Jets: a roller-coaster ride from Earth to Jupiter and back, or:

A tale of three species, or:

Is homogeneous turbulence theory a dangerous idea?

Michael E. McIntyre,
Dept of Applied Mathematics & Theoretical Physics,
University of Cambridge, UK

- 1: Very brief look at standard jet mechanisms and the case of earthly strong jets, following Rosenbluth and Haurwitz lectures
- 2. Hot off the press: results from a new idealized model of **Jupiter's jets** strange new territory!!
- 3. Even hotter (not to say hasty and preliminary) surprises from the extended Hasegawa-Mima equation including a `damn fool experiment' ...

Our life support system



Our life support system



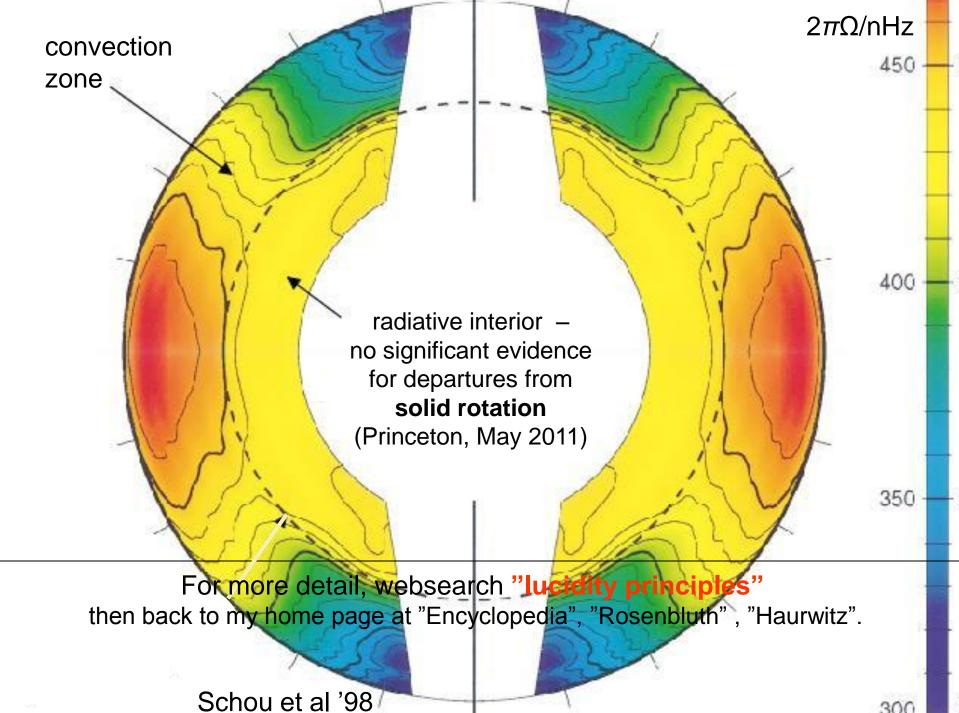
fluid dynamics on a grand scale...

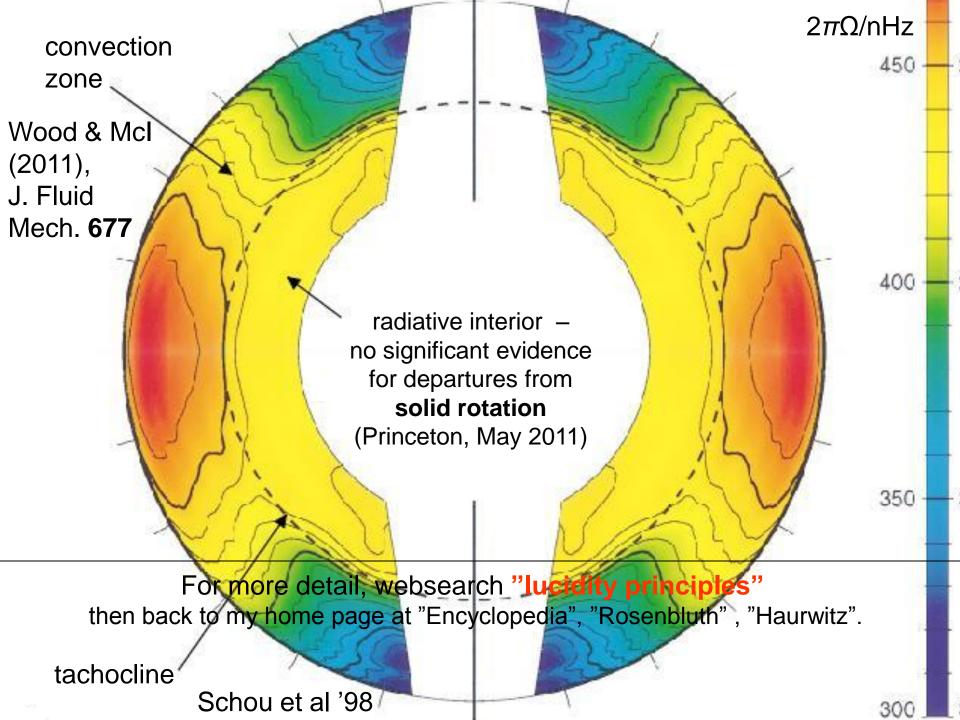
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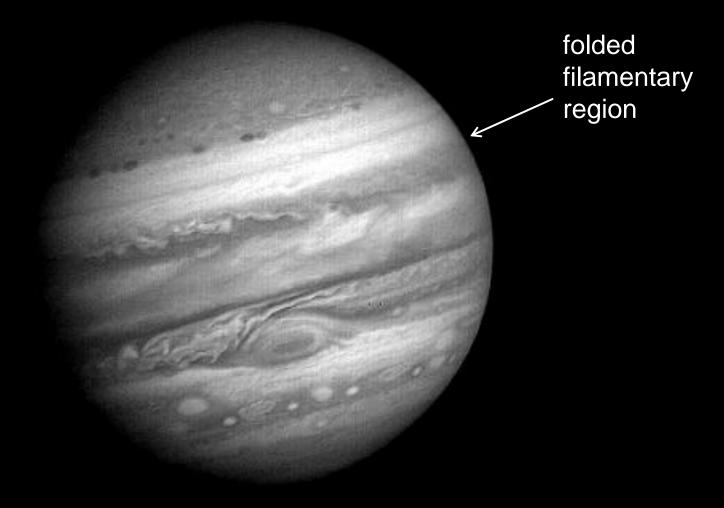
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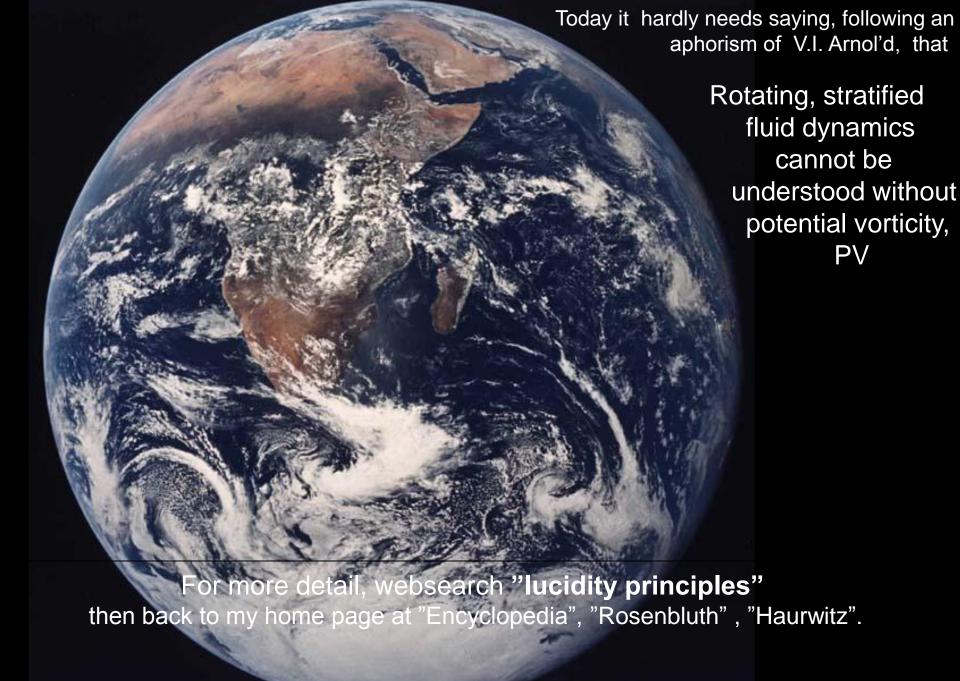
And rotation and stable stratification are important in many cases, including the Sun's interior...





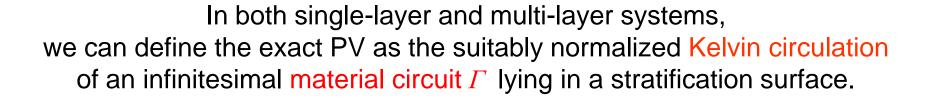
1979: Voyager 1 approaching (60 Jupiter days) - unearthly!



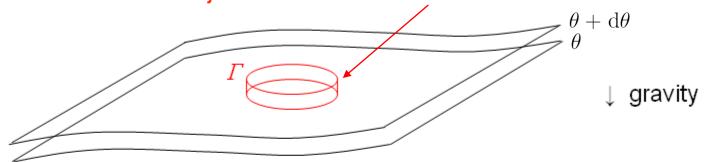




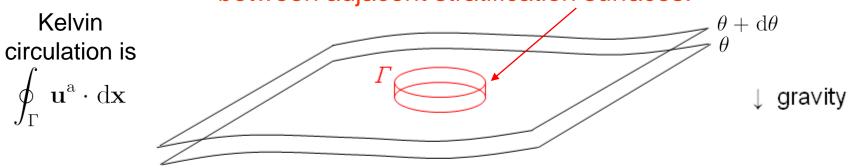




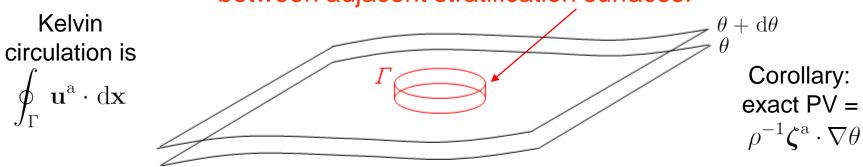
"Suitably normalized": multiply by $d\theta$ / mass of pillbox between adjacent stratification surfaces:



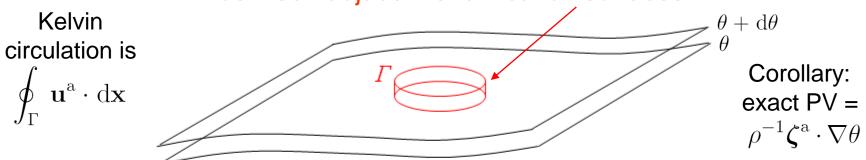
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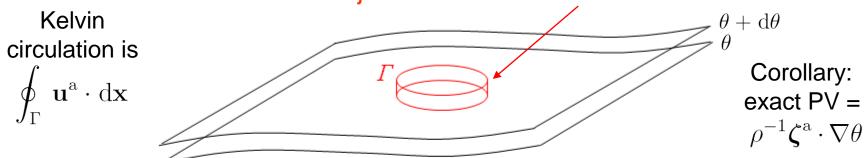


