibria at “intermediate” frictions

**Tau = 12 days, quasi-linear approximation,
two different initial conditions**



Varying surface friction in QL

- Jet structure and poleward migration of the jet with weaker friction not captured below a threshold ($\sim 10^1$ days)
- Abrupt transition to a configuration where momentum flux and zonal winds (but not EKE) are well captured (linear decay of *Held and Philipp, 1987*). Scale separation between waves and NL critical layer?
- Multiple equilibria possible for some intermediate values.

Conclusions

- Relevant irreversible process for wave activity dissipation lies in eddy-eddy interactions. No length scale separation between eddies and surf zone (SWW theory: even if weak, eddy-eddy interactions matter at leading order for waves outside the critical layer).
- Upper tropospheric EMF enhancement not fully explained by a “more linear” upper-troposphere and “more turbulent” lower troposphere, as proposed by linear theories (Held 2000, 2007)
- For weak enough friction, QL captures mean flows and eddy forcing, but not eddy statistics. Why? we are still far from time scale separation....
- Third order closures (parametrizing eddy fluxes of potential enstrophy and capturing NL wave breaking) seem necessary to develop direct statistical simulation of geophysical flows.