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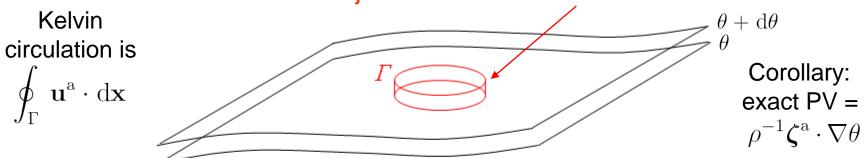
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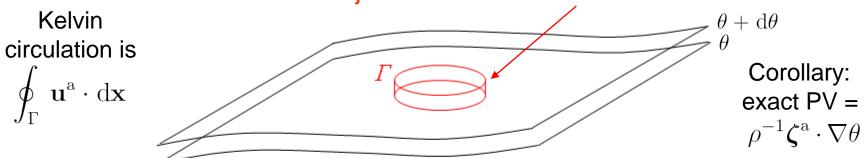


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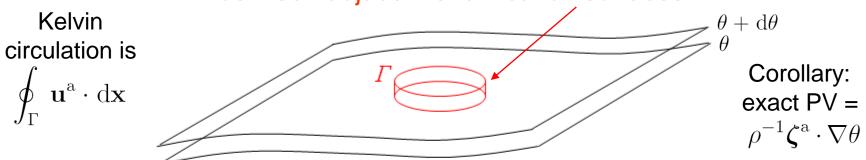
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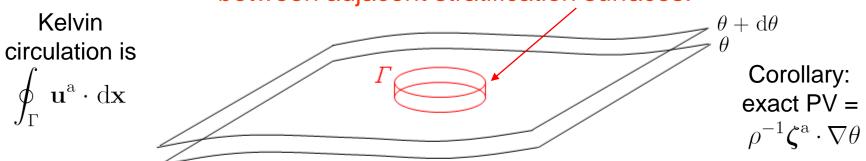
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Note scale effect: small-scale PV anomalies → weak velocities.

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What do real PV fields look like? Here's an example (nostalgic for me):

theorem etc

McIntyre and Palmer (1983), revisited

PV on the 850K stratification surface:

Breaking planetary waves in the stratosphere

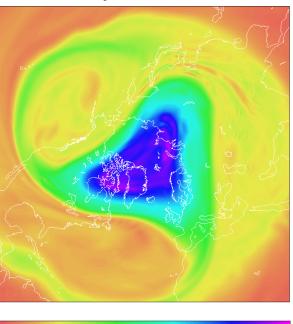
M. E. McIntyre* & T. N. Palmer*

* Department of Applied Mathematics and Theoretical Physics, University of Cambridge, Cambridge CB3 9EW, UK
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Fig. 2 Coarse-grain estimates of Ertel's potential vorticity Q on the 850 K isentropic surface (near the 10-mbar isobaric surface) on 17 (a) and 27 (b) January 1979, at 00 h GMT. The southernmost latitude circle shown is 20° N; the others are 30° N and 60° N. Map projection is polar stereographic. For units see equation (5) onwards. Contour interval is 2 units. Values greater than 4 units are lightly shaded, and greater than 6 units heavily shaded.

Initial state

Potential vorticity at 850K 00UTC 1979/01/17



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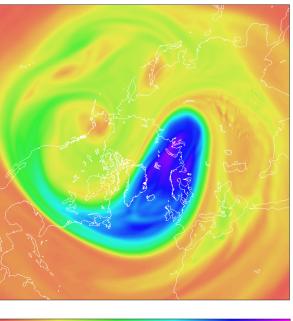
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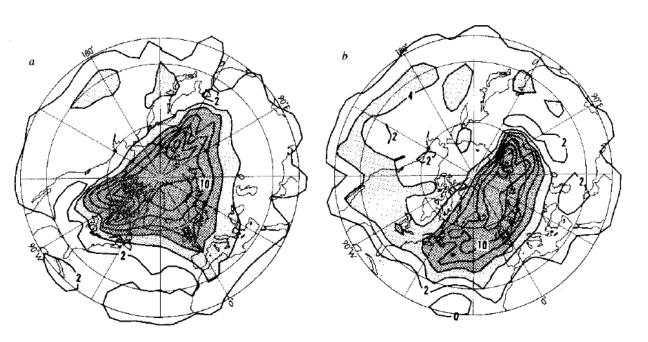
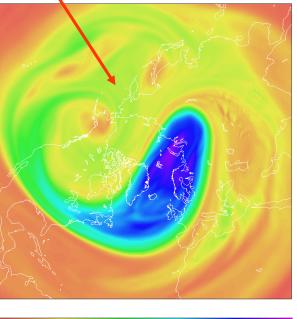


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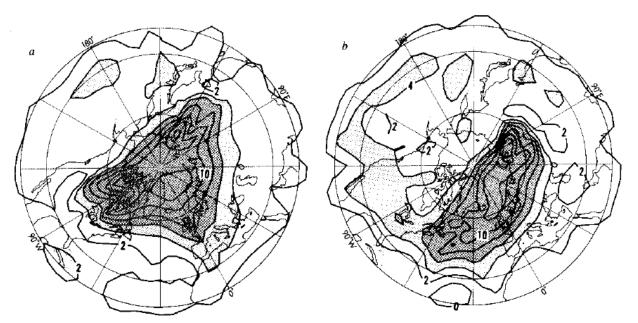
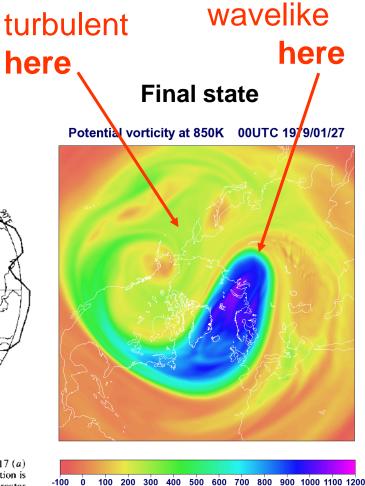
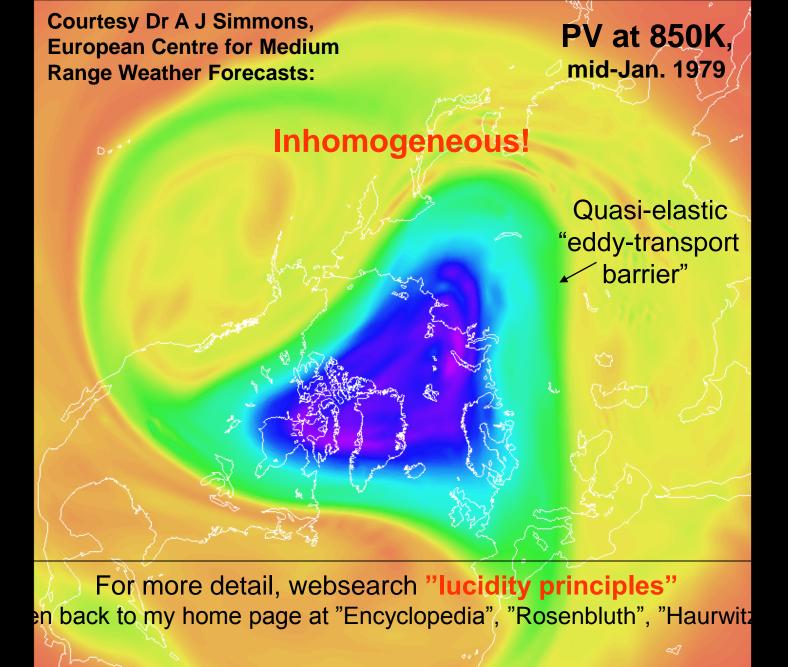
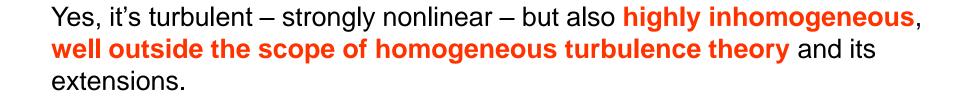


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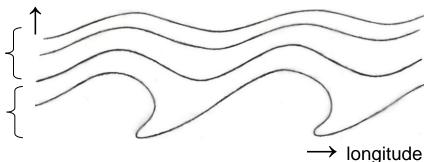
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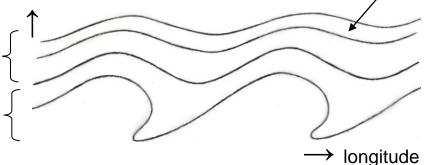


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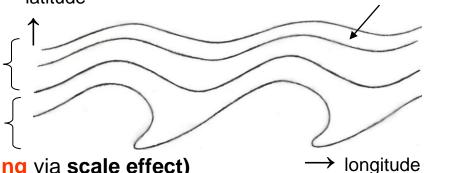
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The simplest **fully self-consistent** model of the structure is the **Stewartson-Warn-Warn (1978) nonlinear critical-layer theory.** And the essence of it was first recognized by Dickinson (1969).

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(Fig 5 from my 1982 review, J. Met. Soc. Japan 60)

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For more detail, websearch "lucidity principles" then back to my home page at "Encyclopedia", "Rosenbluth", "Haurwitz".

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The idea goes back to the work of G. I. Taylor and others on turbulence – thinking about **momentum transport** versus **vorticity transport**. Taylor had already found an **important clue in 1915**, showing that one pattern of vorticity transport can imply a **different pattern** of momentum transport:

 nonlinearly relates eddy fluxes of PV to momentum-flux divergences:



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NB: nonlinear relation: valid at any amplitude!! And valid regardless of whether motion is free, forced, or self-excited.

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Here are 3 more reasons

1. Stagnation effect: For S spanning the globe, easy to prove that homogeneous mixing has an **absurd consequence.** We have

$$\iint_{\mathcal{S}} P \sigma \, \mathrm{d}A = 0 \qquad \text{where } P = \mathbf{exact} \ \mathsf{PV},$$
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on each S. Homogeneous mixing (e.g., constant-diffusivity models) $\rightarrow P = 0$ everywhere on S: inverts to **absolute stagnation!**

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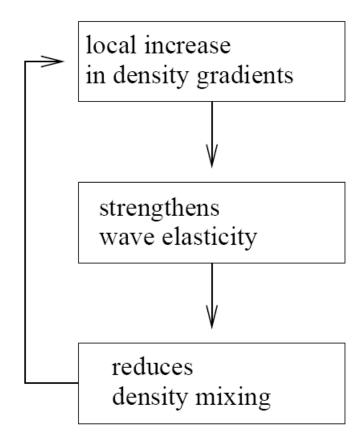
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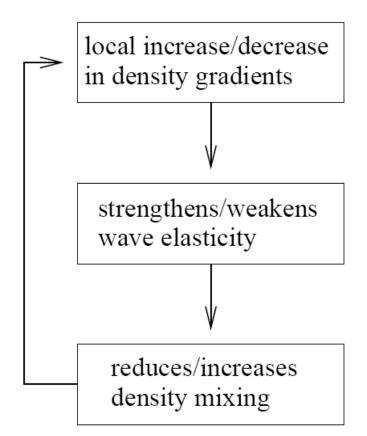
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- 3. PV Phillips effect: a generic positive-feedback effect:

O. M. Phillips (1972 Deep Sea Res).

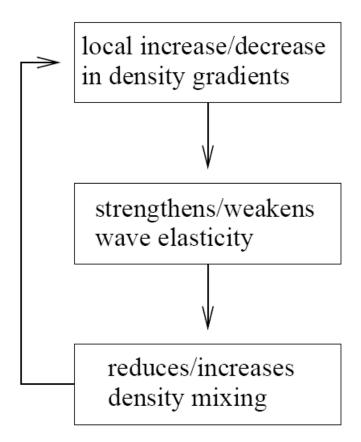
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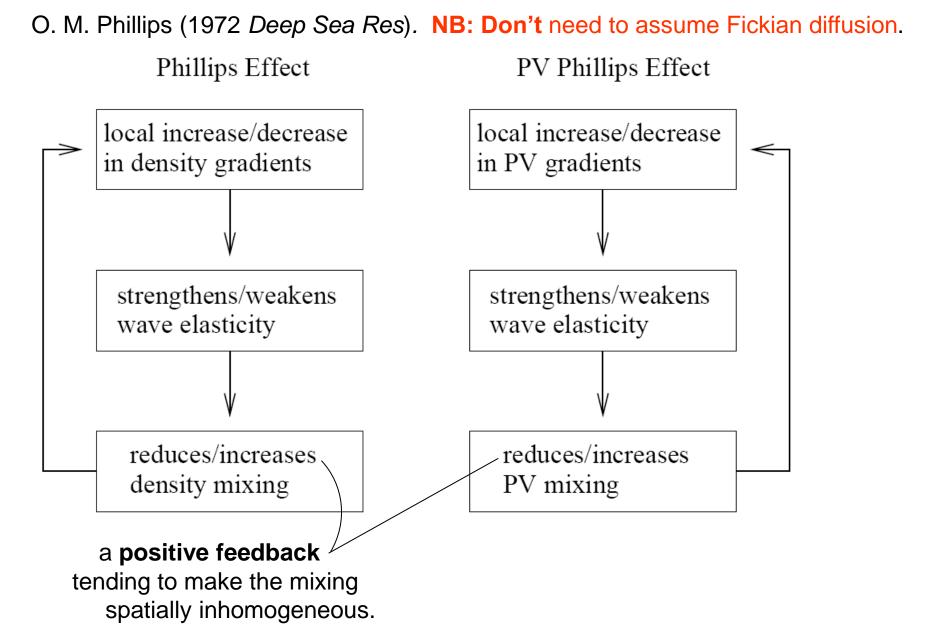


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PV mixing

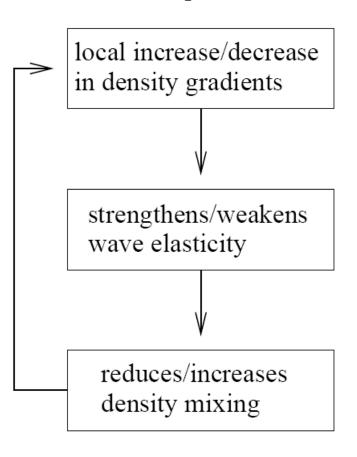
density mixing

[&]quot;staircase" on my home page



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O. M. Phillips (1972 Deep Sea Res). NB: Don't need to assume Fickian diffusion.



Phillips Effect

a **positive feedback** tending to make the mixing spatially inhomogeneous.

local increase/decrease in PV gradients

strengthens/weakens

PV Phillips Effect

reduces/increases
PV mixing

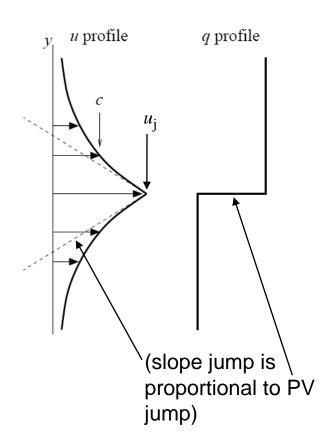
wave elasticity

The more inhomogeneous, the stronger the feedback, bringing in the **shear** effects (Juckes & M, *Nature* 1987).

[&]quot;staircase" on my home page

Simplest quasi-geostrophic **strong-jet model** (Rosenbluth Lecture & refs.). PV contours **infinitely bunched up** to make a **perfect Rossby waveguide**:

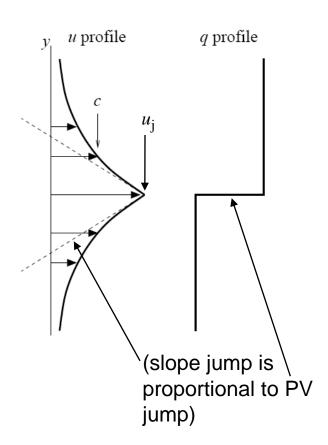
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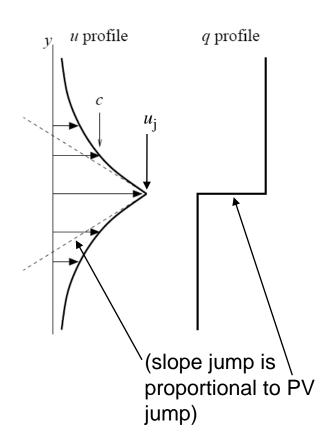


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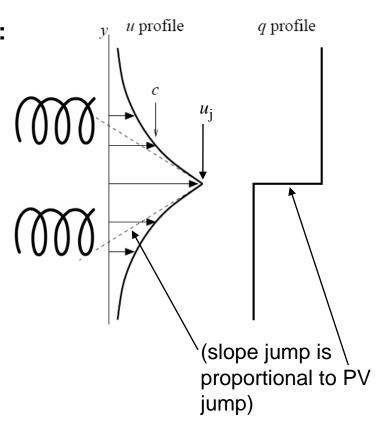
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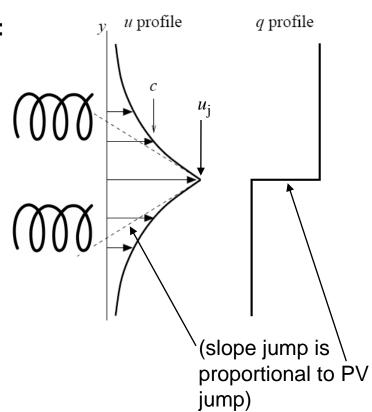
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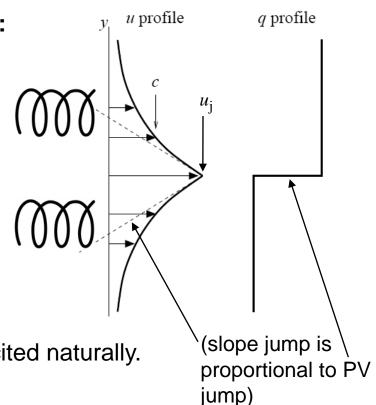
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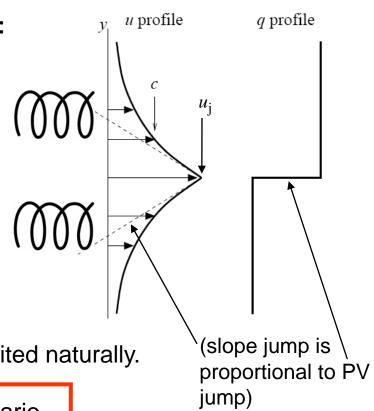
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q profile u profile (slope jump is proportional to PV jump)

Strong support comes from numerical and laboratory experiments, e.g.,

Scott and Dritschel (2012, *J. Fluid Mech.* **711**, 576) – **forced-dissipative** but clarifying how to reach the **more** "**natural**" low-excitation, low-dissipation parameter regimes.

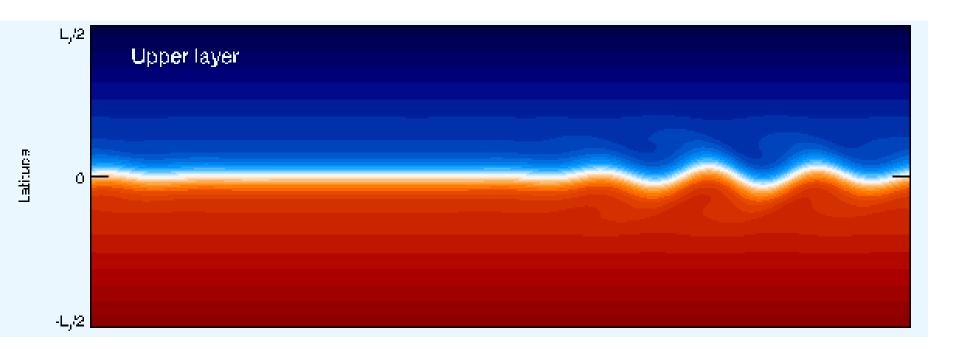
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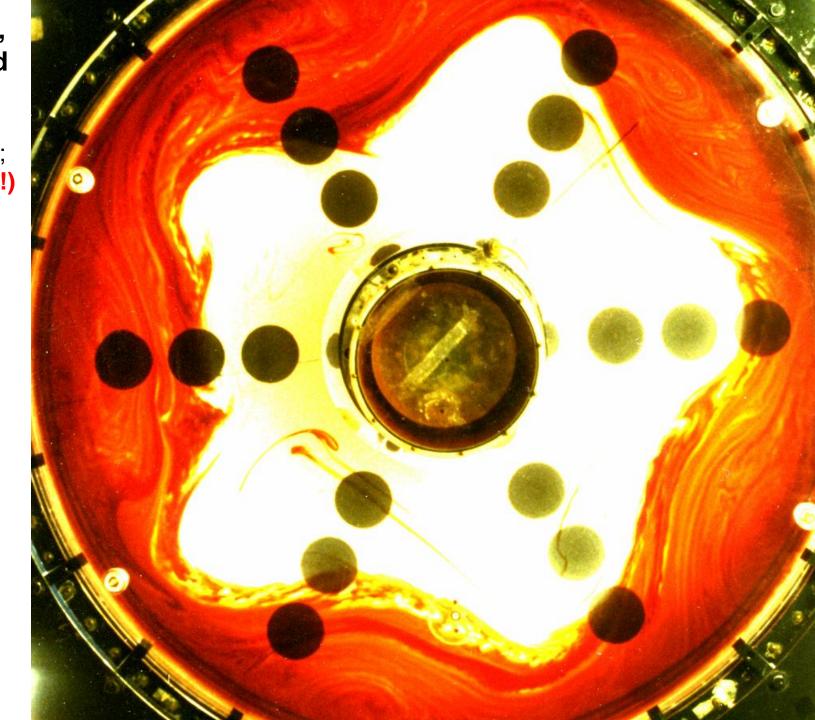
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Here the Taylor identity is satisfied via a **form stress** exerted from below:



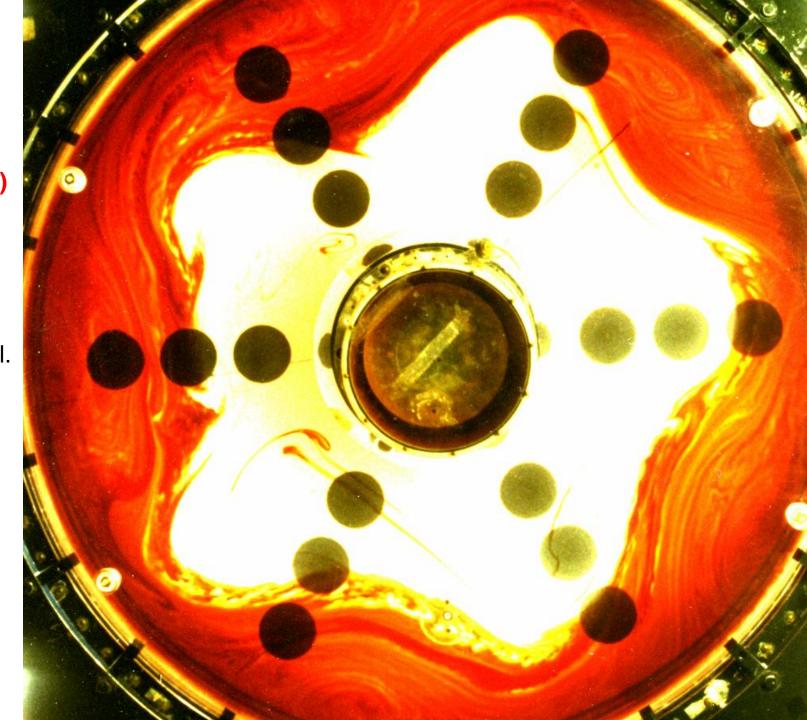
Next is a classic lab experiment conveying the same message:

Sommeria, Myers, and Swinney, Nature 1989 86.4 cm dia.; 3 revs/sec (!)



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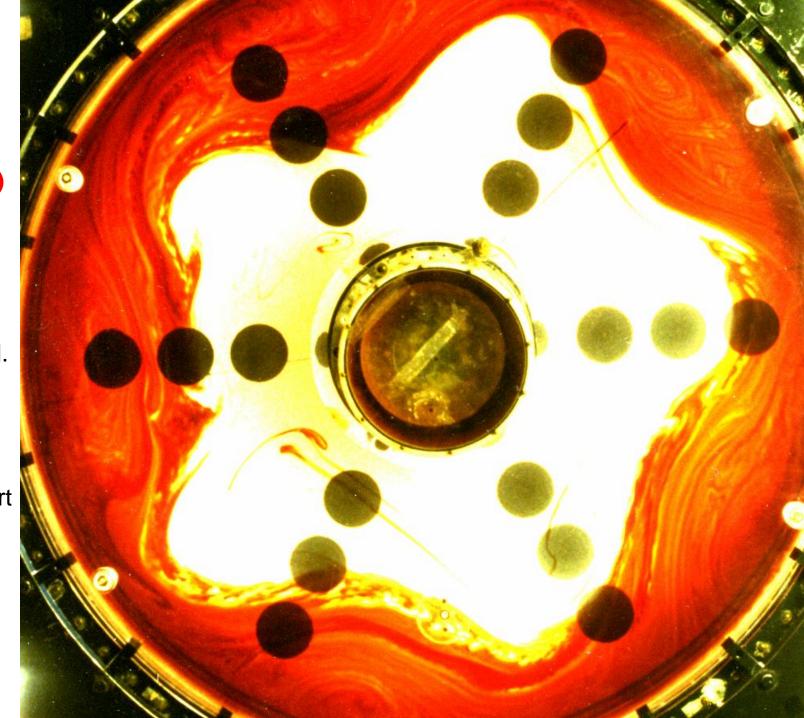
PV map and dye map near-identical.



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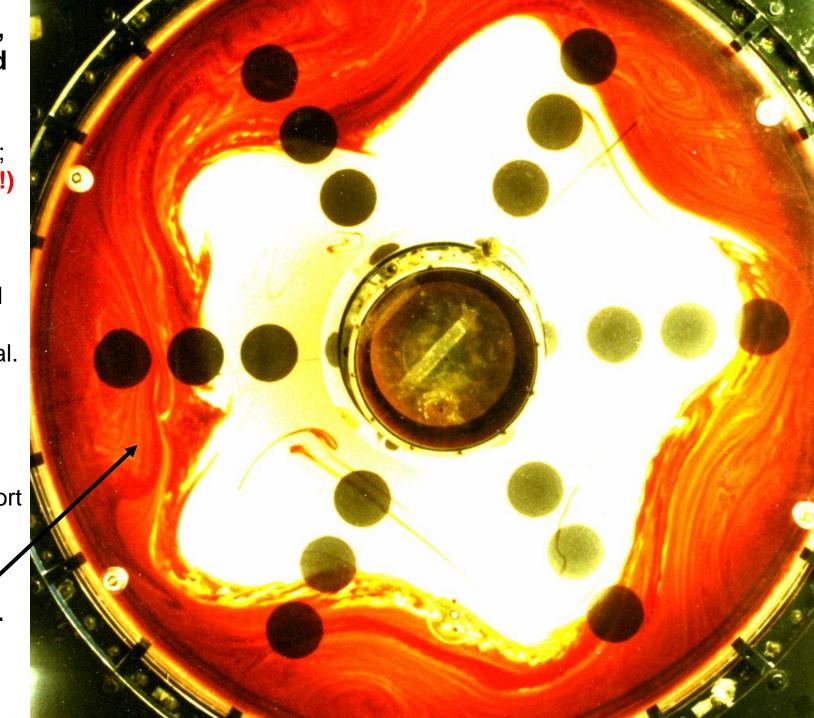


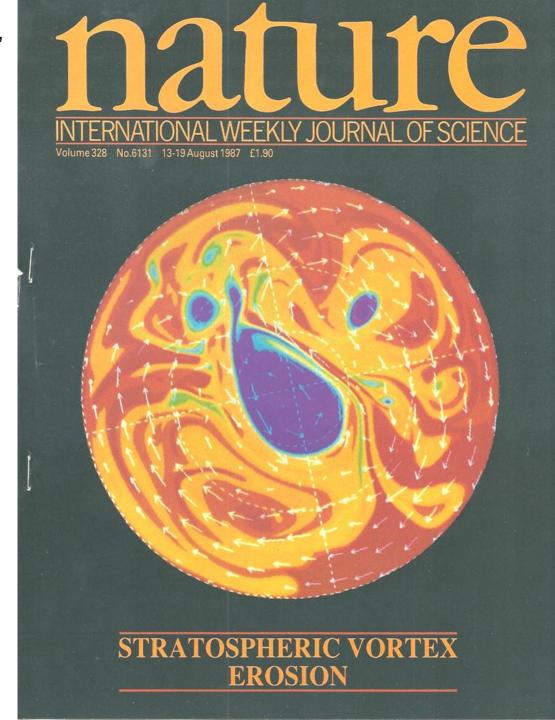
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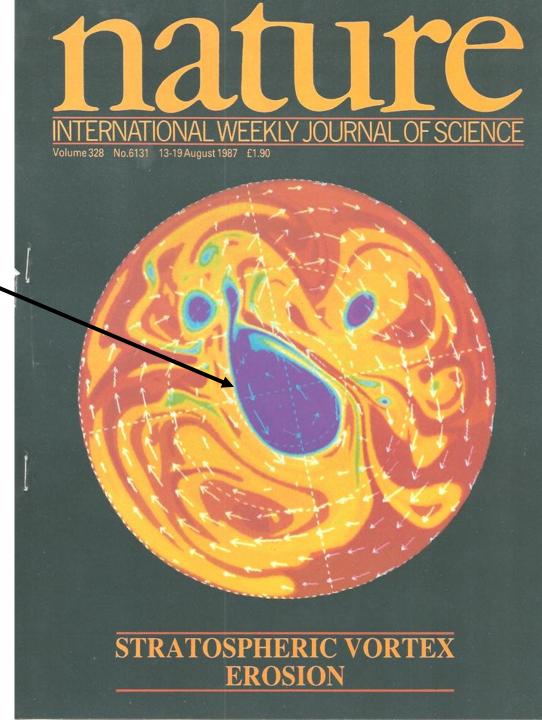
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Dye put in > 500 revs. earlier.





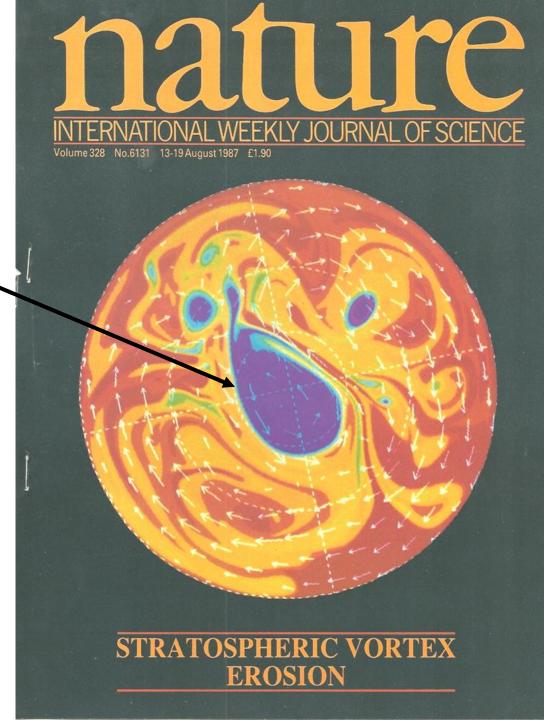
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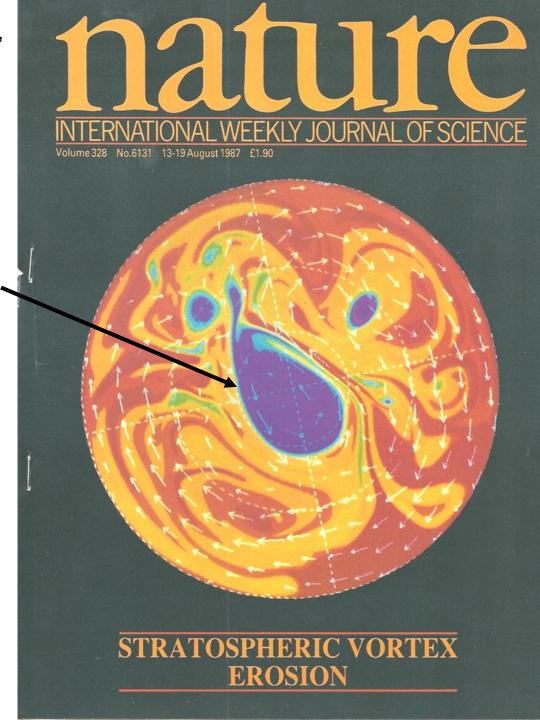


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Many other examples (e.g. nice observational work in Huw Davies' group. So here's the bottom line:



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The paradigm is *not*, of course, the Answer to Everything:

- PV doesn't always mix (e.g., vortex merging)
- Not all jets are strong jets (e.g., "ghost jets" in the Pacific Ocean)
- And there are, of course, other nonlinear mechanisms that show up in different thought-experiments:

Jet dynamics – a complex conceptual landscape wth a 2-level hierarchy of ideas:

1. Generic ideas:

PV Phillips effect PV invertibility

Taylor identity:

 $\overline{v'q'} = - \text{div (eddy momentum flux)}$

Nonlinear, forced/free/self-excited

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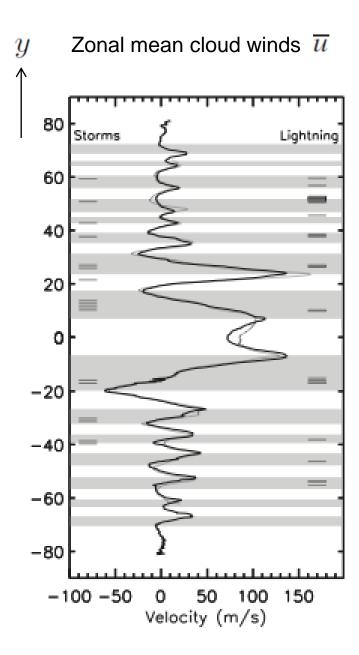
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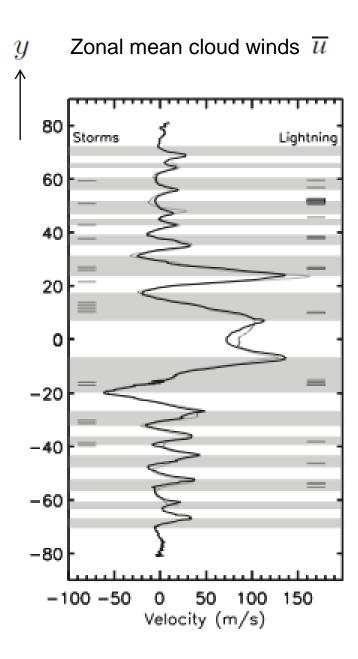
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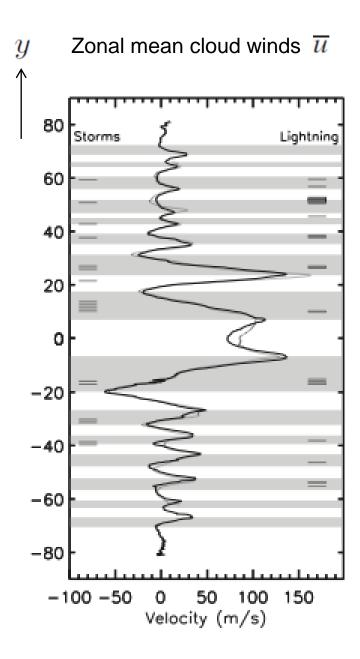
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- (vii) **PV-biased forcing:** critical in the **idealized Jupiter weather-layer model** being studied by Stephen Thomson and myself.



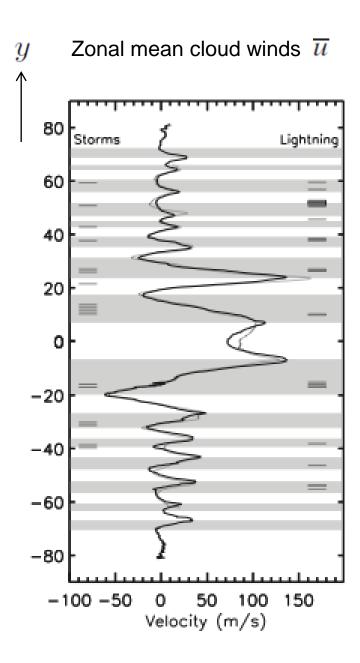


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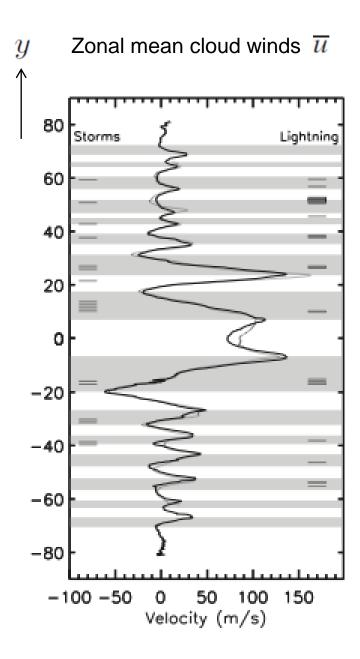
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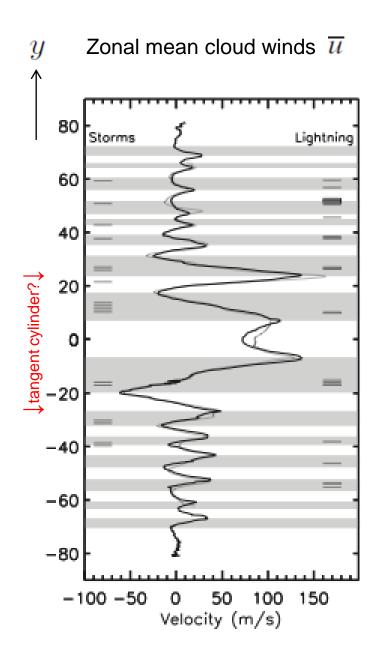
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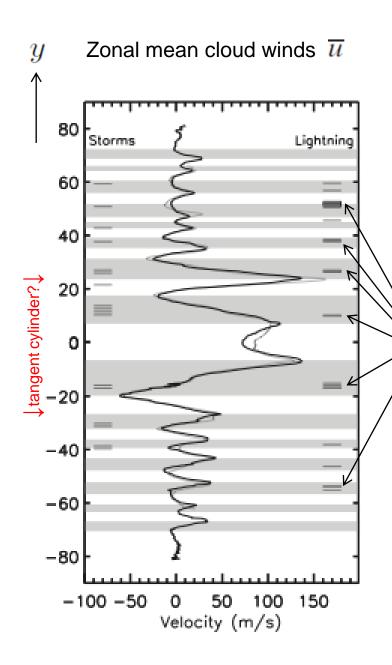
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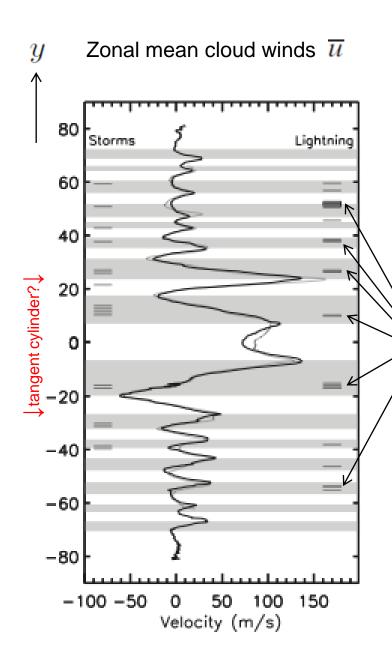


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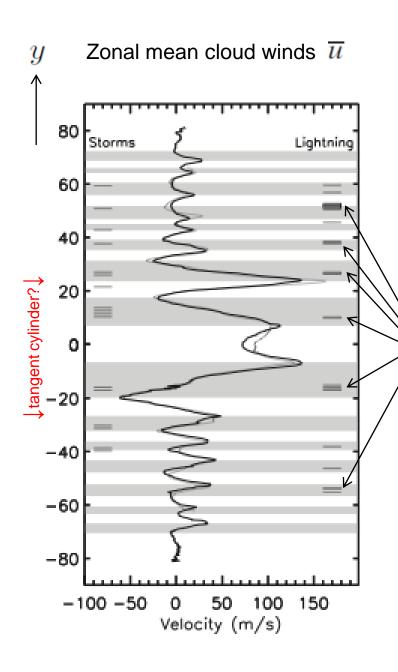
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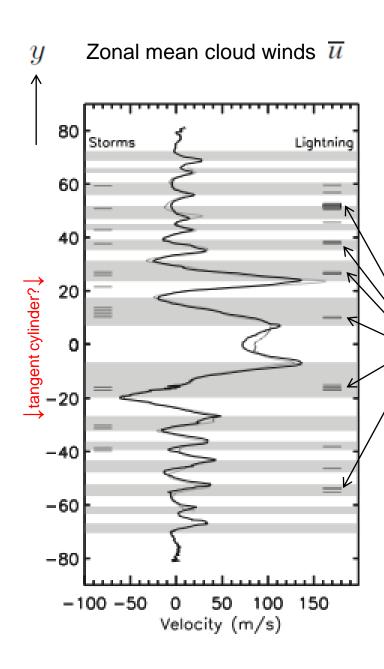
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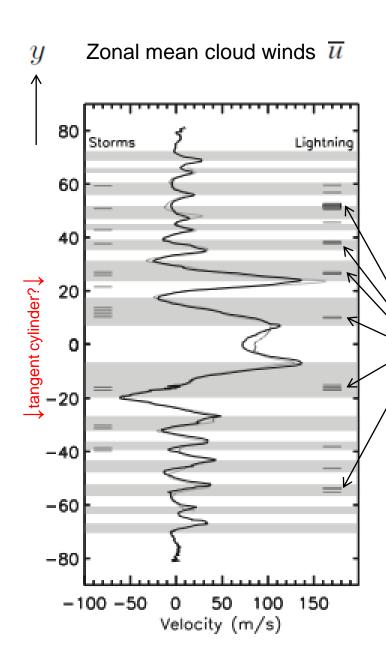
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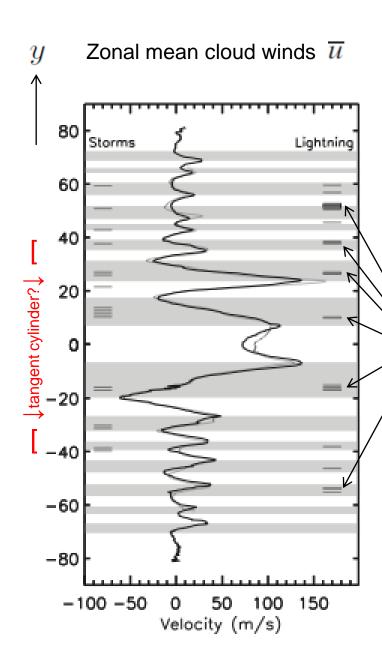
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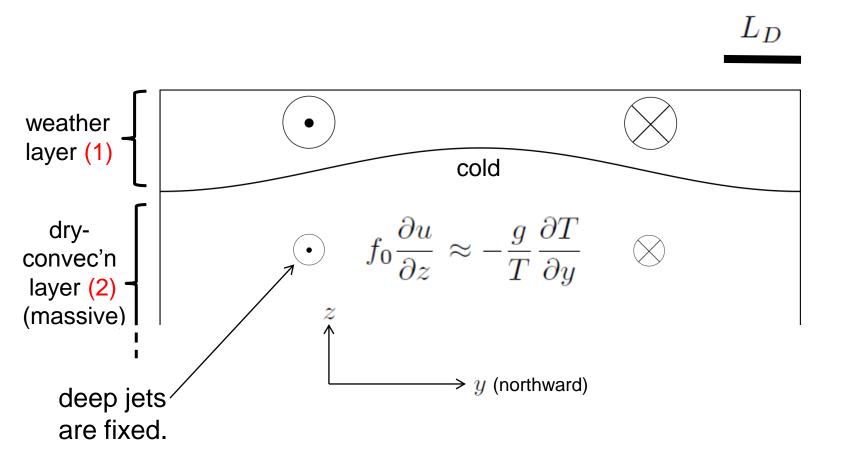
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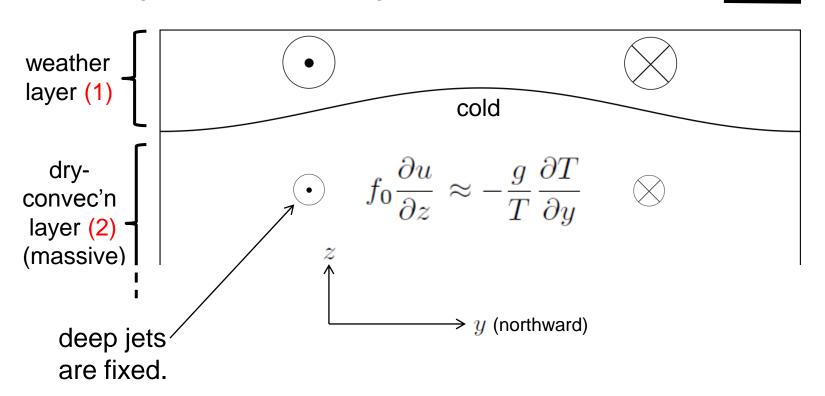
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Start by taking Dowling-Ingersoll '89 seriously,

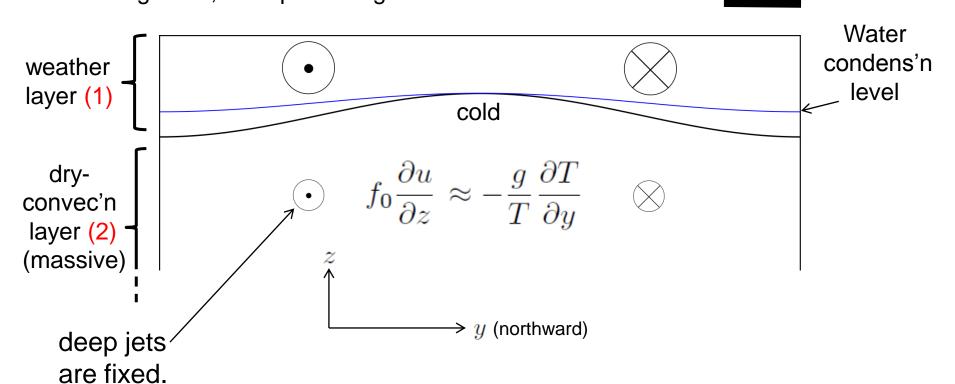


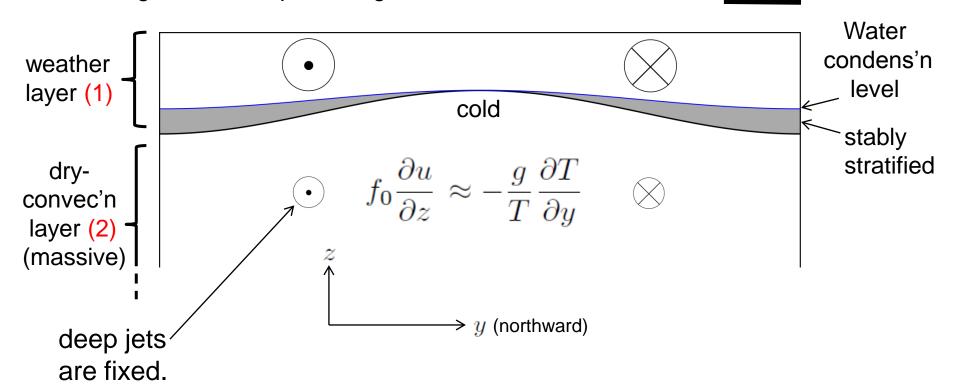
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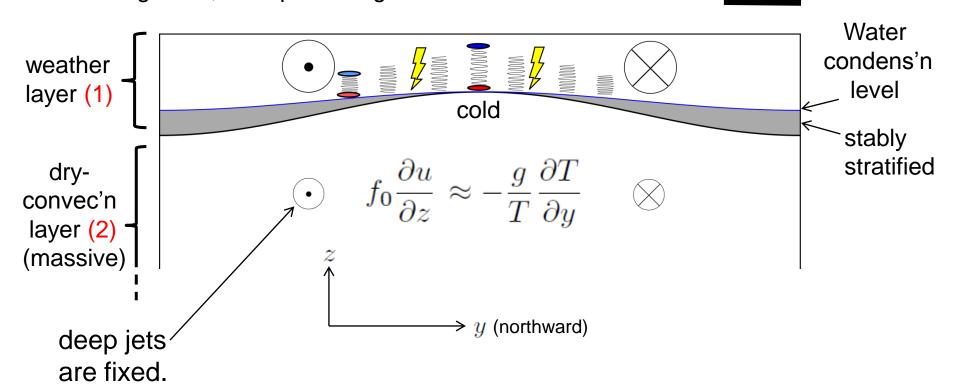


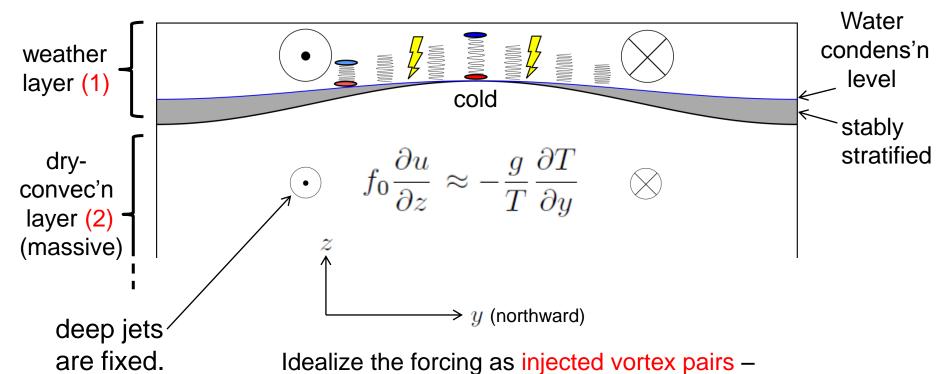


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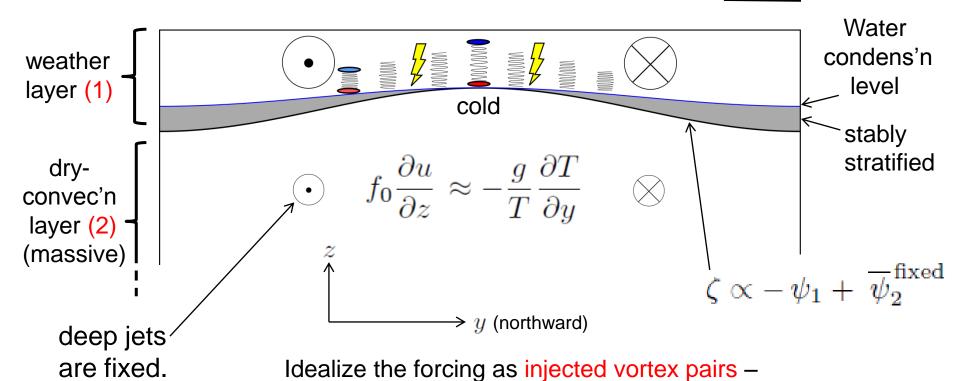






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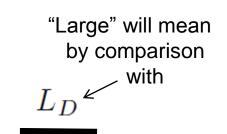
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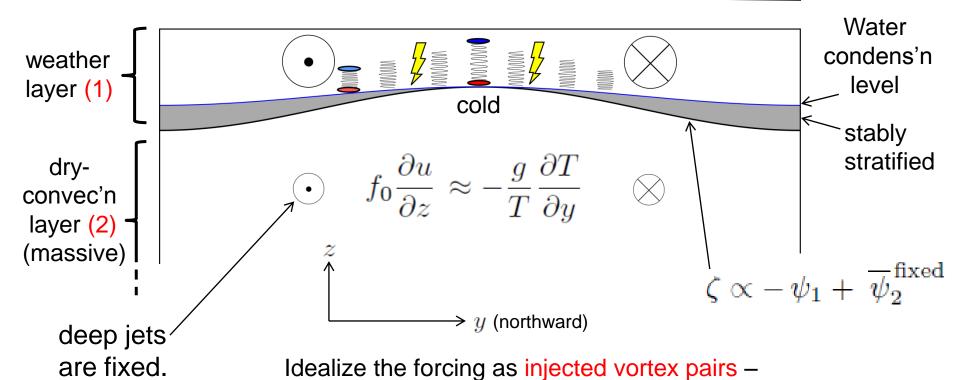


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Note domain-wide "complementary forcing"

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(Realistic radiative damping timescales? *Or* is concept irrelevant?)

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What happens is sensitive to $v'_1q'_1$ & hence to the strength of S:

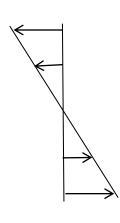
Forcing strengths $|q_{1 \, {
m max}}|$, cf. Shigeo Kida (1981):

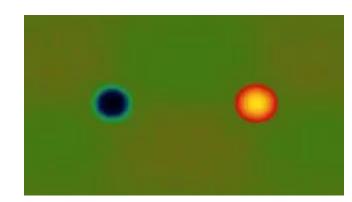
Strong injections $|q_{1 \text{ max}}| = 16$: in units of background shear

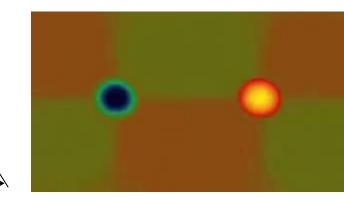


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Semi-strong injec'ns $|q_{1 \max}| = 8$:

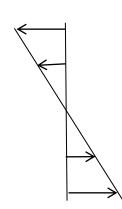


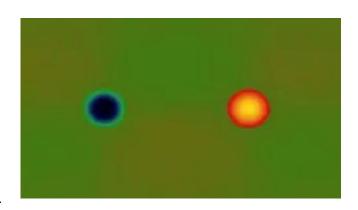




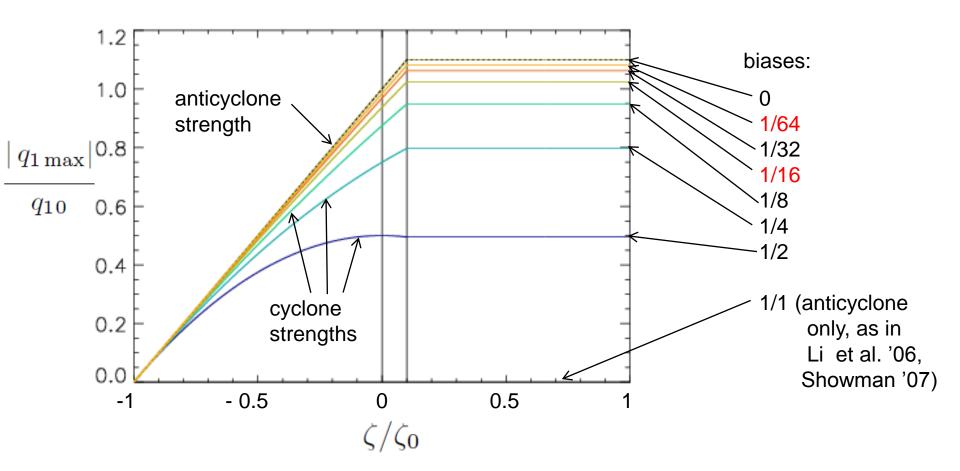


(as required for passive Kelvin shearing)



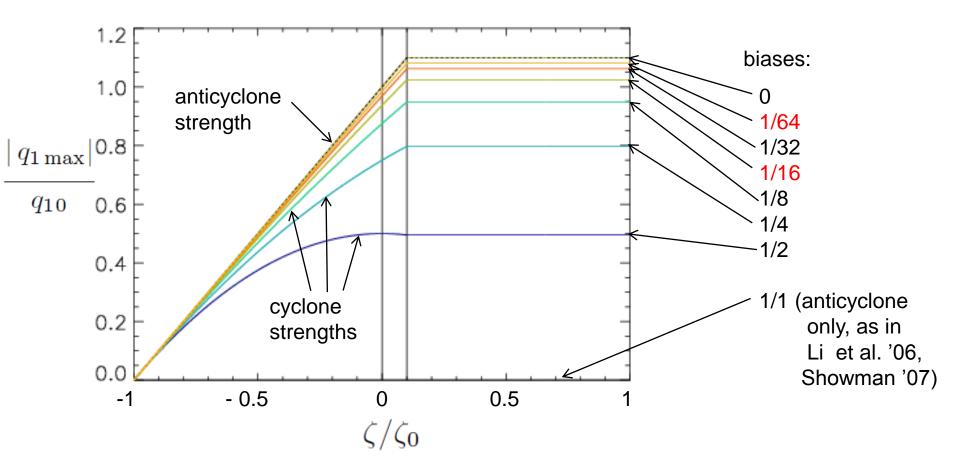


At present we use the following injection-pair algorithm, with "saturation":



"Saturation" constraint prohibits q_1' from exceeding 1.1 q_{10} .

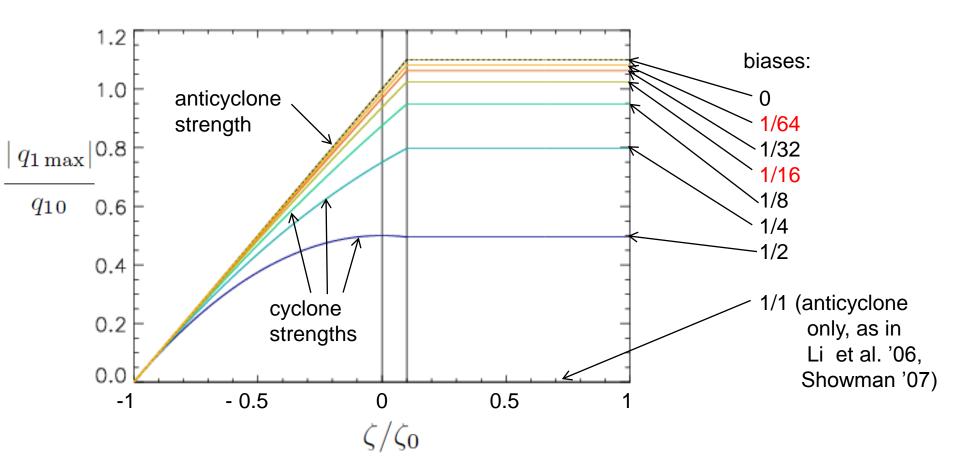
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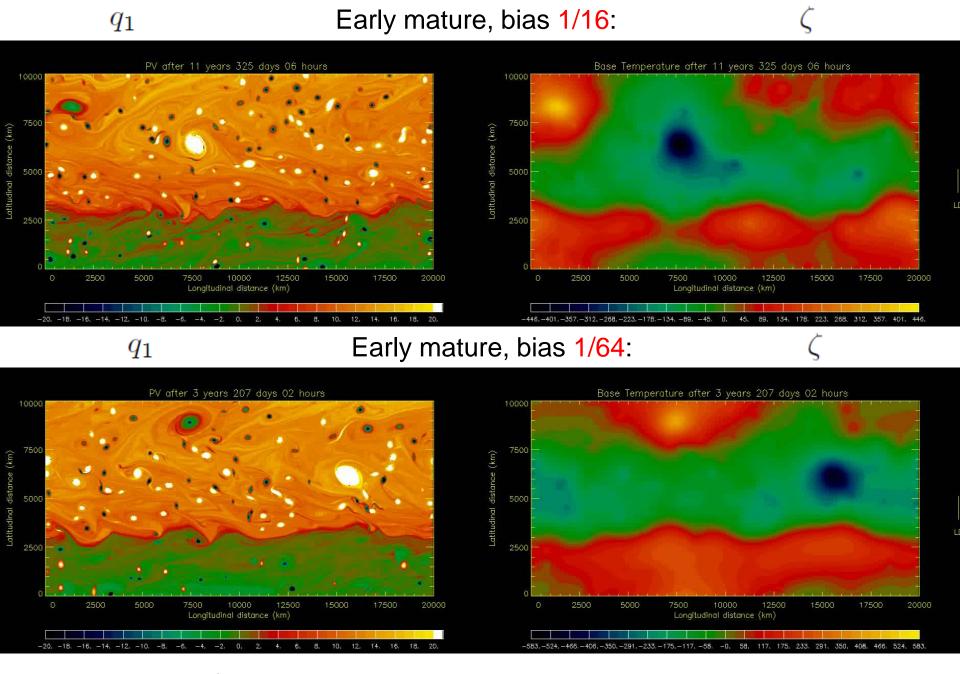
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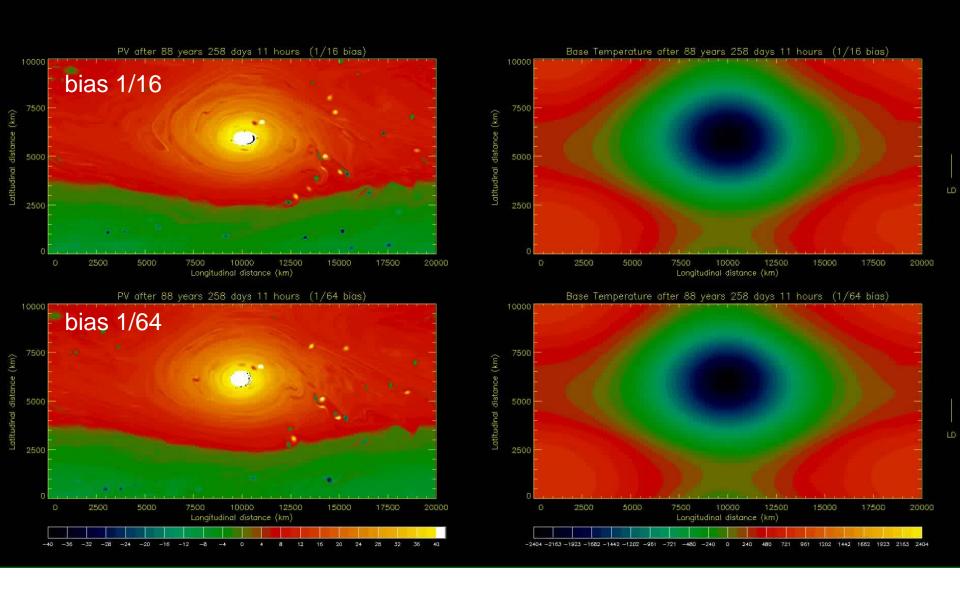
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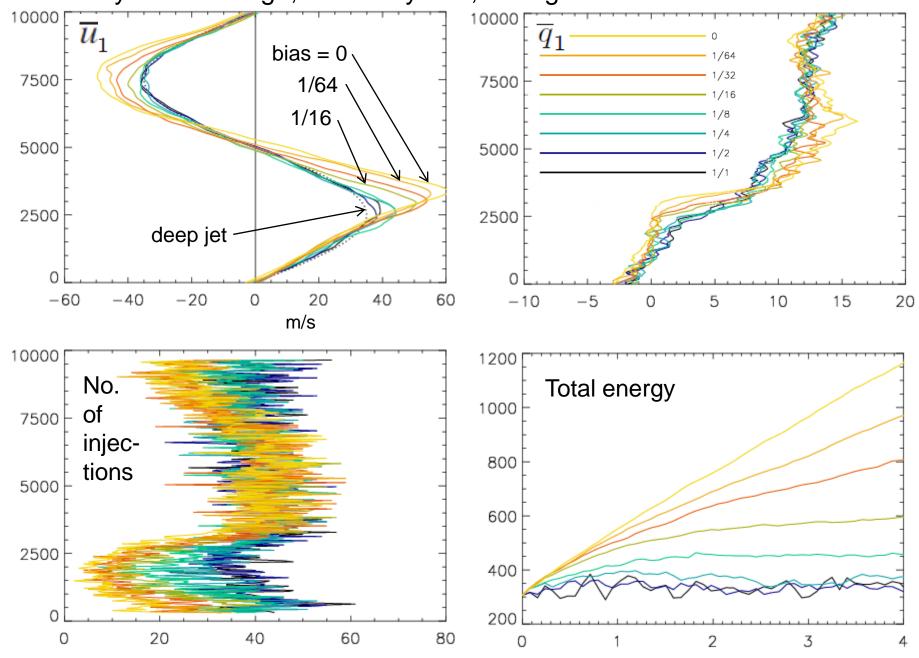
Bias 1/64 is too weak to stop the large cyclone from growing further:



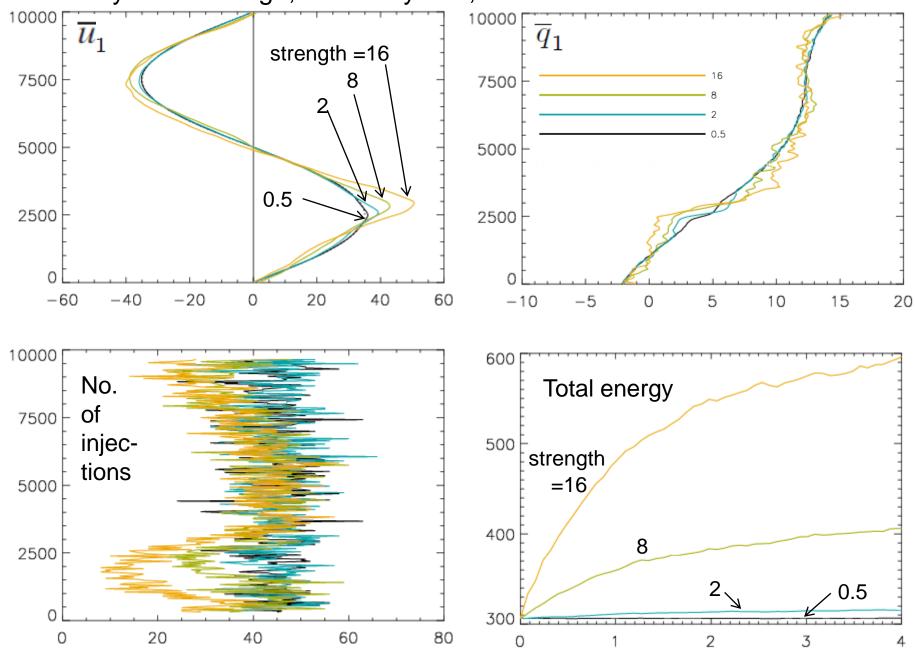


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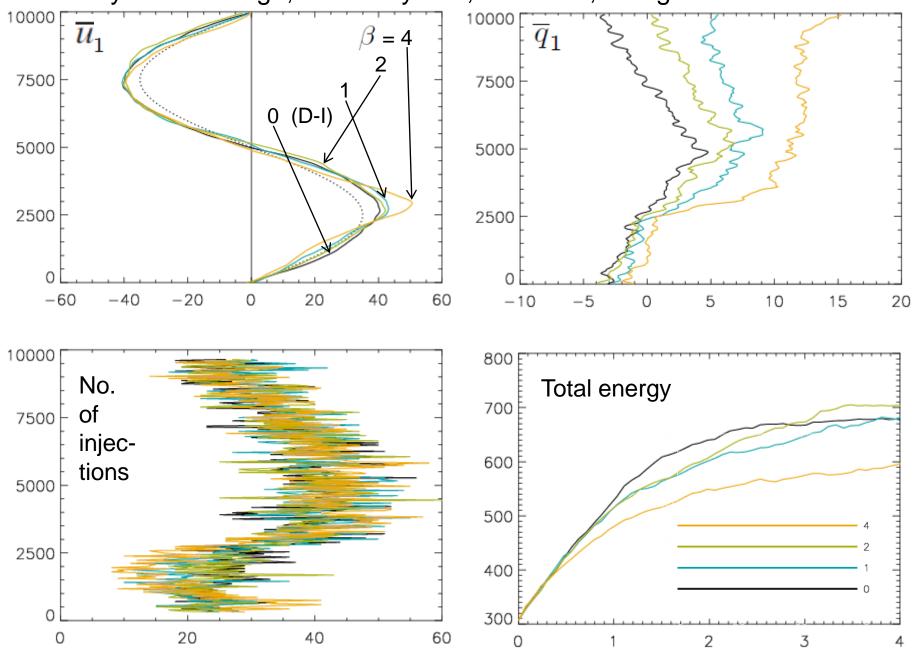
Early mature stage, 4 Earth years, strength still 16:



Early mature stage, 4 Earth years, bias 1/16:



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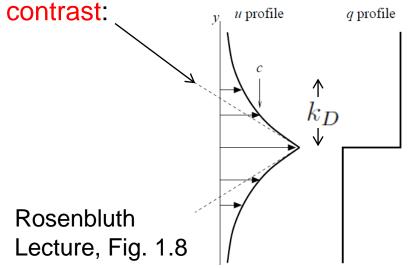
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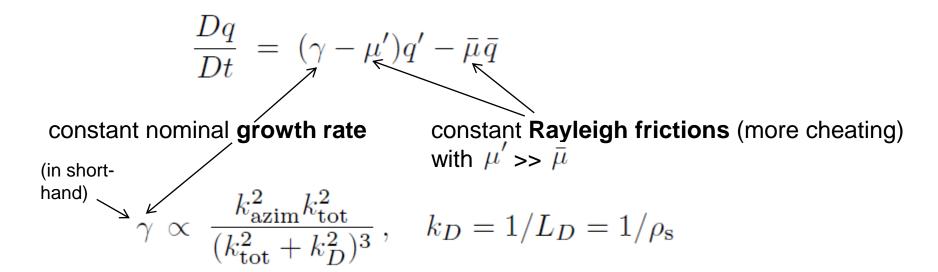
$$rac{Dq}{Dt}=0$$
 with $q=
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Prime is crucial! (Tilde, if you prefer – departure from zonal average.)
Implication: strong jet with given
PV contrast has far greater velocity



Add self-excitation term as crude model – this is cheating! – of resistive drift-wave instability à la Hasegawa-Wakatani

$$\frac{Dq}{Dt} = (\gamma - \mu')q' - \bar{\mu}\bar{q}$$
 constant nominal **growth rate** constant **Rayleigh frictions** (more cheating) with $\mu' >> \bar{\mu}$
$$\gamma \propto \frac{k_{\rm azim}^2 k_{\rm tot}^2}{(k_{\rm tot}^2 + k_D^2)^3} \,, \quad k_D = 1/L_D = 1/\rho_{\rm s}$$



$$\frac{Dq}{Dt} = (\gamma - \mu')q' - \bar{\mu}\bar{q}$$
 constant nominal **growth rate** constant **Rayleigh frictions** (more cheating) with $\mu' >> \bar{\mu}$ with $\mu' >> \bar{\mu}$
$$\frac{k_{\rm azim}^2 k_{\rm tot}^2}{(k_{\rm tot}^2 + k_D^2)^3} \,, \qquad k_D = 1/L_D = 1/\rho_{\rm S}$$

 L_D is chosen *much* smaller than the domain size (square box now) – less Jupiter-like and more tokamak-like – such that

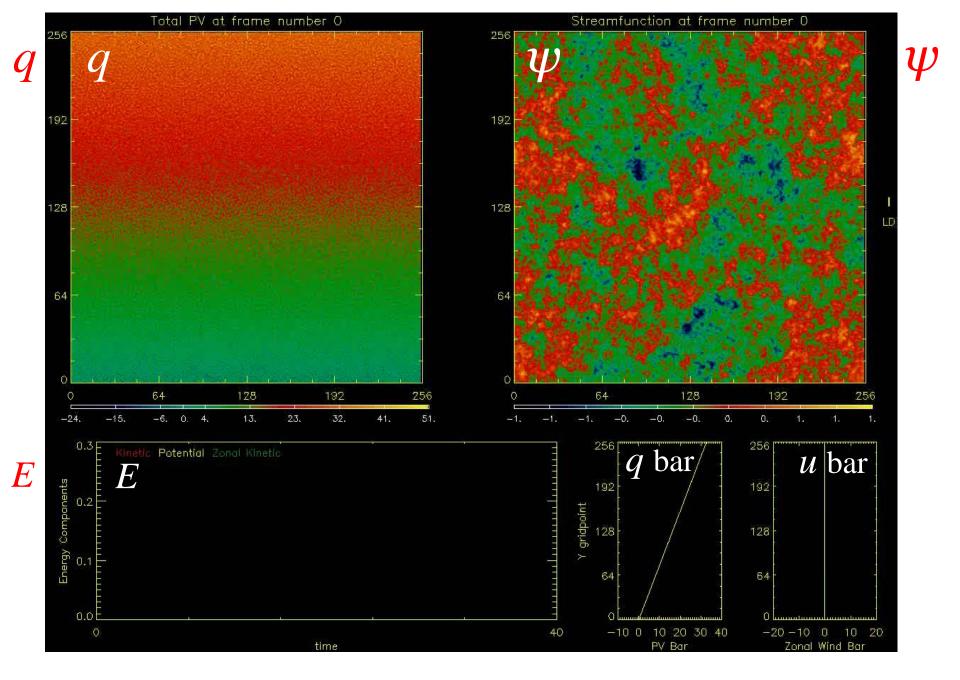
$$\gamma \max \text{ at } 2\pi/k_{\text{azim}} \sim 7, \quad k_{\text{tot}} = k_{\text{azim}}$$

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Get **predator-prey-like situation**, with **hyper-staircasing** of the \overline{q} profile. (My Rosenbluth speculation was **WRONG!!**) *Not* weakly forced/dissipating, PV-mixing-dominated, à la Scott-Dritschel *JFM* 2012.



"Erasmus Darwin held that every so often you should try a damn-fool experiment. He played the trombone to his tulips. This... result... was negative. But other... impudent ideas have succeeded..."

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Set $\beta = 0$, otherwise same problem:

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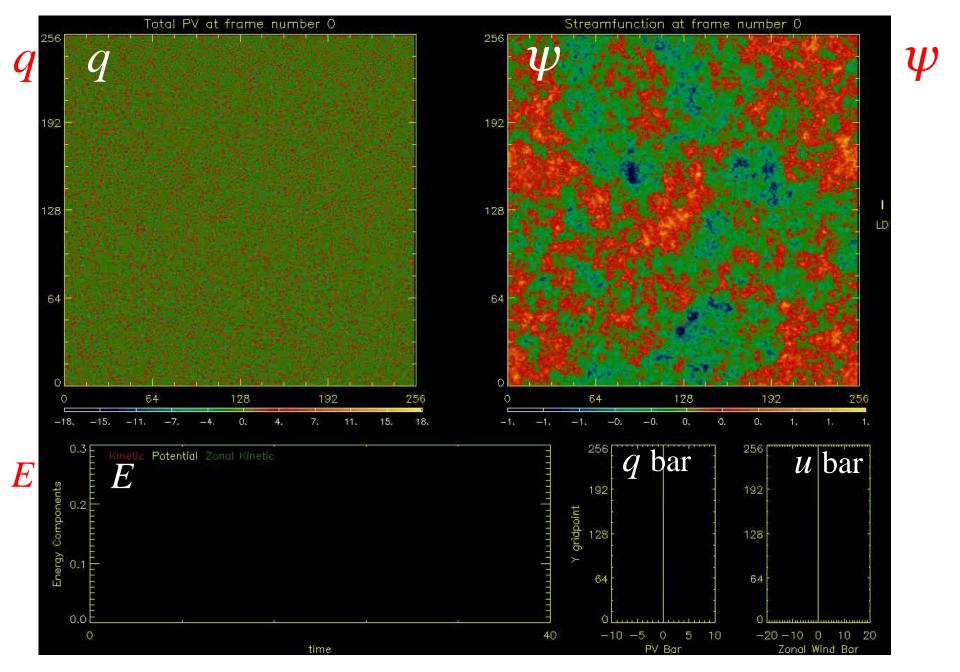


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Conclusions re new work (Jupiter and extended Hasegawa-Mima:

- Jupiter model has no large-scale friction and no radiative damping.
- Despite that, PV bias and zeta-dependence allow statistically-steady
 states but closely tied to the deep jets. Deep jets now essential!!
- Other significant mechanisms in our model include monopole migration, PV mixing (esp. when background beta strong), and cyclone attrition, as well as vortex merging/cannibalism. (Relevance to GRS etc. ???)
- Passive Kelvin (SSST, CE2) almost vanishingly weak despite favourable forcing-anisotropy (Srinivasan & Young).
- EHM results robustly suggest the opposite: Kelvin mech. probably dominant, thanks to perfectly unbiased self-excitation, but easily killed by shear. Reason is the enormously stronger and more extensive shear arising from $k_D = 0$ zonal-mean PV inversion. ("Killed": easy extension of Kelvin 1887.)

For more detail, websearch "lucidity principles" then back to my home page at "Encyclopedia", "Rosenbluth", "Haurwitz".